Out-of-plane heat transfer in Van der Waals stacks





The Institute of Photonic Sciences

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RWITHAACHEN





Thermal management





Ghosh et al, APL 92, 151911 (2008)

Suspended graphene has a very high in-plane thermal conductivity: >1000 W/mK

Supported graphene: ideal substrate = heat sink!

Light-emitting diodes





Conventional GaN LED

Embedded graphene layer GaN LED

Han et al, Nature Comm. 4, 1452 (2013)

Field-effect transistors



 Temperature (°C)
 95

 78
 60

 42
 25

113

AlGaN/GaN FET with graphene/graphite

Yan et al, Nature Comm. 3, 827 (2012)

Fast photodetectors

>50 Gbit/s



Schall et al, ACS Photonics (2014)

Hot graphene electrons



Photo-thermoelectric (PTE) detector

Song et al. *PRL* (2011) Gabor et al. *Science* (2011) Koppens et al. *Nature Nano* (2014)

Hot graphene electrons



Photo-thermoelectric (PTE) detector





$$V_{\rm PTE} = (S_2 - S_1)\Delta T_{\rm e}$$

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Understand thermal transport

Photodetectors: electron cooling?



Thermal management: ideal substrate = heat sink!



Understand thermal transport



Efficient out-of-plane heat transfer:

Hot graphene electrons to hBN hyperbolic phonon polaritons



Hot electrons:

broadened Fermi-Dirac distribution

Electron heat \neq **phonon heat**



Planck radiation: coupling to *far-field* light in vacuum

$$k < \omega/c$$



Planck radiation: coupling to *far-field* light in vacuum

$$k < \omega/c$$

Very inefficient!

=> Cooling to graphene phonons dominates (Governed by deformation potential) Bistritzer and MacDonald, PRL (2009) Song et al. PRL (2012) Graham et al. Nature Phys. (2013) Betz et al. Nature Phys. (2013)



Planck radiation: coupling to *far-field* light in vacuum

$$k < \omega/c$$

Very inefficient!

=> What about *near-field* radiation to encapsulant?







Hyperbolic modes: $sgn(\epsilon_{\perp}) \neq sgn(\epsilon_{\parallel})$



Yoxall et al, Nature Photon. (2015)



Yoxall et al, Nature Photon. (2015)





 k_{v}

hBN: Hyperbolic phonons *Isotropic medium:* $\epsilon_{\perp} = \epsilon_{\parallel}$ $k = \omega/c$ k_{\star}

<u>Hyperbolic mode (Type II):</u> $\operatorname{sgn}(\epsilon_{\perp}) \neq \operatorname{sgn}(\epsilon_{\parallel})$

 $k\gg \omega/c$



Near-field, hyperbolic cooling

<u>Graphene in vacuum</u>

- Light cone with restricted *k*-vectors
- Low thermal energy density
- Blackbody Planck radiation (inefficient)



Near-field, hyperbolic cooling

<u>Graphene in vacuum</u>

- Light cone with restricted *k*-vectors
- Low thermal energy density
- Blackbody Planck radiation (inefficient)

hBN-encapsulated graphene

- Hyperbolic modes with near-infinite range of k-vectors
- High thermal energy density
- Super-Planckian radiation (very efficient)

More efficient than graphene phonon cooling?



Experiment



Photo-thermoelectric (PTE) detector

Experiment



Photo-thermoelectric (PTE) detector

Experiment



Access to temperature dynamics!

 $\Delta V_{\rm PTE}(\Delta t) \sim \Delta T_{\rm e}(\Delta t)$

Ultrafast heating



Cooling: varying carrier density



Cooling: varying carrier density



Disorder-assisted scattering with acoustic graphene phonons:

- Deformation potential
- Disorder density

Graham et al. *Nature Phys* (2013) Song et al. *PRL* (2012)



• Super-collision cooling is faster for *lower* carrier density

Graham et al. Nature Phys (2013)



- Super-collision cooling is faster for *lower* carrier density
- Super-collision cooling requires an unrealistic deformation potential of 65 eV



- Super-collision cooling is faster for *lower* carrier density
- Super-collision cooling requires an unrealistic deformation potential of 65 eV
- Our device has a deformation potential <35 eV





Tielrooij et al, Arxiv:1702.03766 (2017)

Cooling: varying lattice temperature



Tielrooij et al, Arxiv:1702.03766 (2017)

Normal collision cooling?

Normal collision cooling (in overheating) is not dependent on lattice temperature

Bistritzer and MacDonald PRL 102, 206410 (2009)

Normal collision cooling?

- Normal collision cooling (in overheating) is not dependent on lattice temperature
- Normal collision cooling gives non-exponential cooling with a timescale >30 ps at RT

Bistritzer and MacDonald PRL 102, 206410 (2009)





In-plane cooling mechanisms are not consistent with data

Tielrooij et al, *Arxiv*:1702.03766 (2017)

Out-of-plane transport!



In-plane cooling mechanisms are not consistent with data

Tielrooij et al, Arxiv:1702.03766 (2017)

Compare with hyperbolic cooling



Compare with hyperbolic cooling



Cooling to hyperbolic hBN phonon polaritons reproduces the trends

Principi et al, *PRL 118, 126804* (2017) Tielrooij et al, *Arxiv*:1702.03766 (2017)

$$\mathcal{Q} = \iiint_{-\infty}^{\infty} \frac{d\omega dk_x dk_y}{(2\pi)^3} [n_B(\omega, T_{\rm e}) - n_B(\omega, T_{\rm L})] M(\omega, k)$$

Energy transfer rate



Marco Polini



Mark Lundeberg



Alessandro Principi Principi et al, *PRL 118, 126804* (2017)

$$\mathcal{Q} = \iiint_{-\infty}^{\infty} \frac{d\omega dk_x dk_y}{(2\pi)^3} [n_B(\omega, T_{\rm e}) - n_B(\omega, T_{\rm L})] \mathcal{M}(\omega, k)$$

Energy transfer rate

$$M(\omega, k) = 4 \frac{\mathcal{R}\{Y(\omega, k)\} \mathcal{R}\{\sigma(\omega, k)\}}{|Y(\omega, k) + \sigma(\omega, k)|^2}$$

Impedance matching function between graphene electrons and hBN phonon polaritons



Marco Polini



Mark Lundeberg



Alessandro Principi Principi et al, *PRL 118, 126804* (2017)



$$\mathcal{Q} = \iiint_{-\infty}^{\infty} \frac{d\omega dk_x dk_y}{(2\pi)^3} [n_B(\omega, T_{\rm e}) - n_B(\omega, T_{\rm L})] M(\omega, k)$$

Energy transfer rate

$$\Gamma_{\rm cool} = \frac{\partial Q}{\partial T_{\rm e}} \bigg|_{T_{\rm e} = T_{\rm L}}$$

Interfacial heat conductivity

$$\tau_{\rm calc}^* = C_{\rm n} / \Gamma_{\rm cool}$$

Near-equilibrium cooling time



Marco Polini



Mark Lundeberg



Alessandro Principi Principi et al, *PRL 118, 126804* (2017)

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- Origin of the low-density peak?
- Effect of laser power?



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Tielrooij et al, Arxiv:1702.03766 (2017)



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Tielrooij et al, Arxiv:1702.03766 (2017)



- Origin of the low-density peak?
- Effect of laser power?



$$\tau_{\rm calc}(T_{\rm e},T_{\rm L}) = C_{\rm n} \frac{T_{\rm e} - T_{\rm L}}{\mathcal{Q}}$$

Cooling time in overheating regime

Principi et al, PRL 118, 126804 (2017)

Longer cooling time with increasing laser power
 Consistent with hyperbolic phonon cooling

Summary

<u>Graphene encapsulated by hBN:</u> Cooling of hot graphene carriers through **out-of-plane heat transfer** to hBN hyperbolic phonons



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Experimental agreement for varying carrier density, lattice temperature and laser power



Tielrooij et al, Arxiv:1702.03766 (2017)

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