

Electronic and optical properties of 2D (atomically thin) InSe crystals

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GRAPHENE FLAGSHIP

Zoo of 2D Materials

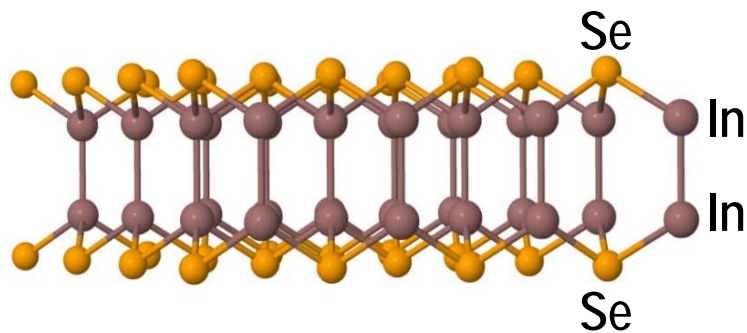
Graphene	hBN 'white graphene'	BCN	Fluorographene	Graphene oxide
MoS ₂ , WS ₂ , MoSe ₂ , WSe ₂	Semiconducting dichalcogenides: MoTe ₂ , WTe ₂ , ZrS ₂ , ZrSe ₂ and so on	Metallic dichalcogenides: NbSe ₂ , NbS ₂ , TaS ₂ , TiS ₂ , NiSe ₂ and so on		
		Layered semiconductors: GaSe, GaTe, InSe, Bi ₂ Se ₃ and so on		
Micas, BSCCO	MoO ₃ , WO ₃	Perovskite-type: LaNb ₂ O ₇ , (Ca,Sr) ₂ Nb ₃ O ₁₀ , Bi ₄ Ti ₃ O ₁₂ , Ca ₂ Ta ₂ TiO ₁₀ and so on		Hydroxides: Ni(OH) ₂ , Eu(OH) ₂ and so on
Layered Cu oxides	TiO ₂ , MnO ₂ , V ₂ O ₅ , TaO ₃ , RuO ₂ and so on			Others

layered substances with covalent bonding within the layers and van der Waals coupling between the layers

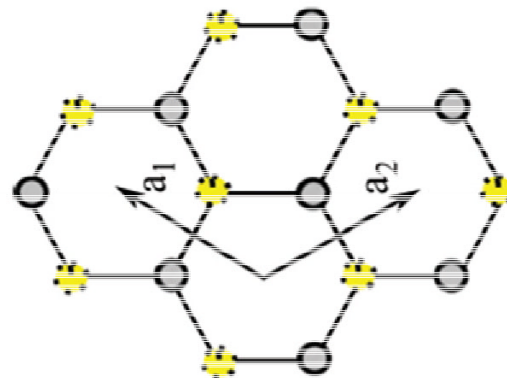
Electronic and optical properties of 2D (atomically thin) InSe crystals

- band and gaps for mono- and few-layer γ -InSe
- optical and transport properties (th+exp)

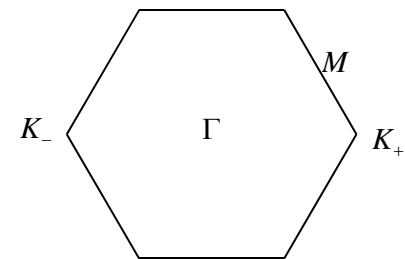
side view



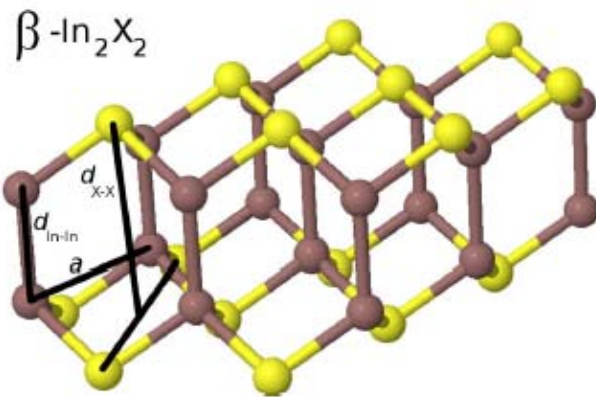
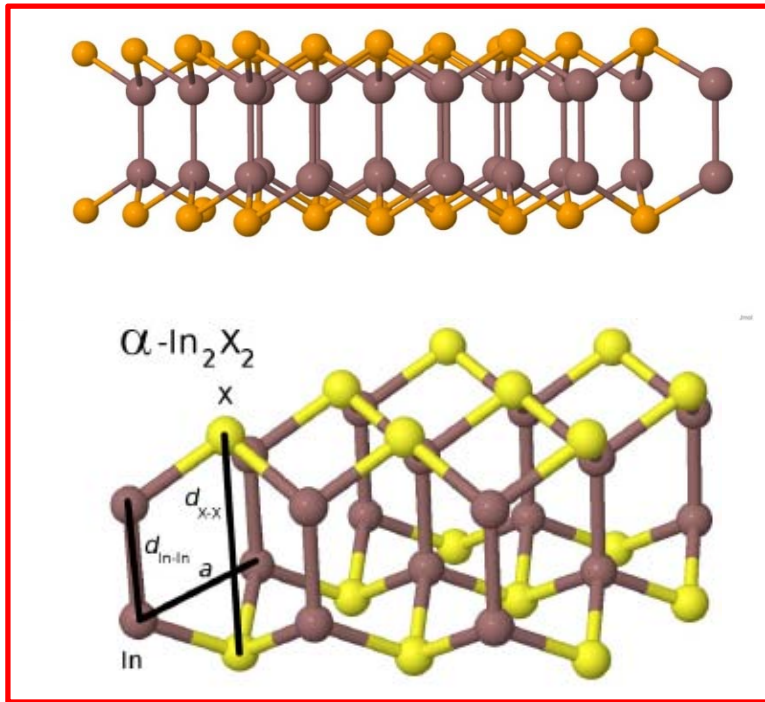
top view



Brillouin zone



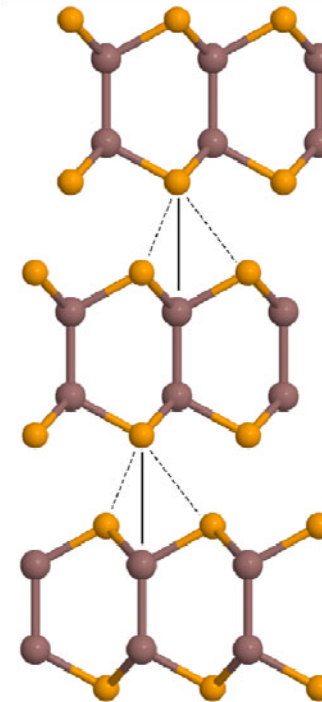
Monolayer (In_2Se_2)



Zolyomi, Drummond, Fal'ko
PRB 89, 205416 (2014)

N-layer $\text{In}_{2N}\text{Se}_{2N}$

$z \rightarrow -z$ mirror symmetry



broken
 $z \rightarrow -z$ mirror
symmetry

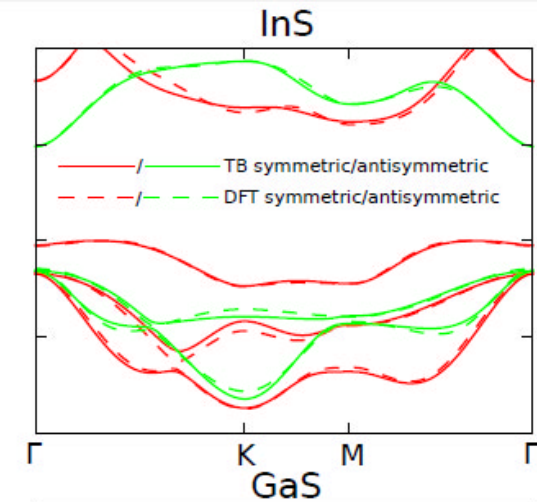
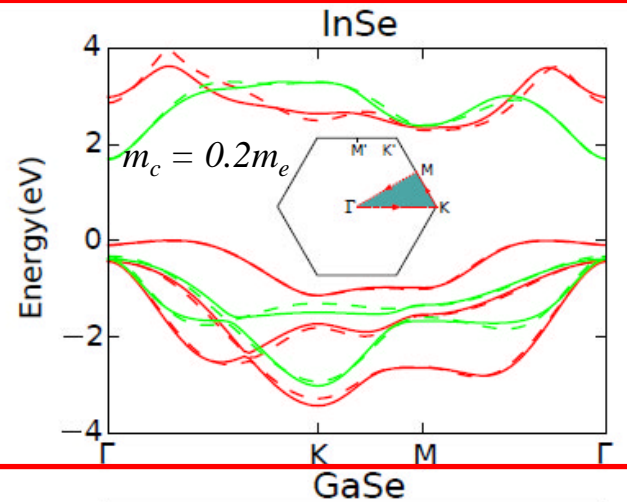
Monolayers In_2S_2 , Ga_2X_2 have qualitatively very similar properties, but multilayer films have different lattice, bands/gaps and selection properties for optical transitions)

Zolyomi, Drummond, Fal'ko
PRB 87, 195403 (2013)

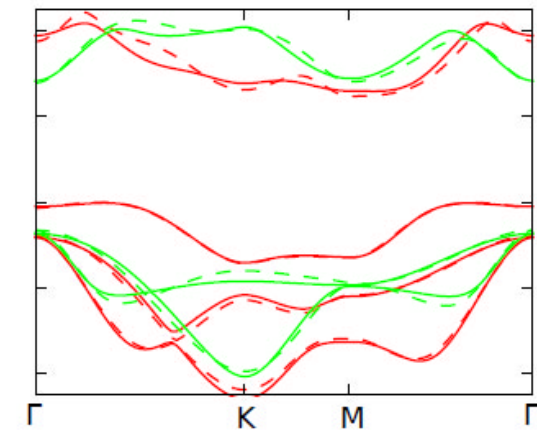
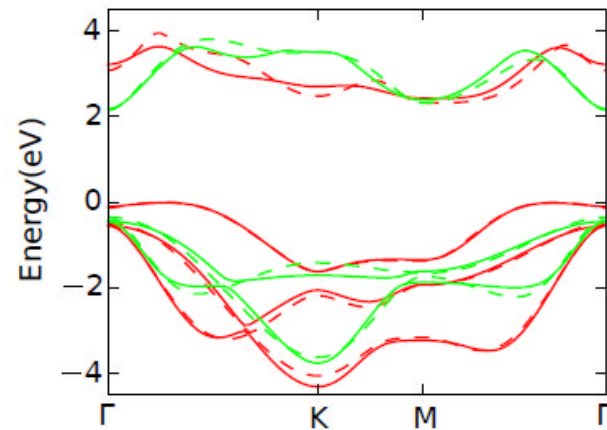
why InSe

DFT bands for monolayer M_2X_2

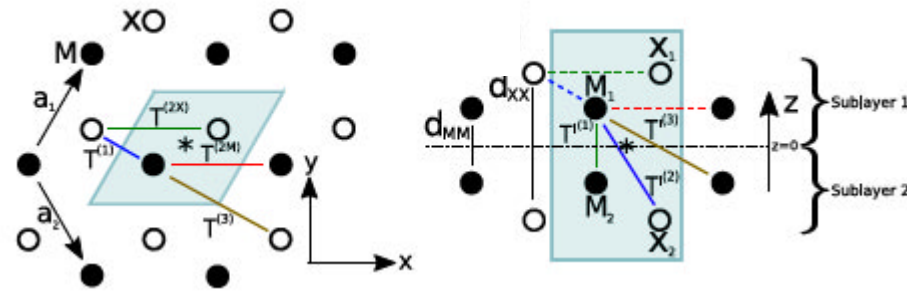
most stable in ambient conditions
bulk γ -InSe



Relatively light electrons in conduction band:
good for high mobility
almost flat edge of the valence band:
opportunity to get strong correlations in hole-doped material

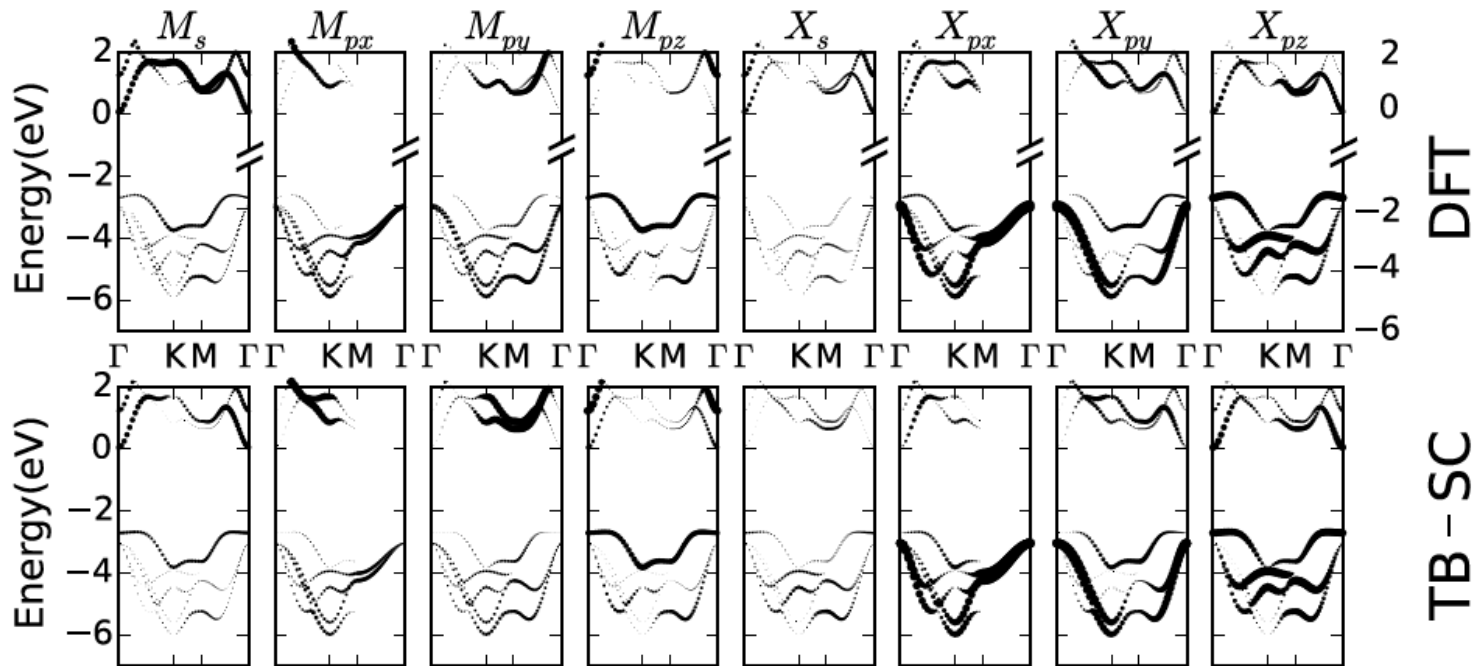
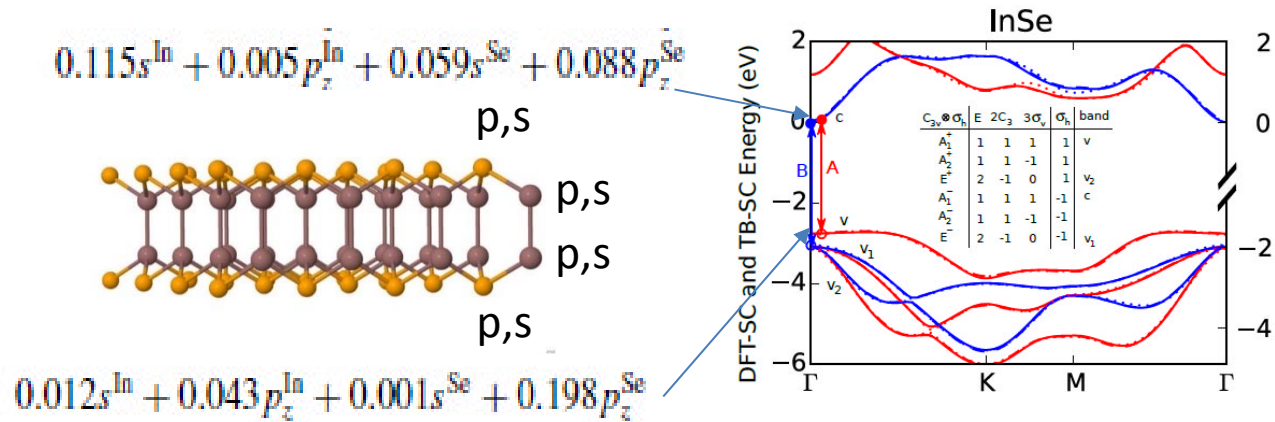


- Formulate tight binding model with all s and p orbitals and all nearest neighbour hoppings (MX, MM, XX) for monolayer (In_2Se_2) and bulk

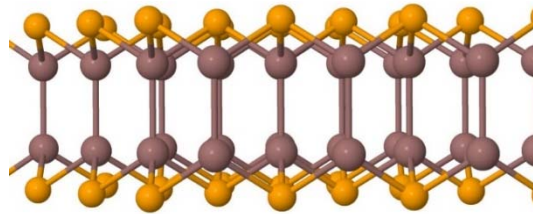


- Fit TB parameters to reproduce DFT bands in monolayer and bulk, after implementing a scissor correction to the band gap determined by comparing experimental and DFT gap values for bulk γ -InSe (1.45eV at low T and 1.25eV at room T)
- Compute spectra of N-layer InSe ($\text{In}_{2N}\text{Se}_{2N}$) and matrix elements for z- and in-plane polarised optical transition, to compare with the experiment

DFT-parametrised tight-binding model for In_2Se_2

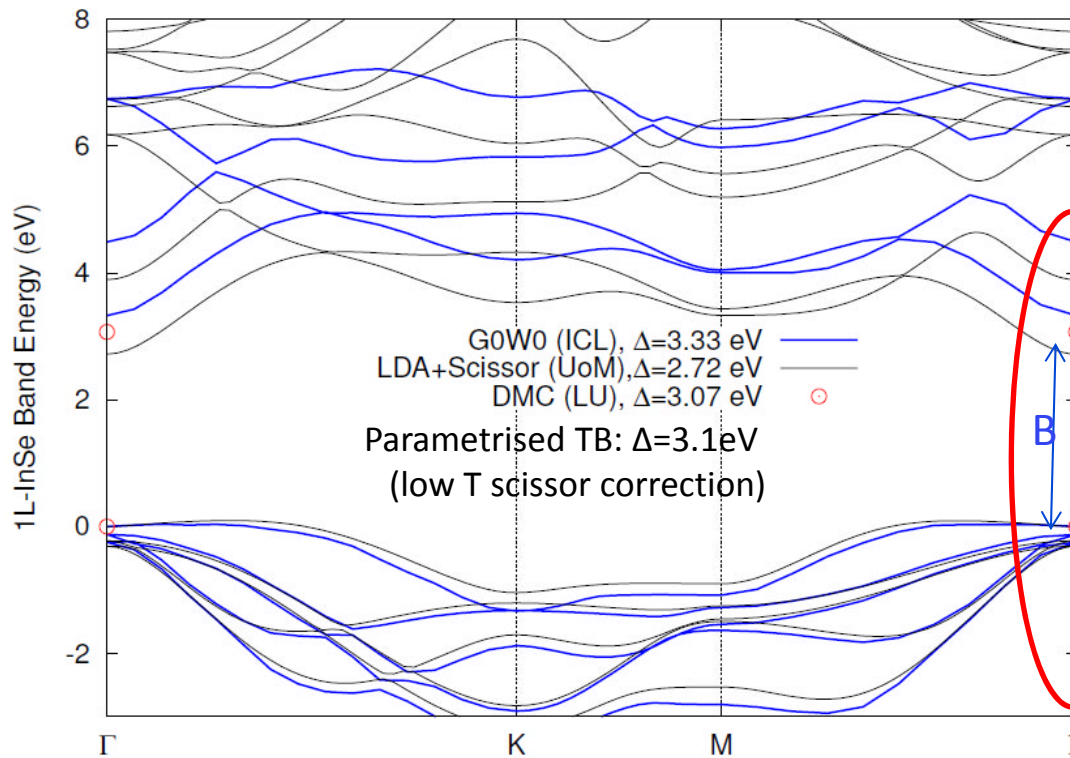


DFT-parametrised tight-binding model for $\text{In}_2\text{NSe}_2\text{N}$



$z \rightarrow -z$ mirror symmetry

VB-CB transition 'A' across the gap is active in z-polarisation;
 a higher-energy transition 'B' to a deeper double degenerate at the Γ -point valence band is active in the x-y polarisation



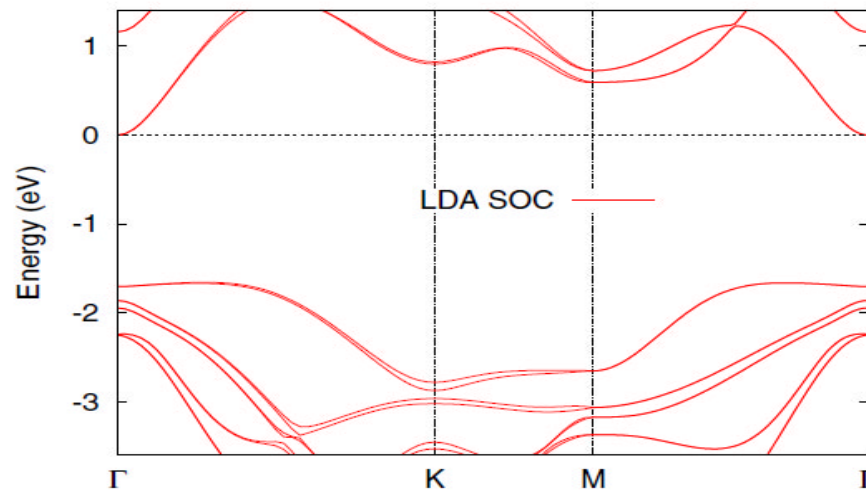
$0.115s^{\text{In}} + 0.005\bar{p}_z^{\text{In}} + 0.059s^{\text{Se}} + 0.088\bar{p}_z^{\text{Se}}$
 antisymmetric (odd)

symmetric (even)

$0.012s^{\text{In}} + 0.043p_z^{\text{In}} + 0.001s^{\text{Se}} + 0.198p_z^{\text{Se}}$

Role of spin-orbit coupling in monolayer InSe

$$\hat{H}_c = \hbar^2 k^2 / 2m_c + \gamma_c s_z k^3 \cos(3\phi)$$



$$\hat{H}_v(k, \phi) = E_0 + E_2 k^2 + E_4 k^4 + E_6 k^6 + E_8 k^8 + E'_6 k^6 \cos(6\phi) + \gamma_v s_z k^3 \cos(3\phi)$$

odd/even $z \rightarrow -z$
conduction/valence
band states allow for

$$\frac{\hbar e}{cm_e} \beta_{sf} \vec{A} \times \vec{s}$$

spin-flip
(due to atomic SO
coupling) interband
transition coupled with
in-plane xy-polarised
photons

4-band k·p theory for monolayer InSe

$$\hat{H}_c = \hbar^2 k^2 / 2m_c + \gamma_c s_z k^3 \cos(3\phi)$$

$$\hat{H} = \begin{pmatrix} H_c & & & \\ E_z d_z + \frac{\hbar e}{cm_e} \beta_{sf} \vec{A} \times \vec{s} & & & \\ \frac{\hbar e}{cm_e} \beta \vec{A} & & & \\ 0 & & & \end{pmatrix}$$

$$\hat{H}_v(k, \phi) = E_0 + E_2 k^2 + E_4 k^4 + E_6 k^6 + E_8 k^8 + E'_6 k^6 \cos(6\phi) + \gamma_v s_z k^3 \cos(3\phi)$$

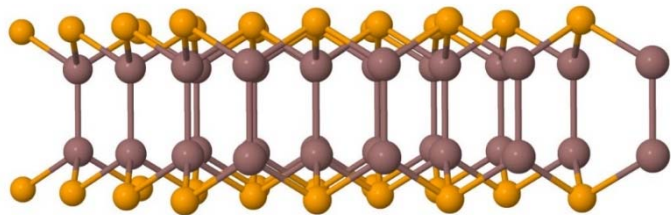
$$\begin{pmatrix} E_z d_z + \frac{\hbar e}{cm_e} \beta_{sf} \vec{A} \times \vec{s} & \frac{\hbar e}{cm_e} \beta \vec{A} & 0 \\ H_v & 0 & \frac{\hbar e}{cm_e} \tilde{\beta} \vec{A} \\ 0 & H_{v_1} & 0 \\ \frac{\hbar e}{cm_e} \tilde{\beta} \vec{A} & 0 & H_{v_2} \end{pmatrix}$$

$C_{3v} \otimes \sigma_h$	E	$2C_3$	$3\sigma_v$	σ_h	band
A_1^+	1	1	1	1	v
A_2^+	1	1	-1	1	
E^+	2	-1	0	1	v_2
A_1^-	1	1	1	-1	c
A_2^-	1	1	-1	-1	
E^-	2	-1	0	-1	v_1

$$H_{v_{1(2)}} = E_0 + \frac{k^2}{2m} + \frac{k_x^2 - k_y^2}{2m'} \sigma_x + \frac{2k_x k_y}{2m'} \sigma_y$$

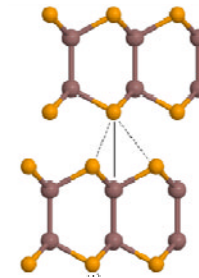
Electronic bands in $\text{In}_{2N}\text{Se}_{2N}$

Monolayer (In_2Se_2)



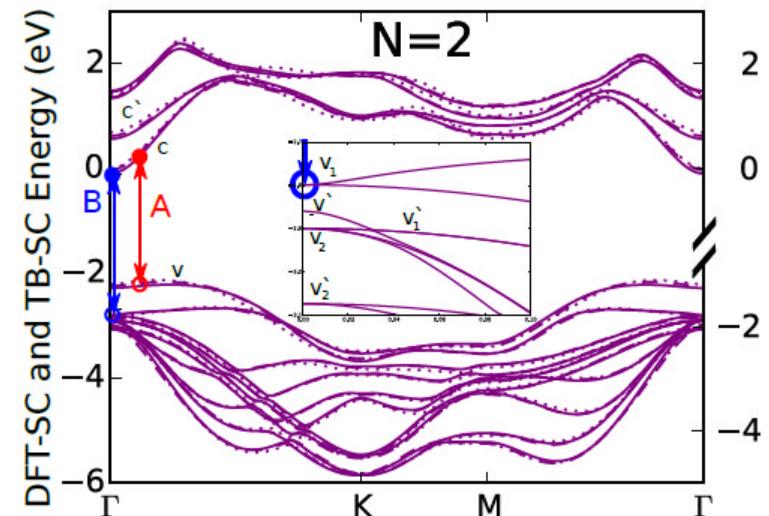
$z \rightarrow -z$
mirror
symmetry

Bilayer (In_4Se_4)

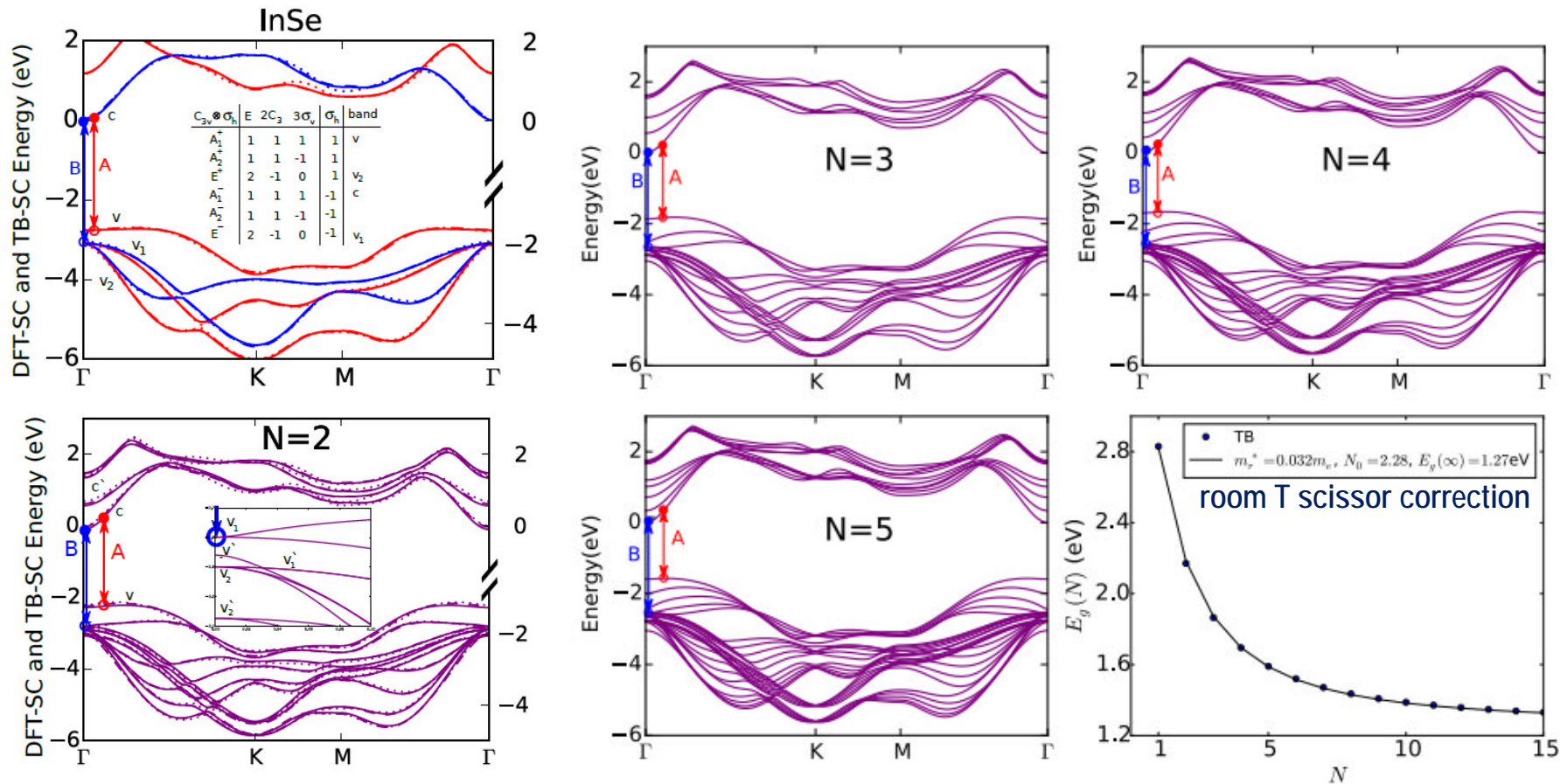


broken
mirror
symmetry

Magorrian, Zólyomi, Fal'ko - PRB 94, 245431 (2016)



Electronic bands in $\text{In}_{2N}\text{Se}_{2N}$

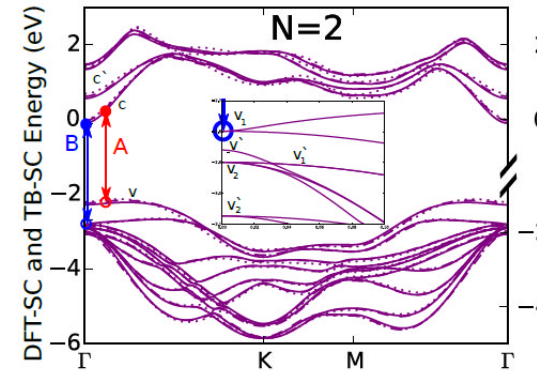
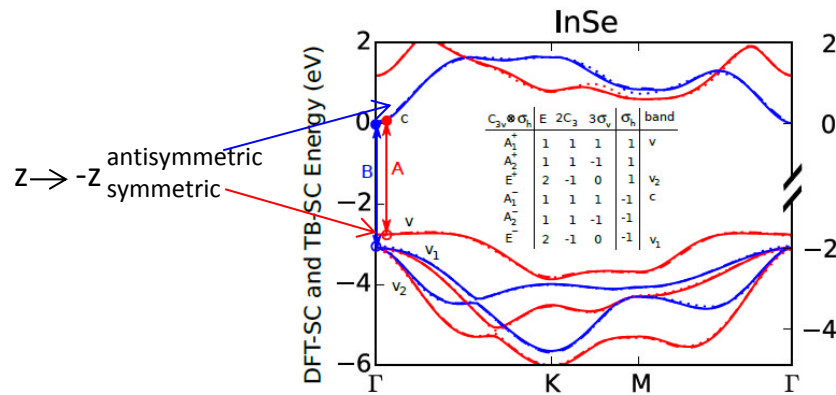


Relatively light electrons CB and wide interval of an almost flat VB edge also appear in few-layer InSe

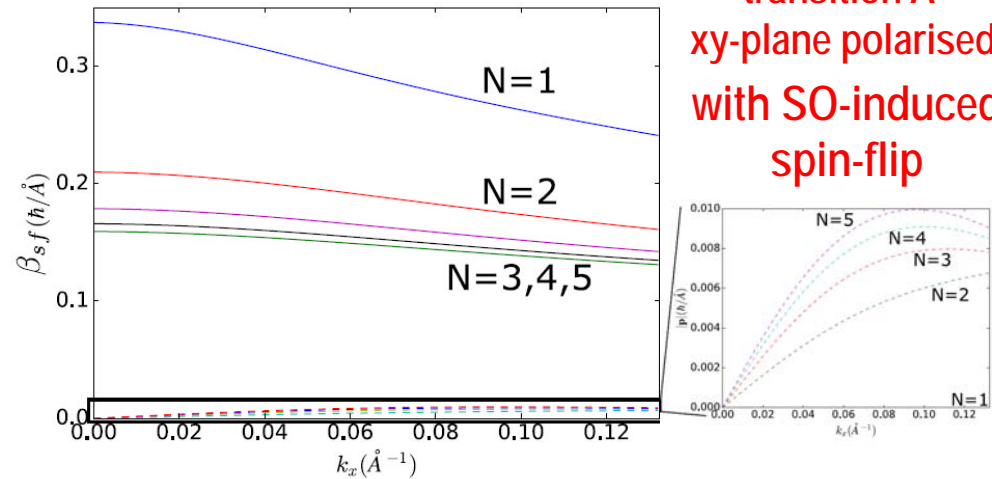
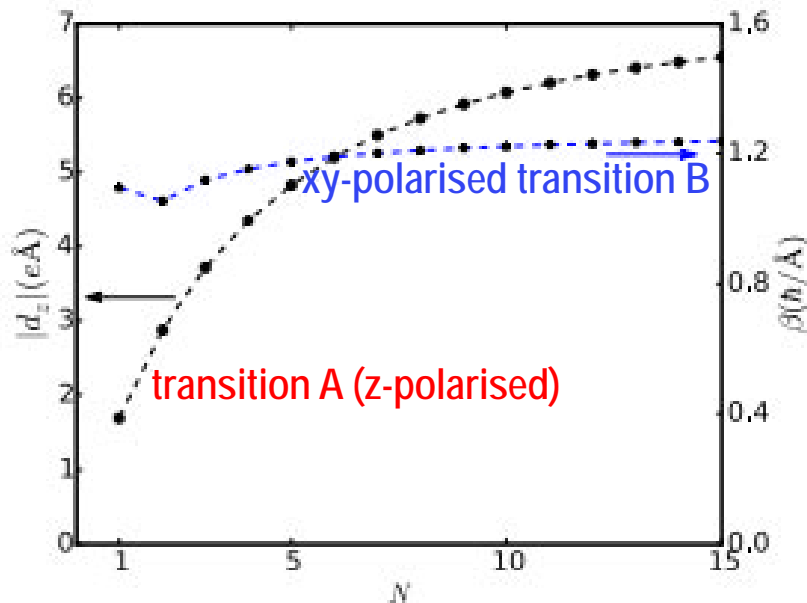
Magorrian, Zólyomi, Fal'ko - PRB 94, 245431 (2016)

large variation of the band gap as a function of number of layers.

Optical transitions in $\text{In}_{2N}\text{Se}_{2N}$

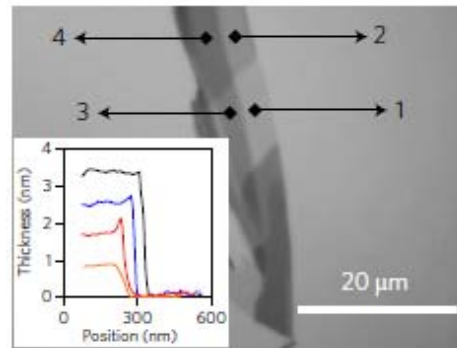
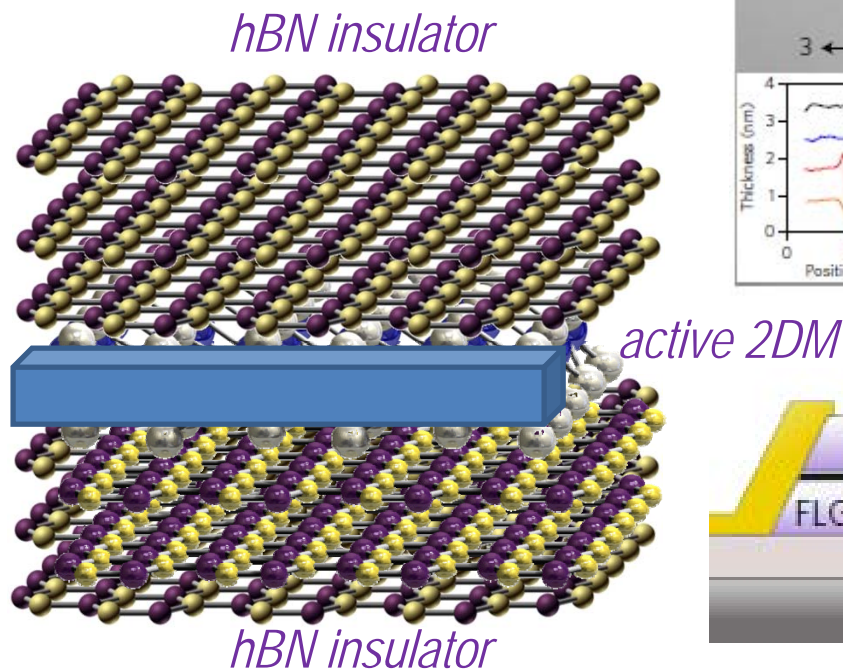


broken
 $Z \rightarrow -Z$ mirror
 symmetry



hBN-encapsulated In_2NSe_2

- hBN provides atomically flat substrate and clean encapsulation environment (low-resistance side contacts are possible)

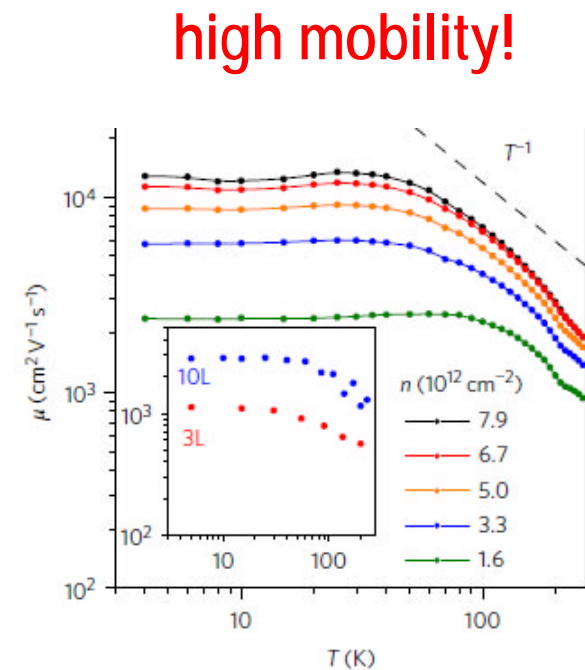
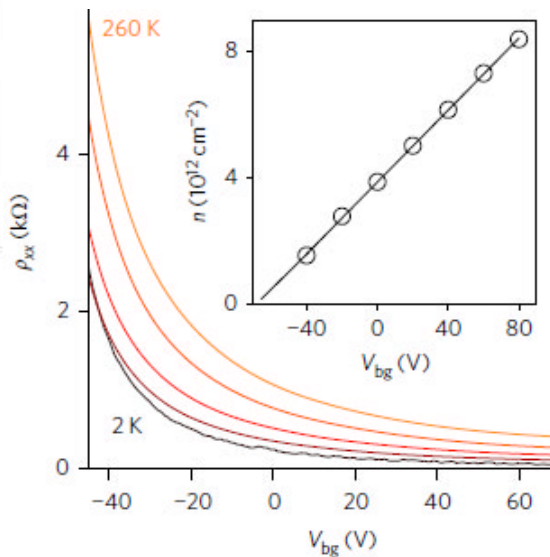
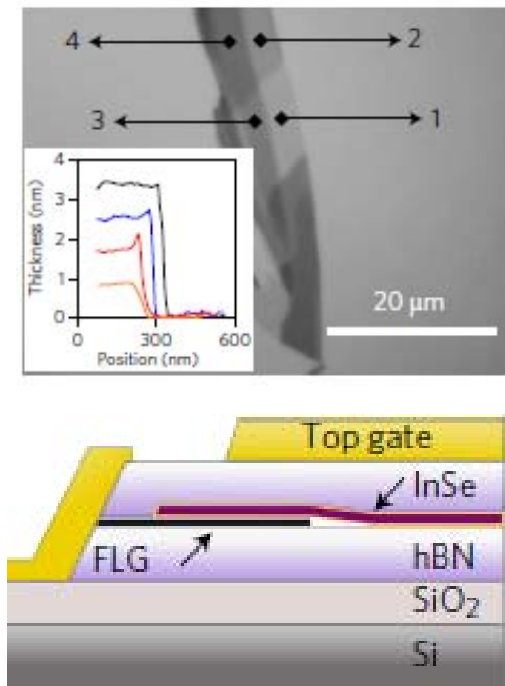


hBN is sp^2 – bonded insulator with a large band gap
it is transparent and takes high voltage drop (useful for electrostatic gating)

- hBN and 2DM are glued together by weak van der Waals attraction, hence, they preserve their lattice structure and, hence, retain their basic physical properties.

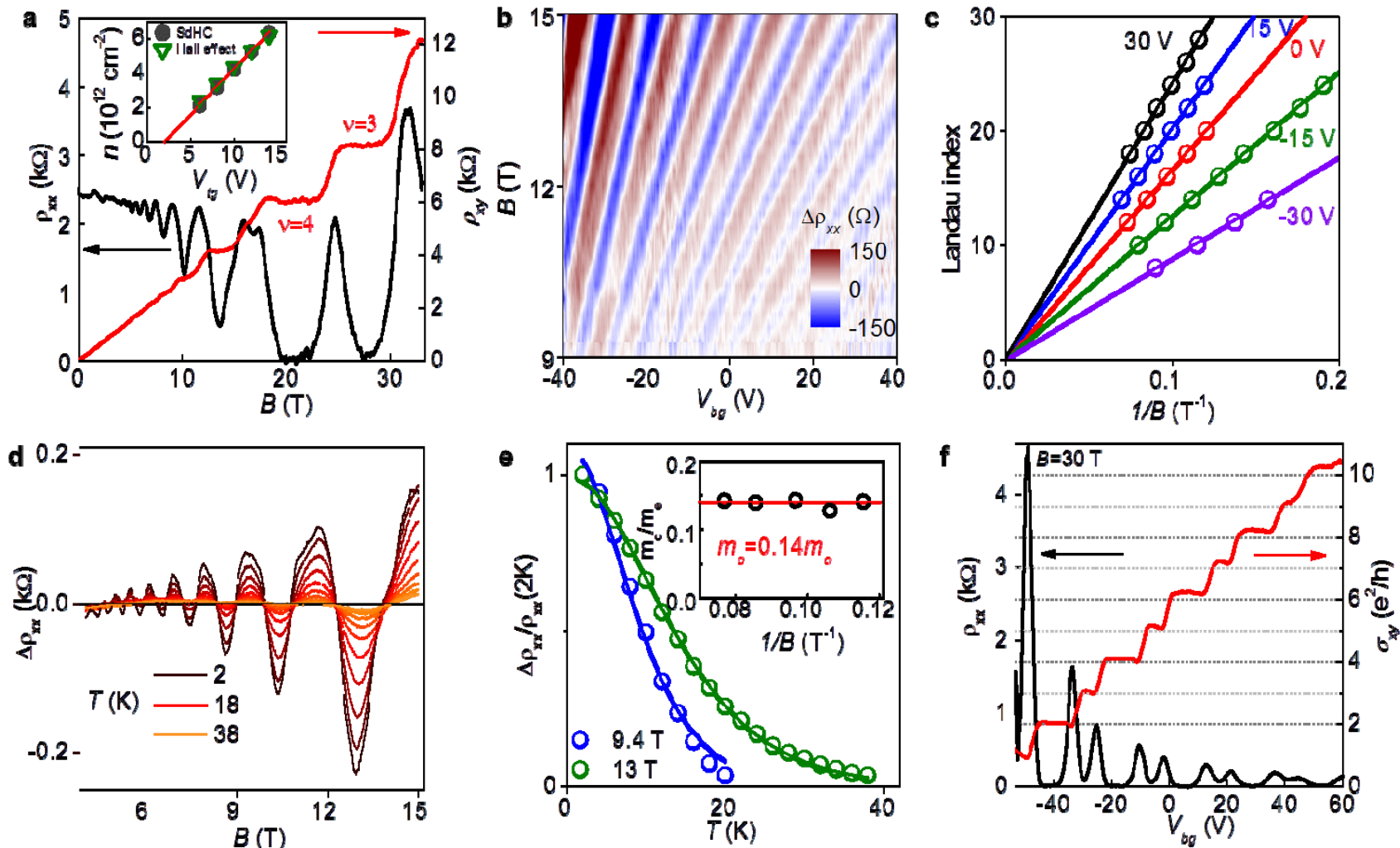
hBN-encapsulated In_2NSe_2

- hBN provides atomically flat substrate and clean encapsulation environment (low-resistance side contacts are possible)



- hBN and 2DM are glued together by weak van der Waals attraction, hence, they preserve their lattice structures and, hence, retain their basic physical properties.

hBN-encapsulated 2D crystal: QHE in 6L-InSe

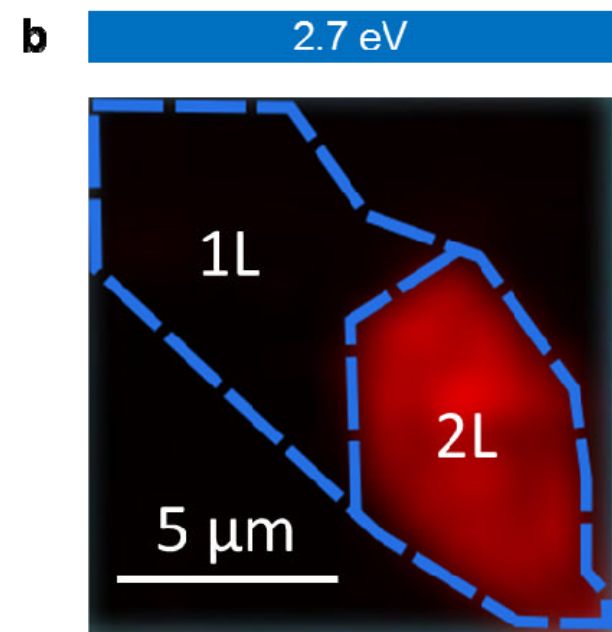
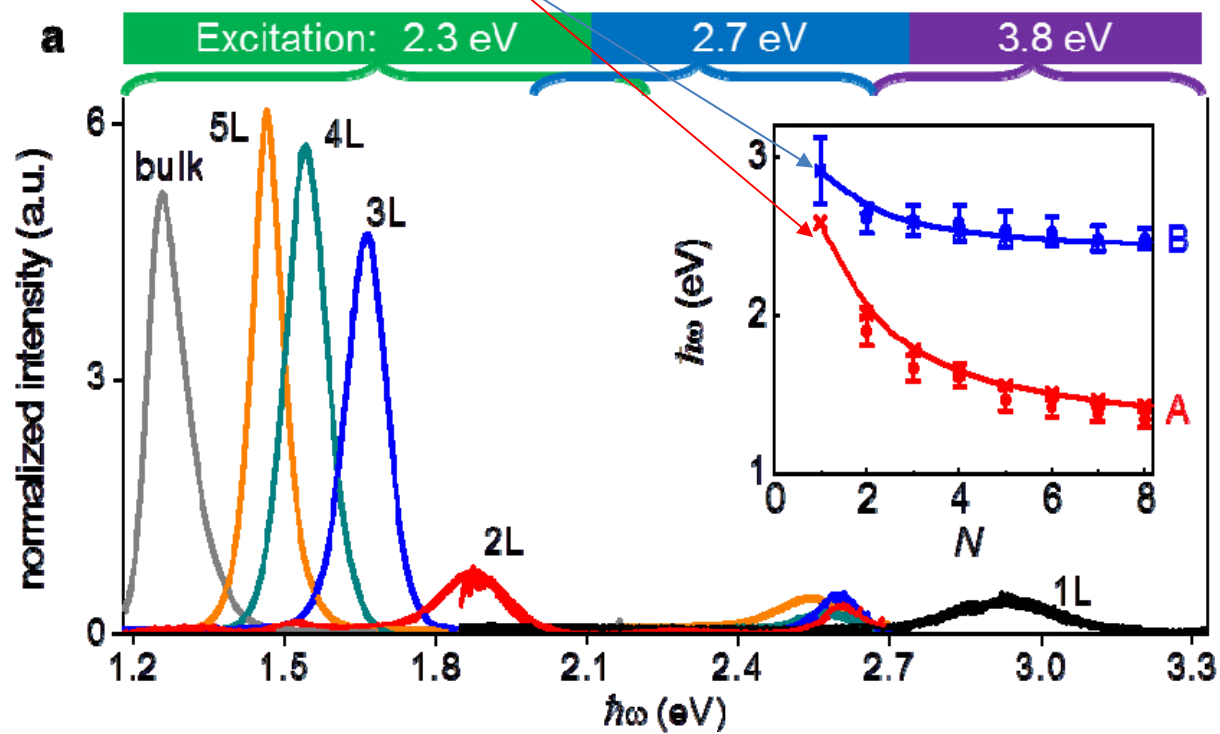


$$m_c(N=3)=0.17m_e$$

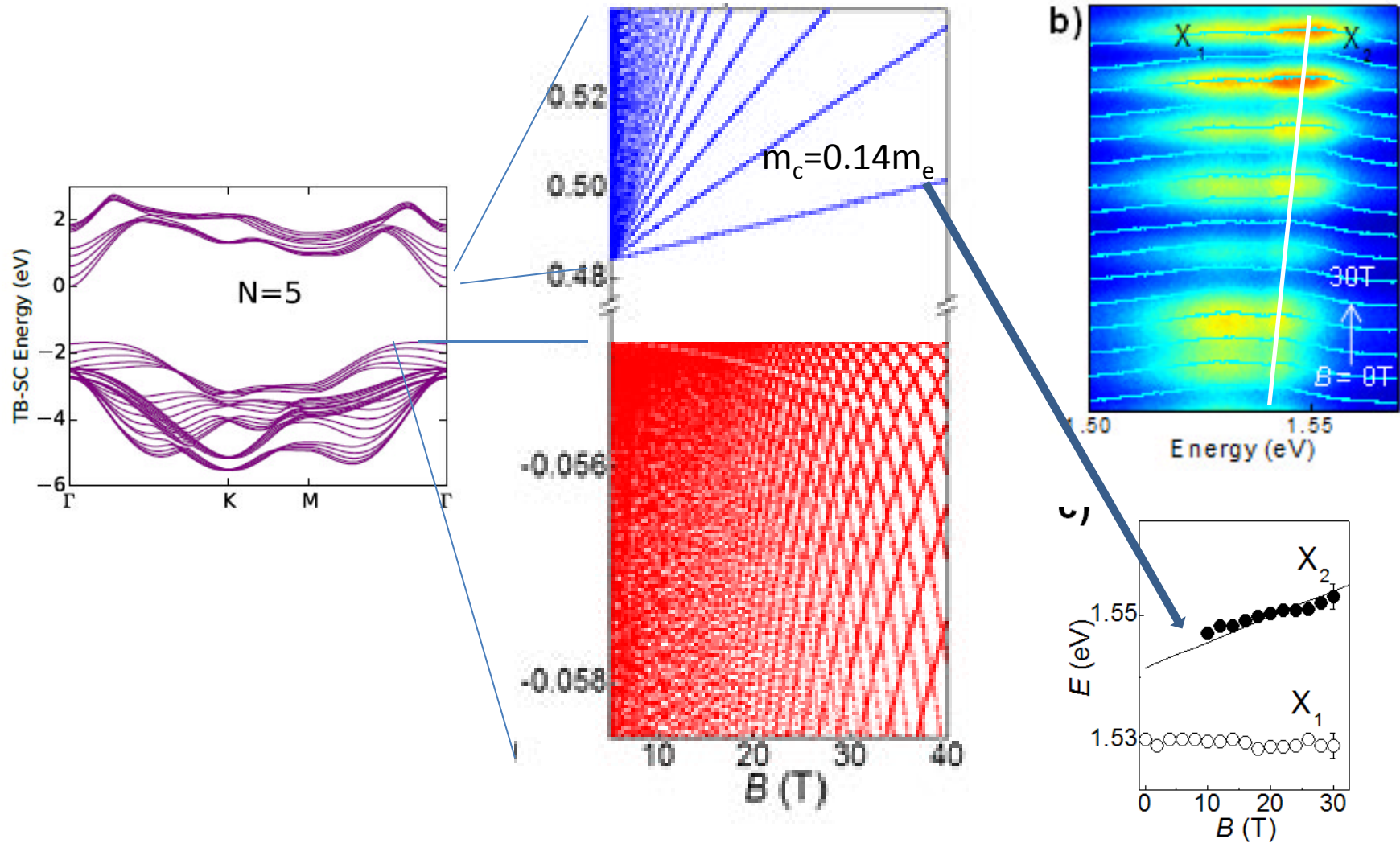
$$m_c(N=6)=0.14m_e$$

Bandurin, Tyurnina, Yu, Mishchenko, Zólyomi, Morozov, Krishna Kumar, Gorbachev, Kudrynskiy Pezzini, Kovalyuk, Zeitler, Novoselov, Patanè, Eaves, Grigorieva, Fal'ko, Geim, Cao

DFT/TB with scissor correction



Magnetoluminescence (non-encapsulated 5L)

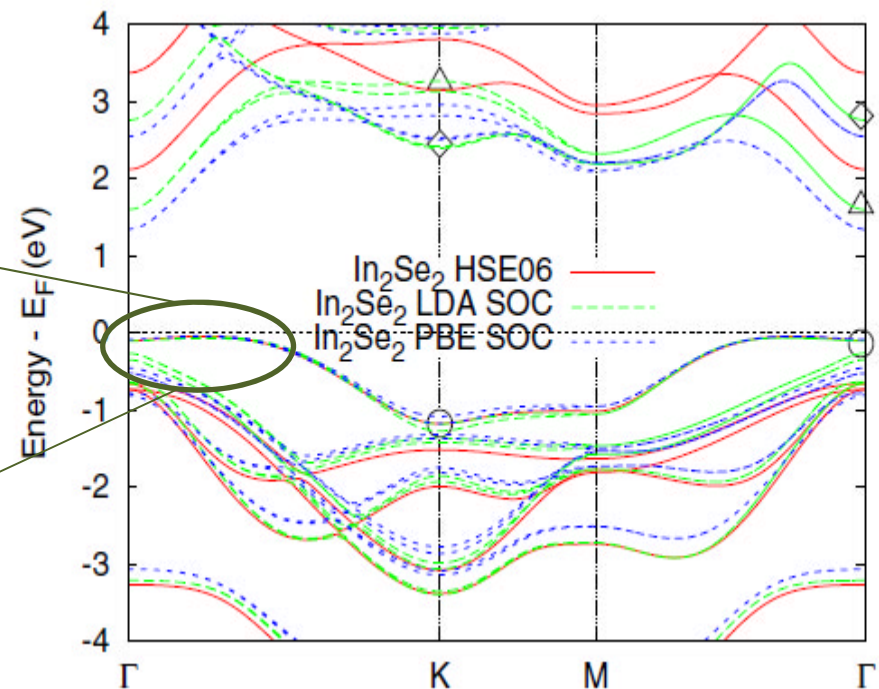


Mudd, Molas, Chen, Zolyomi, Nogajewski, Kudrynskiy, Kovalyuk, Yusa, Makarovskiy, Eaves, Potemski, Fal'ko, Patane - Scientific Reports 6, 39619 (2016)

Two-dimensional InSe

- strong band gap variation with the number of layers
- light conduction band mass:
potential for high mobility and quantum circuits
- potential for strongly correlated states
in p-doped 2D InSe

P-doped: Wigner crystal?
Mott insulator? ferromagnetic state?
Peierls instability? superconductivity?





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1824



S Magorrian (NGI)

V Zolyomi (NGI)

X Chen (NGI)

N Drummond (Lancaster)

J Lischner (Imperial)

A Tyurnina (NGI)

D Bandurin (NGI)

R Gorbachev (NGI)

K Novoselov (NGI)

A Geim (NGI)

A Patane (Nottingham)

L Eaves (Nottingham)

Z Kovalyuk (IPMS-UAS)

U Zeitler (Nijmegen)

M Potemski (Grenoble)