

Strain and correlations on flat bands of TBLG in STM Measurements



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raphene Conline



LDOS calculations

(tight binding)

0.0

Energy (meV)

-50

SC

0.0

50

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Introduction

In twisted bilayers of graphene (TBLG), the two layer's graphene periodicity interfere in a moiré pattern that is characteristic of their precise stacking arrangement [1, 2]. For low angles (~1.1°), the charge carriers posses no kinetic energy, thus enhancing correlations between the electrons and yielding the formation of novel correlated states [3,4]. Here we survey recently published STM (Scanning Tunneling) Microscopy) measurements Ref. [5-13] showing that for such low angles the physics are controlled first by heterostrain *i.e.* the relative strain between the two layers, small angle variations acting on top of that. We show that even small details in heterostrain can have dramatic effect on the flat bands, and enable STM experiments for highly doped samples. When the energy of electrons lie in the flat bands however, the system enters a new regime that is not captured by Hartree Fock calculations, though they include strain and electronic interactions.





28.8

7.2

0.0

0.0



When the angle of application of uniaxial

Meta-analysis of recent STM experiments

Experimental vHs spacings as a function of uniaxial strain, compared to theoretical predictions

strains varies, 3 different behaviours of the LDOS can be distinguished :

 $-\theta_{s} \sim 8.5^{\circ}$: a zero energy peak of high density of states appears;

 $-\theta_{s} \sim 30^{\circ}$: two simple vHs appear ;

- $\theta_{s} \sim 51^{\circ}$: the two vHs are splitted appart.

21.6 Heterostrain alters the flat bands of low 14.4 angle TBLG because it occurs in a Van der Waals stacking that has no interlayer chemical bonds (unlike for cuprates for ex), creating a relative displacement of 28.8 🗐 the atomic lattice that accumulates over large superlattices (> 100 nm²), and 21.6 making its effect non negligible even for 14.4 0 small values. 7.2 9

vHs spacing increases linearly when 28.8 uniaxial strain is increased $\sim 50 \text{meV}/\%$ 21.6



Experimental vHs spacings as a function of θ_{int} *, and tight binding* predictions for several values of strain



show a linear trend well described by tight binding (ΔE_{TB}) and Hartree Fock calculations ($\Delta E^{\epsilon=\infty}_{HF}$) that include strain.

For high dopings, experiments (ΔE^0_{ex}

Close to CNP, Hartree Fock calculations $(\Delta E^{\epsilon=5}_{HF})$ including strain and coulomb interactions underestimate the vHs spacing. Experiments (ΔE^*_{exp}) reveal increasing spacing with strain and large sample to sample variation.

Experiments show no correlation with small interlayer angle variation.

TB calculations including heterostrain show low influence of interlayer angle for a given strain. On the contrary, a small variation of strain at a given angle has quite strong effect on the vHs spacing.



outlook

Identifying heterostrain as a critical parameter for this highly tunable system explains the ubiquitous sample to sample variability. Its implementation in the existing many models could be prolific. Yet, high variability remains when electron-electron interactions are strongest, pointing the need to renew the structural parametrization of flat band in that regime, when we know that combining heterostrain and twist angle is too simplistic. Sample environment is a good candidate because it can engineer the electrostatic environment that affects coulomb interactions and can also distort the flat bands through substate quality. Moreover, in STM experiments the tip can induce additional strain, locally modifying the relative stacking as well as the spacing between the two layers. This calls for a more systematic study of distorted flat bands, both in a local manner by the tip and by the substrate morphology and dielectric constant.

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