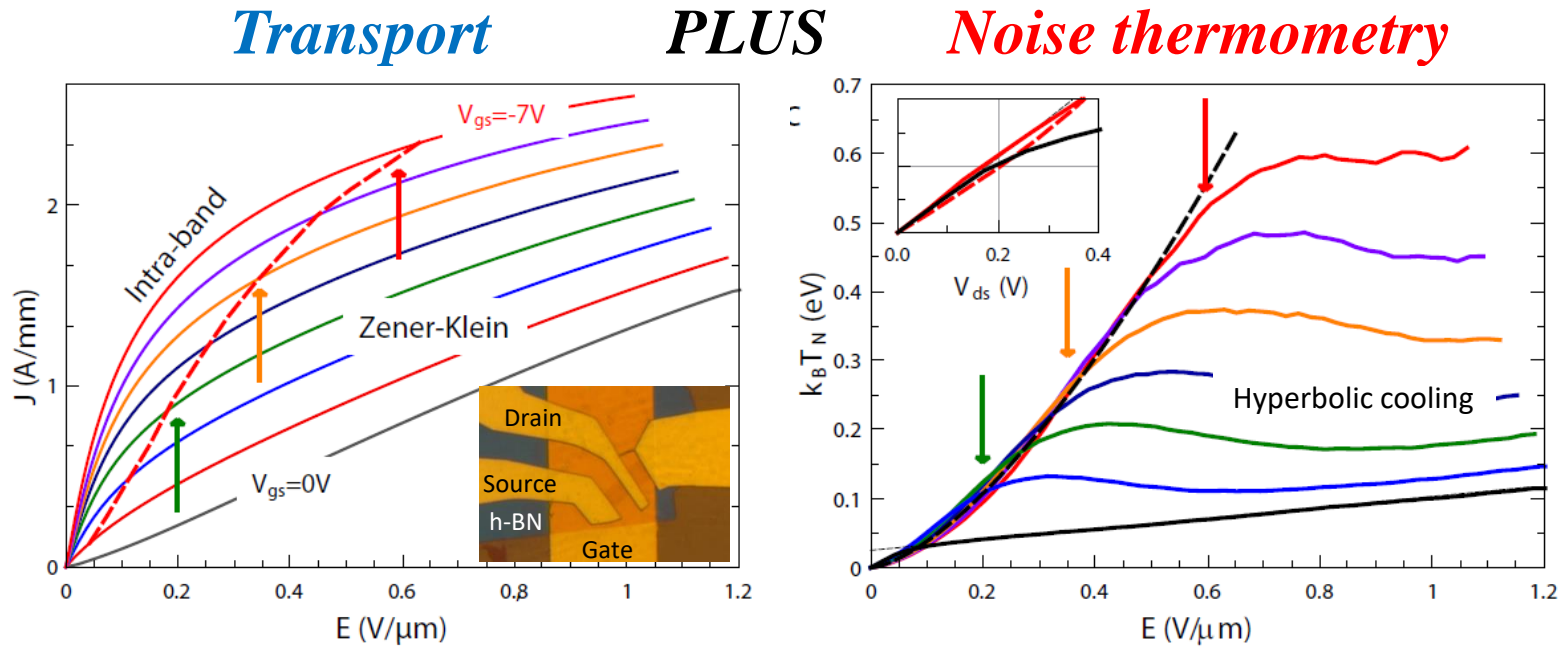
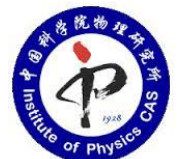




Hyperbolic cooling of graphene Zener-Klein transistors



W. Yang, S. Berthou, X. Lu, Q. Wilmart, A. Denis, M. Rosticher, T. Taniguchi, K. Watanabe, G. Fève, J.M. Berroir, G. Zhang, C. Voisin, E. Baudin, and B. Plaçais



*Noise thermometry brings new information
on scattering and relaxation of graphene carriers.*

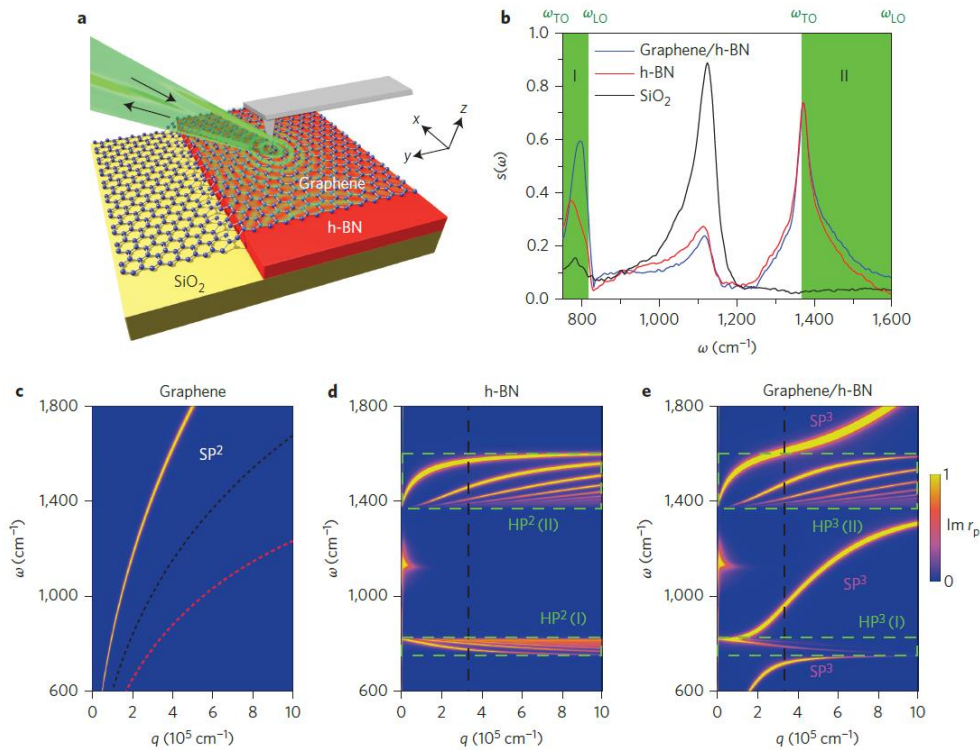
Current saturation regime is investigated here

Introduction

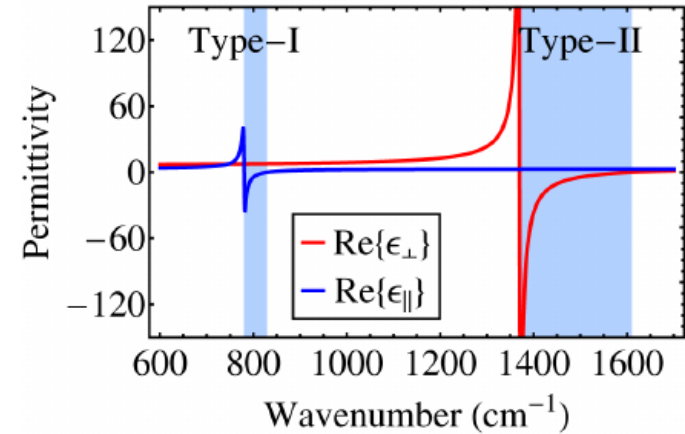
Hyperbolic Phonon Polaritons of uniaxial hBN

Graphene on hexagonal boron nitride as a tunable hyperbolic metamaterial

S. Dai¹, Q. Ma², M. K. Liu^{1,3}, T. Andersen², Z. Fei¹, M. D. Goldflam¹, M. Wagner¹, K. Watanabe⁴, T. Taniguchi⁴, M. Thiemens⁵, F. Keilmann⁶, G. C. A. M. Janssen⁷, S-E. Zhu⁷, P. Jarillo-Herrero², M. M. Fogler¹ and D. N. Basov^{1*}



Propagating HPPs

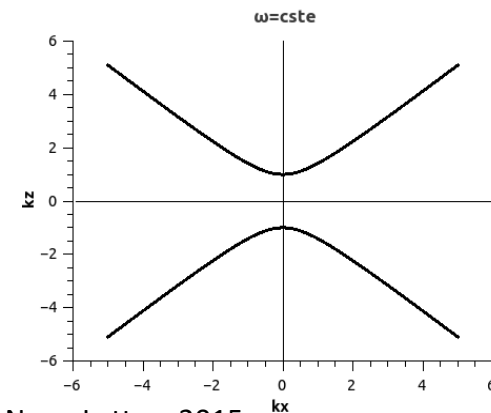


$$\epsilon_m = \epsilon_{\infty,m} + \epsilon_{\infty,m} \times \frac{(\omega_{LO,m})^2 - (\omega_{TO,m})^2}{(\omega_{TO,m})^2 - \omega^2 - i\omega\Gamma_m}$$

$$\frac{kx^2}{\epsilon_{\perp}} + \frac{kz^2}{\epsilon_{\parallel}} = \frac{\omega^2}{c^2}$$

$$(\epsilon_{\perp} < 0, \epsilon_{\parallel} > 0)$$

$$e^{i(k_x x - \Omega t)} \times e^{i(k_z z)}$$



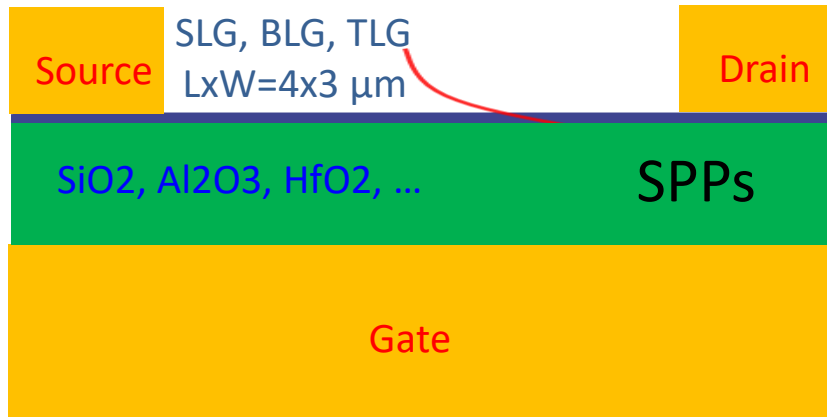
Dai et al. Nat. Nano. 2015

Caldwell et al. Nat Comm. 2014 ; Brar et al., Nano Letters 2014 ;

Kumar et al., Nano Letters 2015

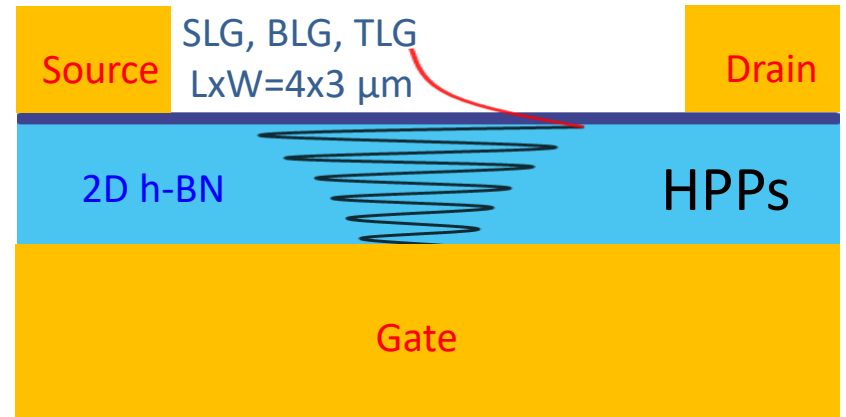
Near field coupling of graphene hot electrons with substrate phonons

Graphene on 3D oxide



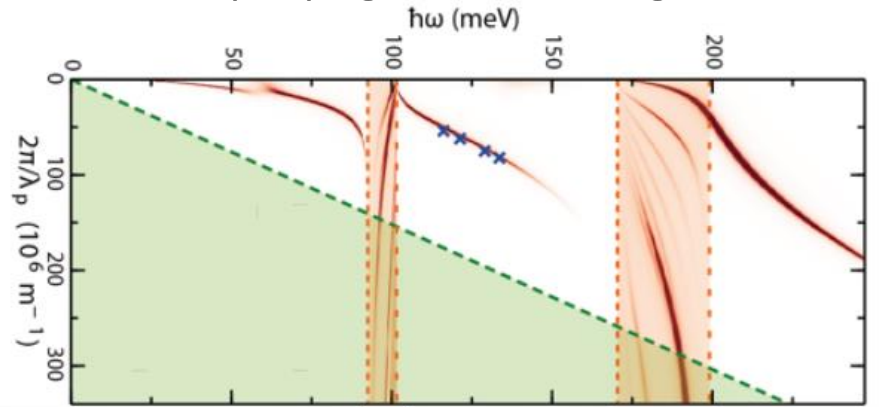
heat diffusion to the gate

Graphene on 2D h-BN



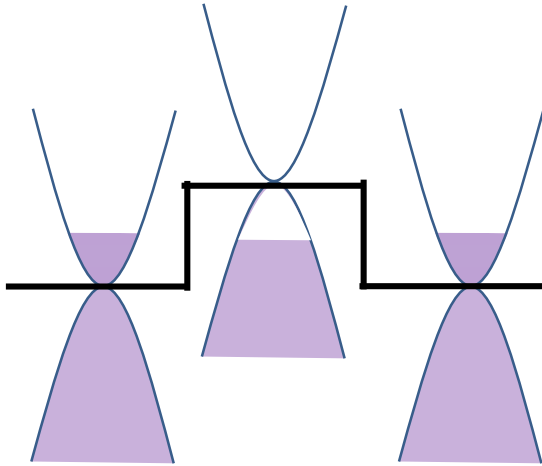
heat propagation to the gate

Graphene current fluctuations emit HPP radiations deep into hBN bulk

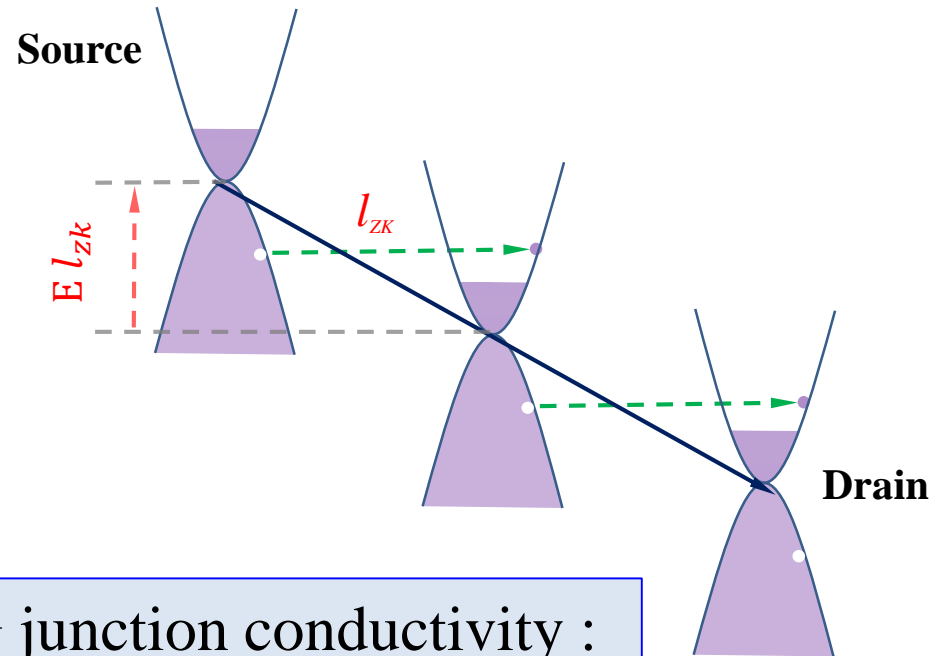


Klien and Zener Tunneling

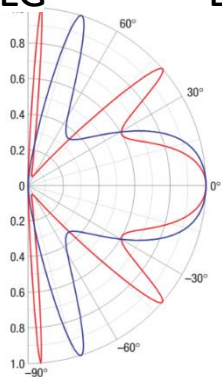
Klein Tunneling across n-p-n barriers



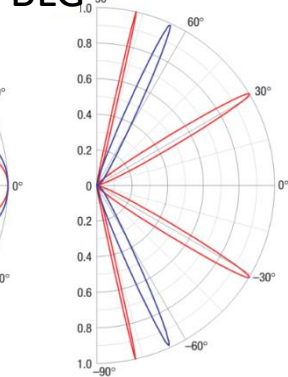
Electric field induced Zener tunneling zero bandgap semiconductor



An_c SLG



BLG



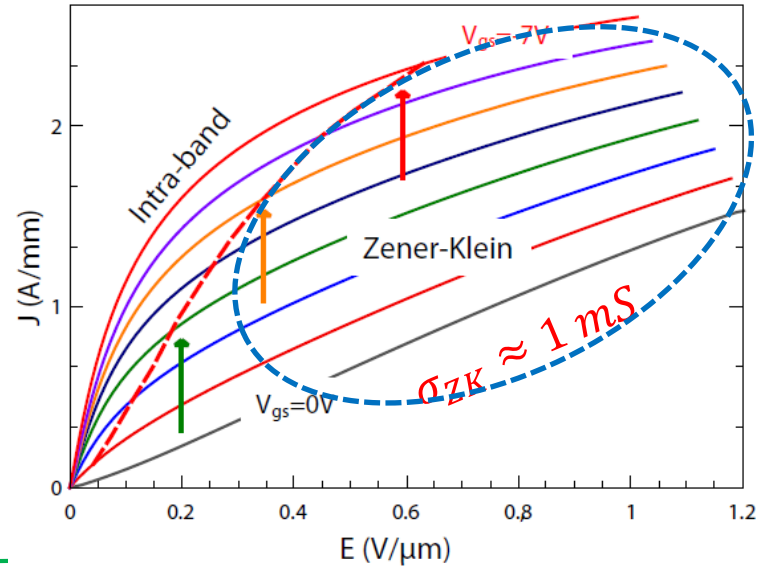
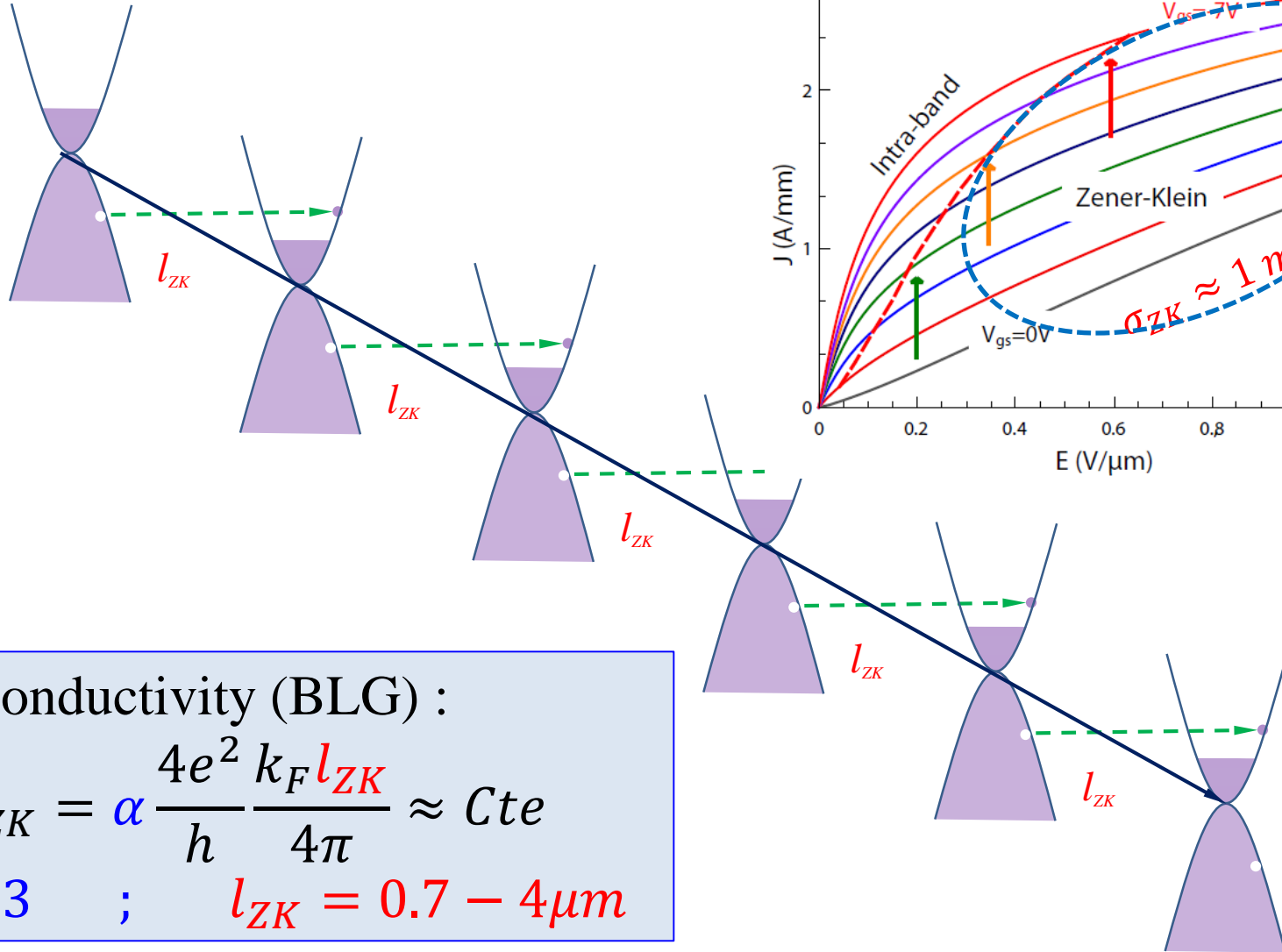
junction

BLG junction conductivity :

$$\sigma_{ZK} = \frac{4e^2}{h} \frac{k_F l_{ZK}}{4\pi}$$

Zener-Klein Tunneling (ZKT)

Source

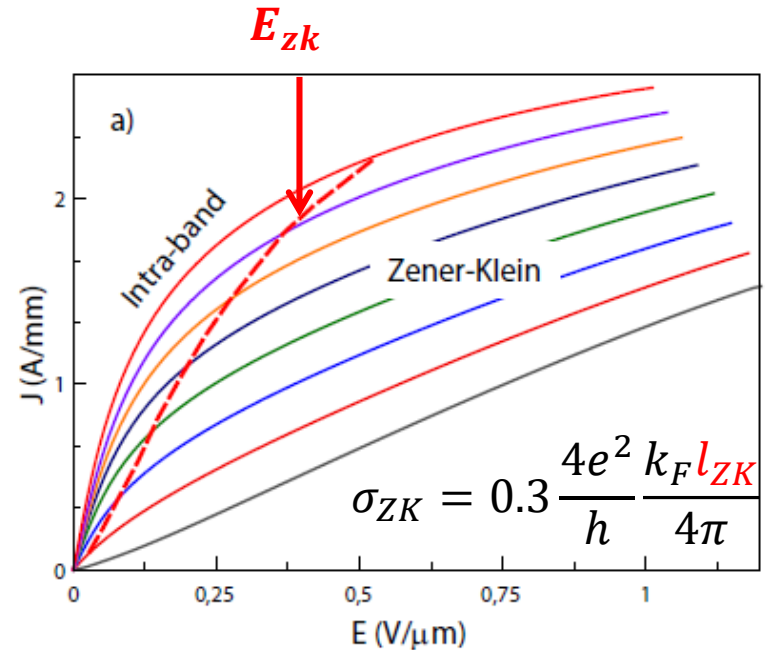
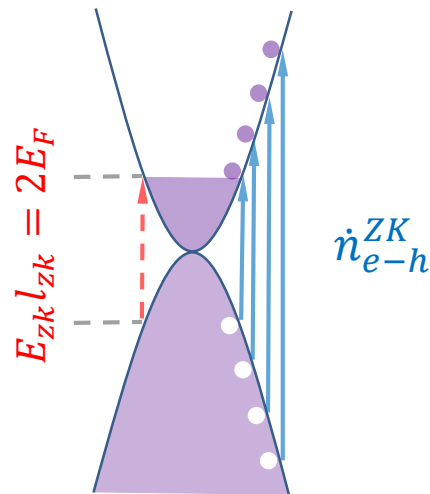


ZKT conductivity (BLG) :

$$\sigma_{ZK} = \alpha \frac{4e^2 k_F l_{ZK}}{h 4\pi} \approx Cte$$

$$\alpha \sim 0.3 \quad ; \quad l_{ZK} = 0.7 - 4 \mu m$$

Drain

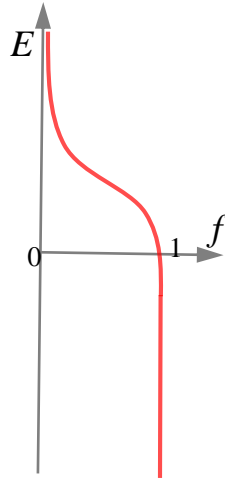
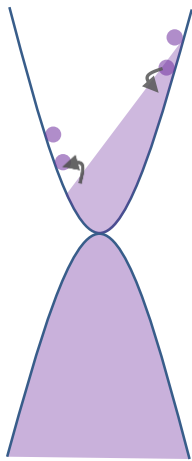
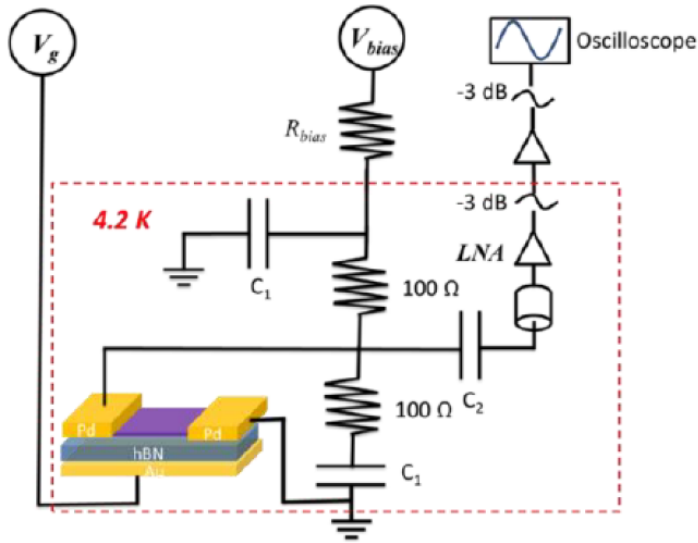


Pauli blocking of ZKT

$$E > E_{zk} = \frac{2E_F}{el_{zk}} \sim k_F^3$$

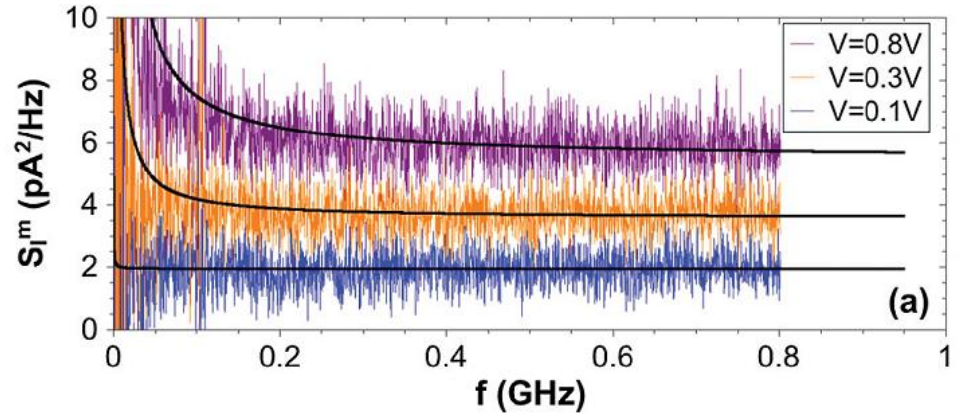
e-h creation by ZKT

$$\dot{n}_{e-h}^{ZK} = \frac{e k_F}{\pi^2 \hbar} (E - E_{ZK})$$



Graphene 2017, opto-electronics

High frequency to overcome 1/f noise



Betz et al. / Phys. Rev. Lett. 109 (2012) 056805

Betz et al. / Nat. Phys. 9 (2013) 109

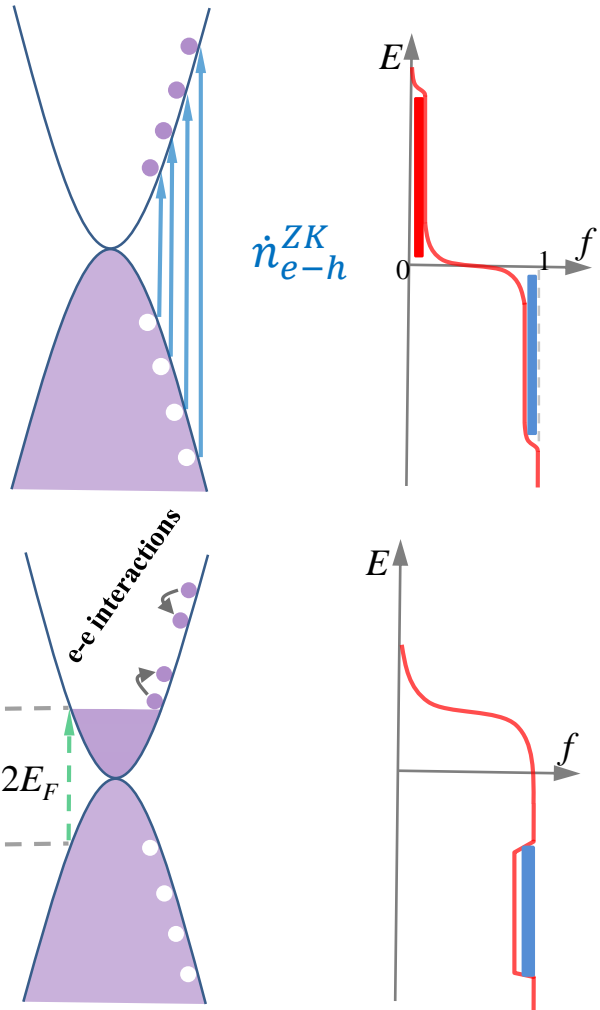
Brunel et al. / J. Phys. : Condens. Matter 27 (2015) 164208

Laitinen et al. / Phys. Rev. B. 91 (2015) 121414(R)

Thermal current noise $S_I = 4 G k_B T_N$

Fast e-e thermalisation (20 fs)

$$k_B T_N = \int_{-\infty}^{\infty} f(1-f) dE = k_B T_e$$



Out-of-equilibrium e-h population

$$k_B T_N = \int_{-\infty}^{\infty} f(1-f) dE \approx \frac{n_e + n_h}{DOS}$$

$$n_e = \int_0^{\infty} DOS \times f dE ;$$

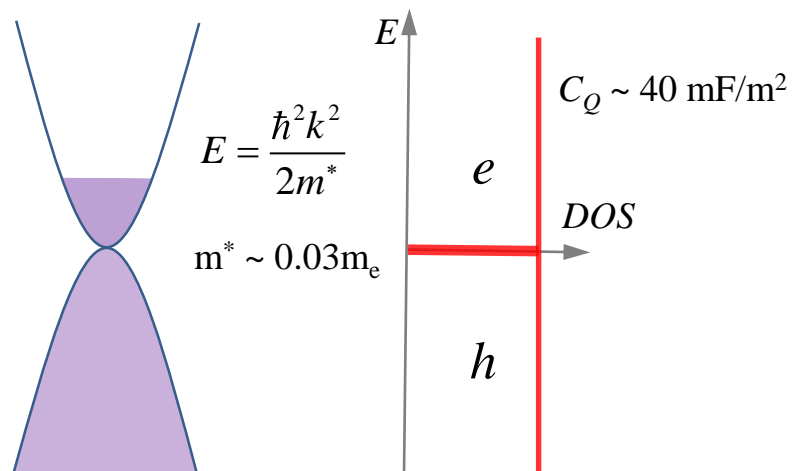
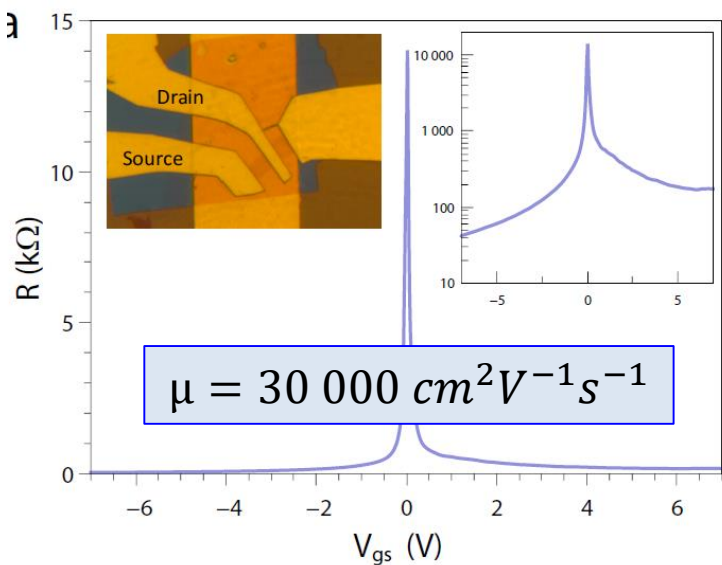
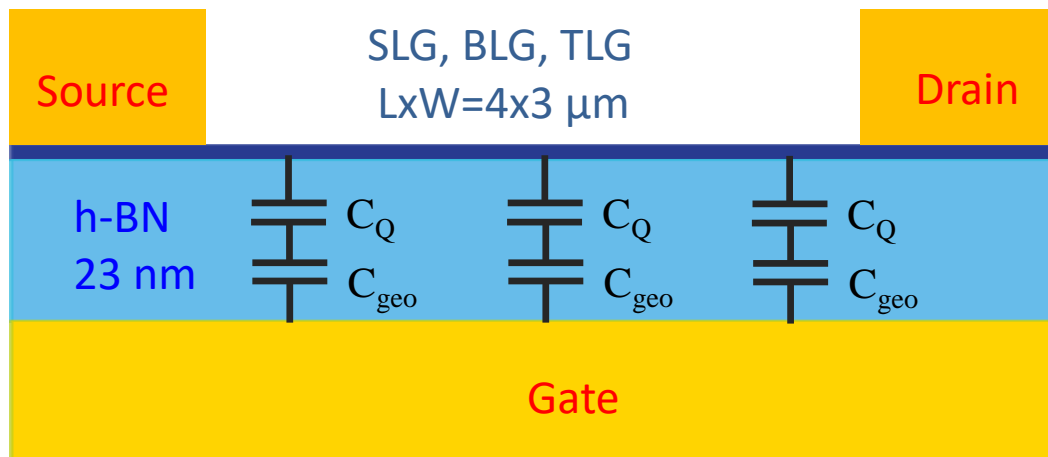
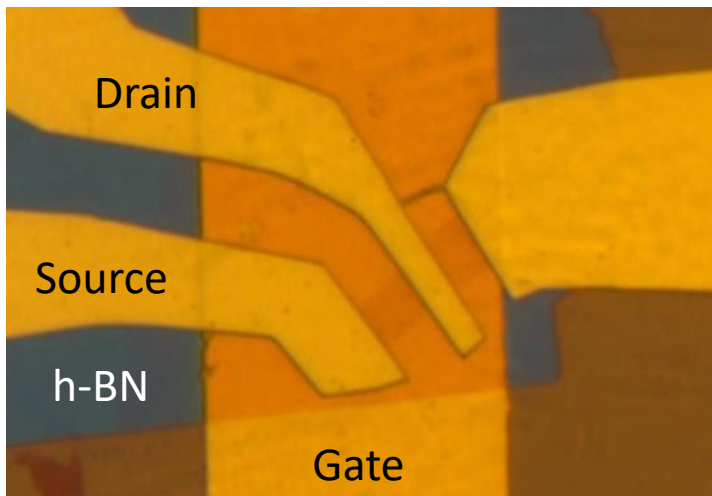
$$n_h = \int_0^{\infty} DOS \times (1-f) dE$$

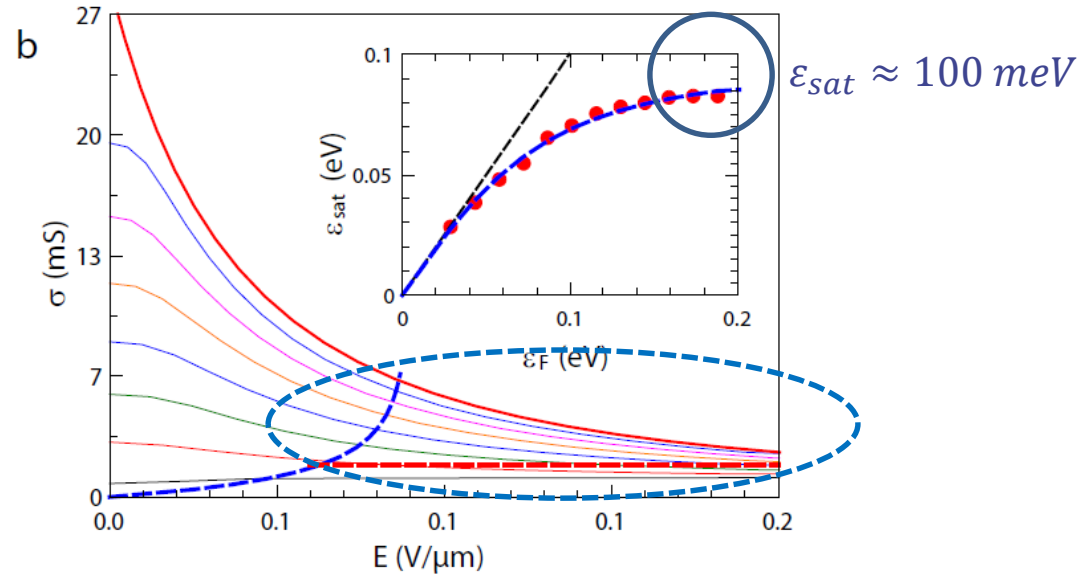
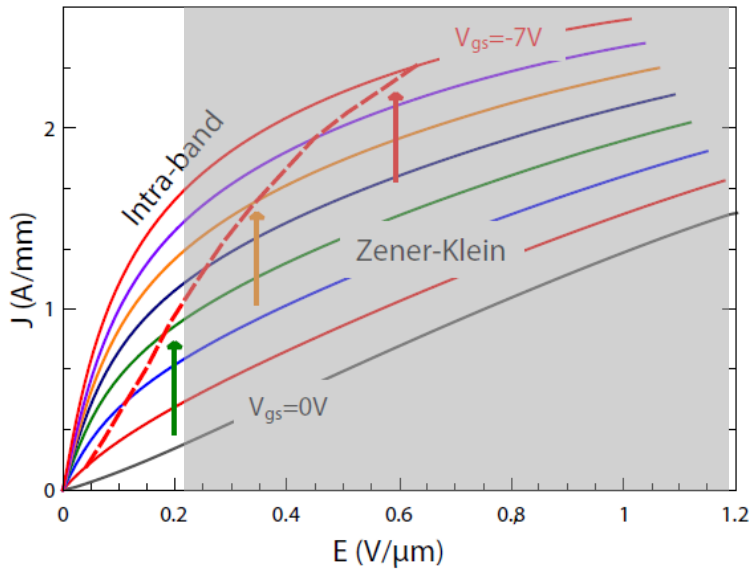
Hot electrons + holes

$$\int_{-\infty}^{\infty} f(1-f) dE$$

$$k_B T_N \approx k_B T_e + \frac{n_h}{DOS}$$

Experiment



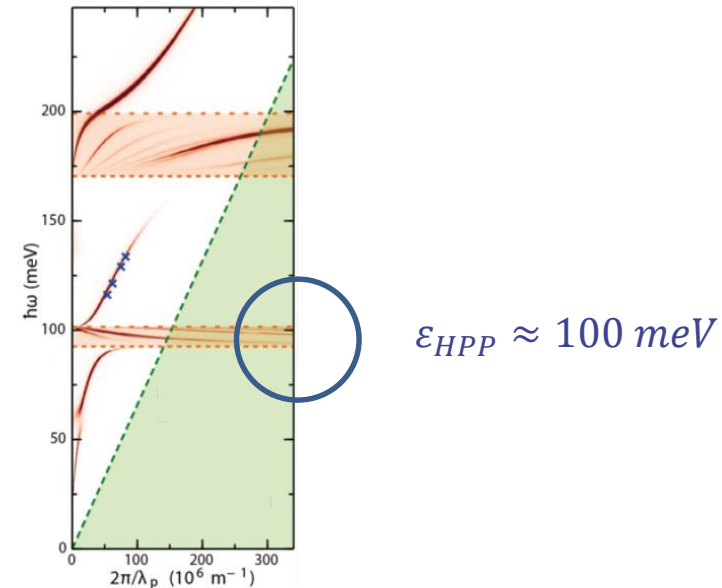


$$\sigma = \frac{ne\mu}{(1 + E/E_{sat})^2} ; \quad v_{sat} = \mu E_{sat} \leq 3 \cdot 10^5 \text{ m/s}$$

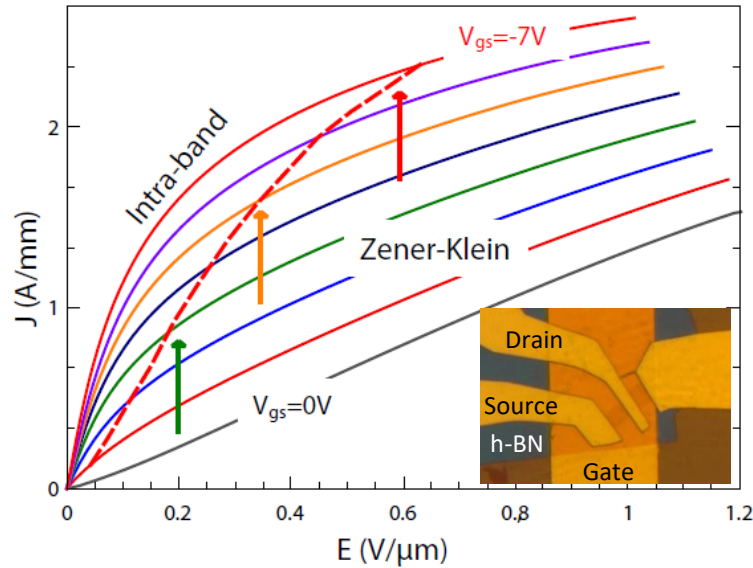
Importantly : $E_{sat} \ll E_{ZK}$

$$J_{sat} = nev_{sat} ; \quad \epsilon_{sat} = \frac{\pi}{2} \hbar k_F v_{sat}$$

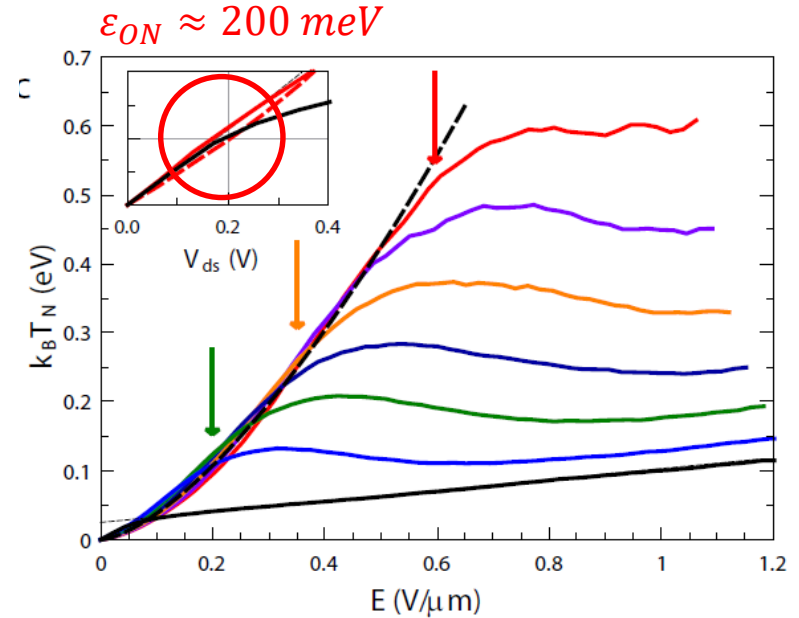
Velocity saturation by type-I hBN phonons



Zener-Klein Transport

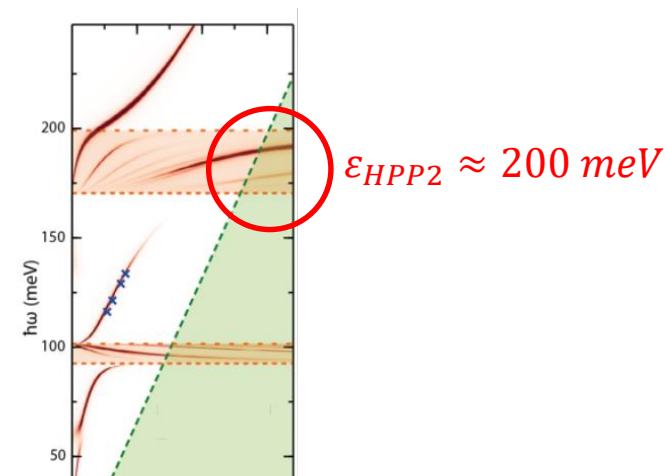


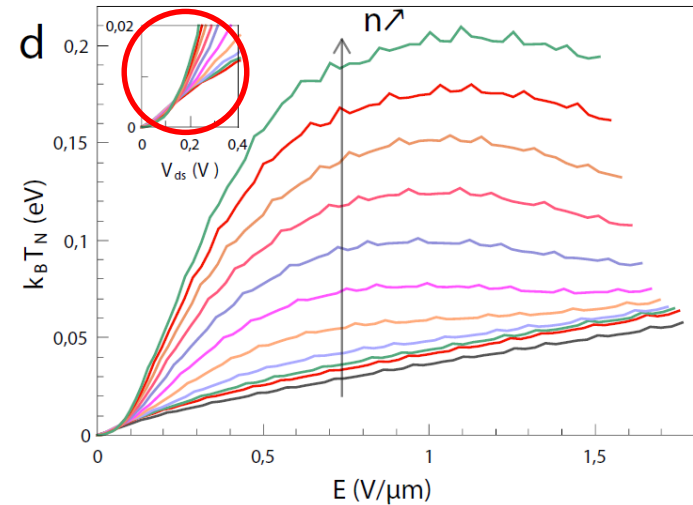
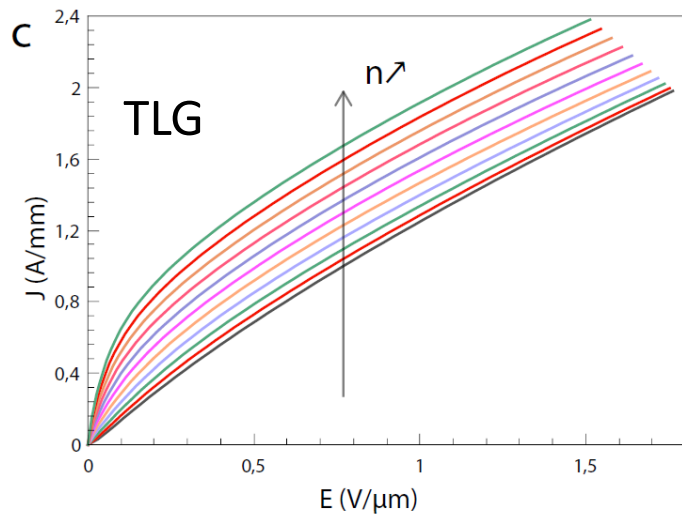
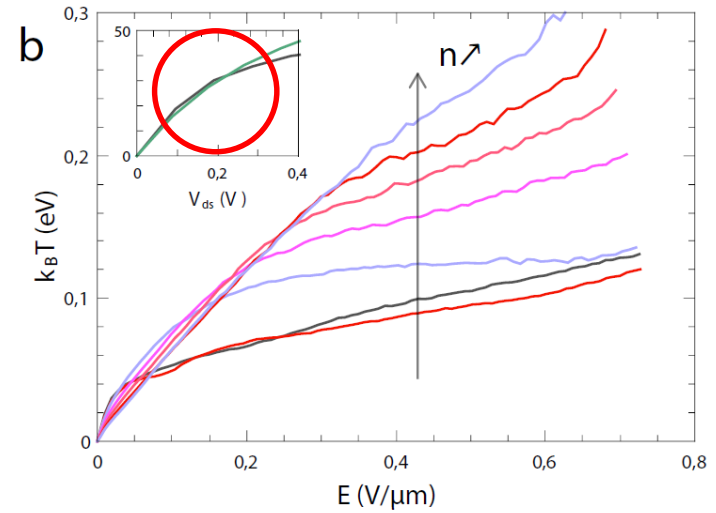
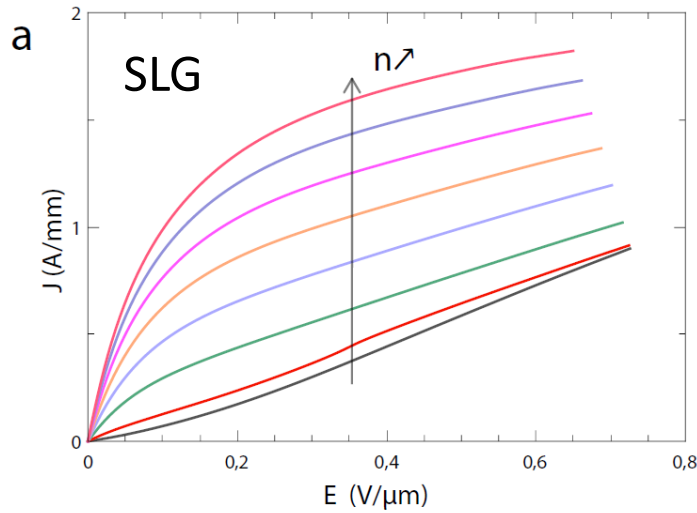
Noise thermometry



Transport is featureless. Main noise features are :

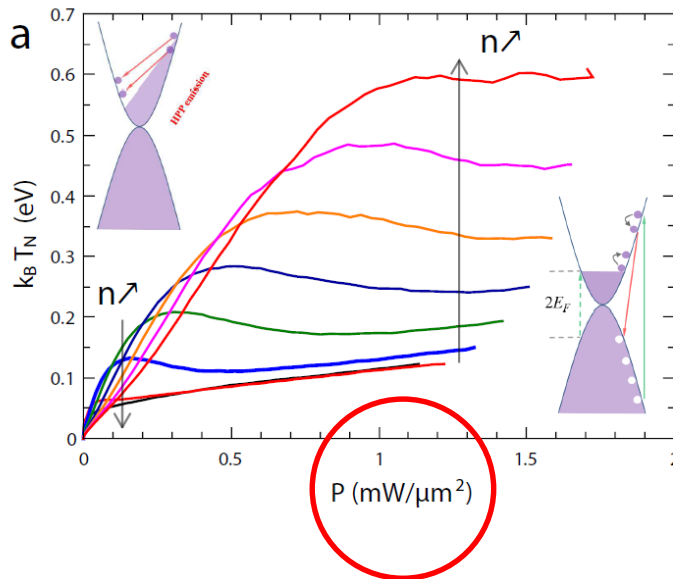
- 1) Superlinear $T_N(E) \Leftrightarrow$ current saturation
- 2) Temperature plateaus in ZKT regime
- 3) Thershold at ZKT onset (arrows)
- 4) Linear $T_N(E)$ at neutrality (ZKT e-h creation)
- 5) Voltage threshold \Leftrightarrow activation energy 200 meV



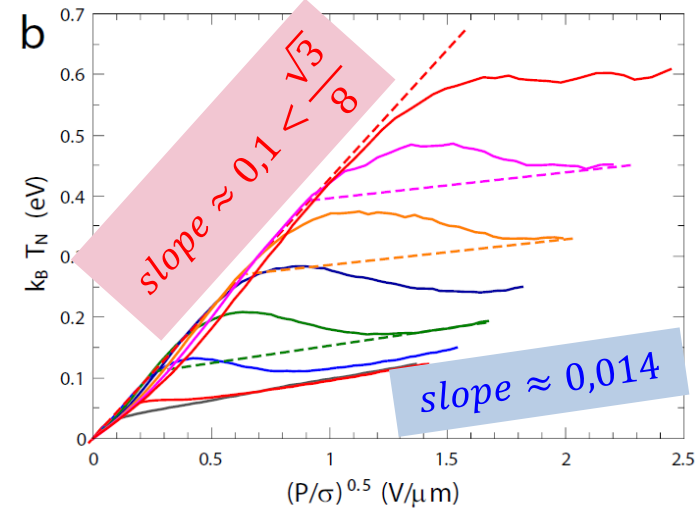


Conventional cooling mechanism ?

Raw noise thermometry



Wiedemann-Frantz analysis



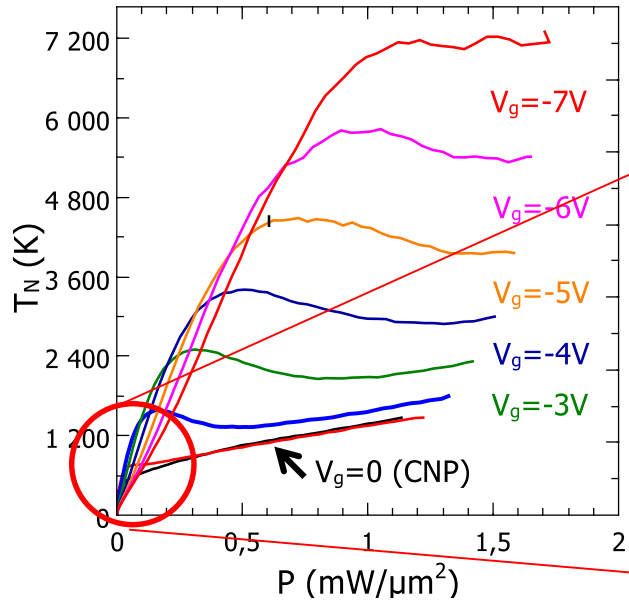
e-e interactions (thermalisation) $\rightarrow \tau_{ee} \sim 20 \text{ fs}$

Wiedemann-Frantz heat conduction $k_B T_N \equiv \langle k_B T_e \rangle = \frac{\sqrt{3}}{8} \times \text{Length} \times \sqrt{P/\sigma}$

~~electron conduction~~

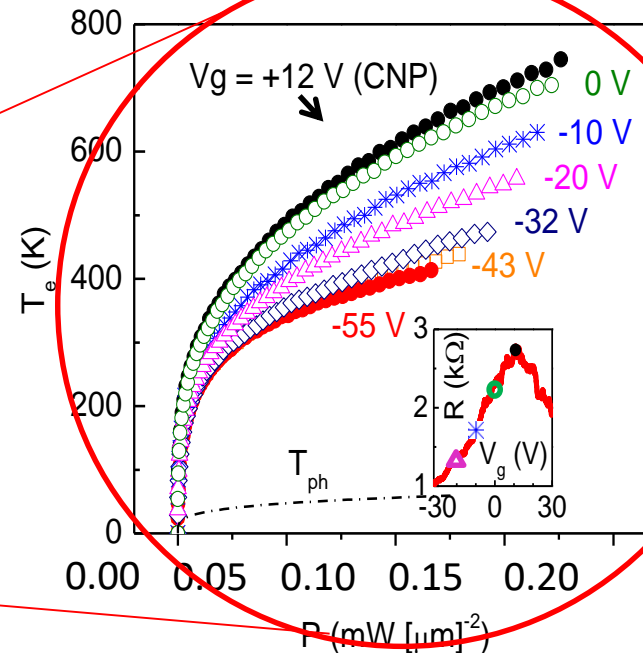
AC phonon cooling ?

this work



Yang et al. / arXiv:1702.02829v1 (2017)

AC phonon cooling

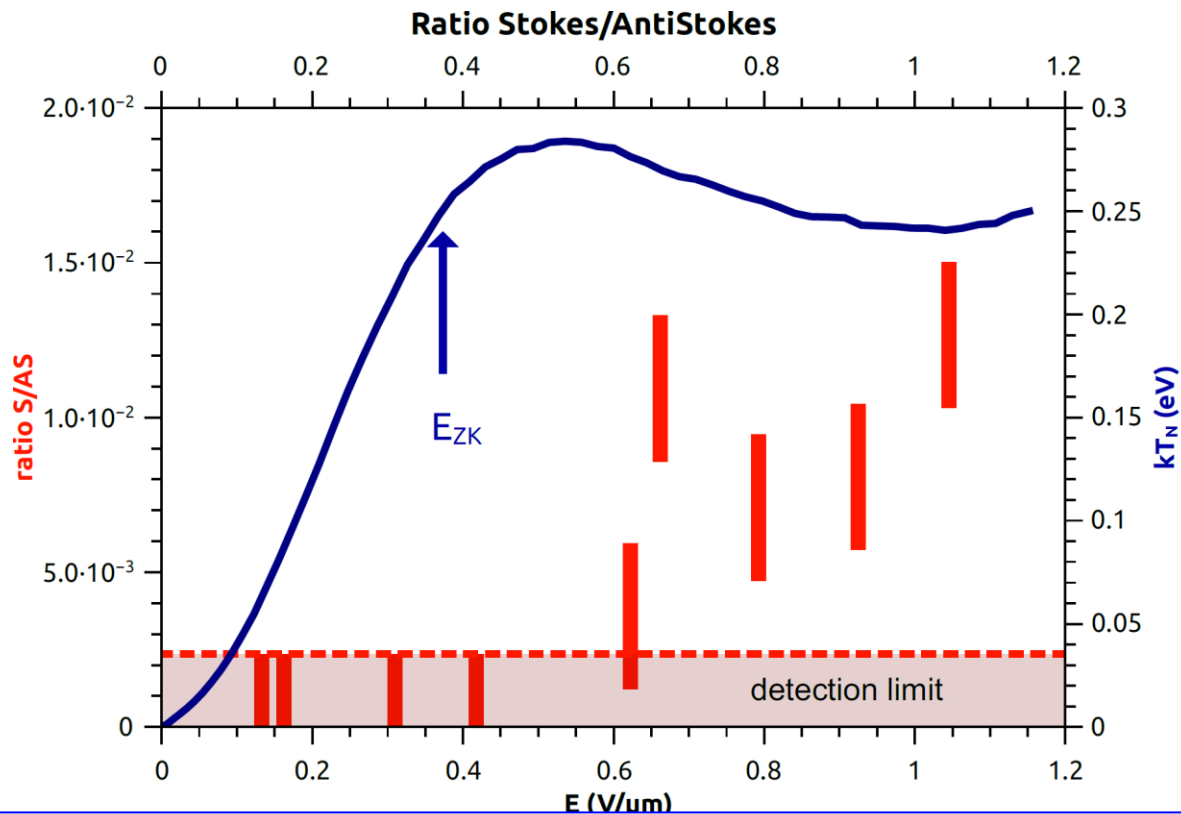


Betz et al. / Phys. Rev. Lett. 109 (2012) 056805;
Betz et al. / Nat. Phys. 9 (2013) 109

Neutral graphene cools better than doped graphene at high bias !

AC phonons

OP phonon cooling ?



e-OP interaction = deformation potential $\Rightarrow \Gamma_{OP} \ll \Gamma_{HPP,SPP}$

OP-phonon Raman thermometry \Rightarrow OP cooling negligible

HPP cooling !

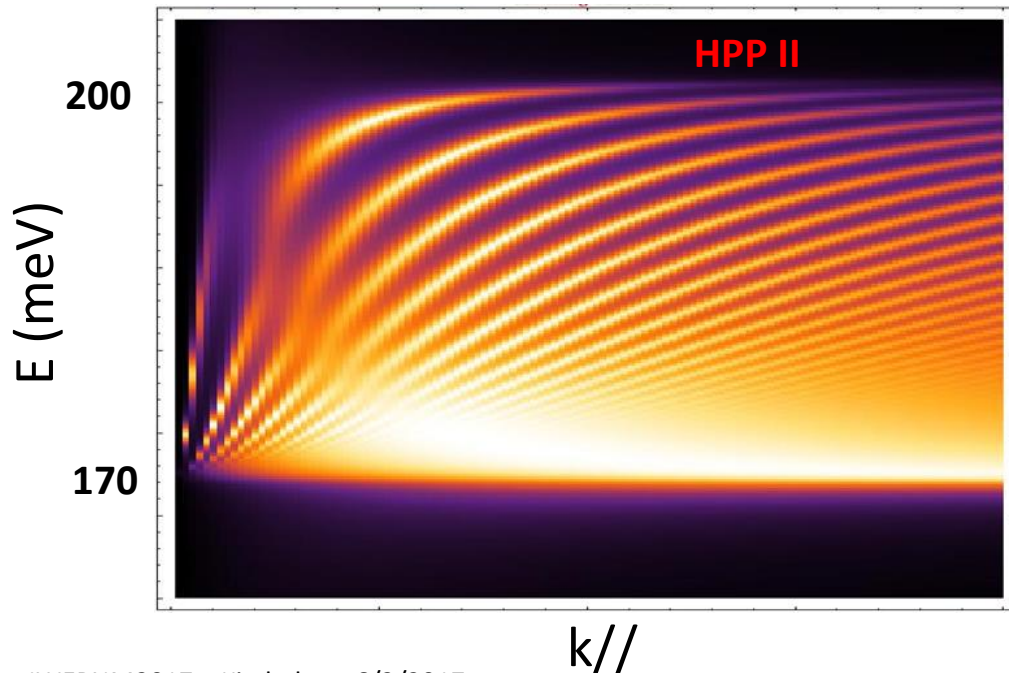
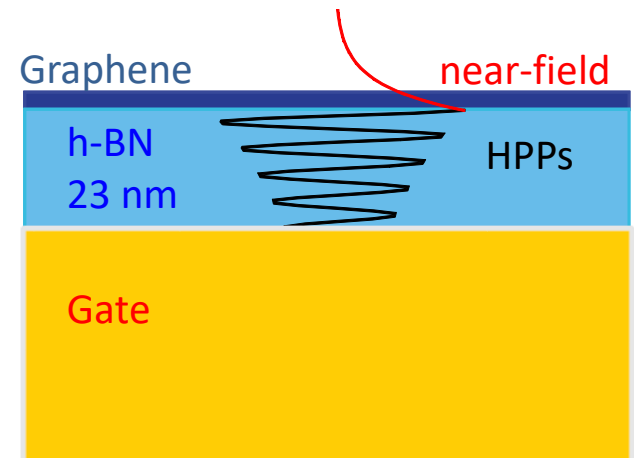
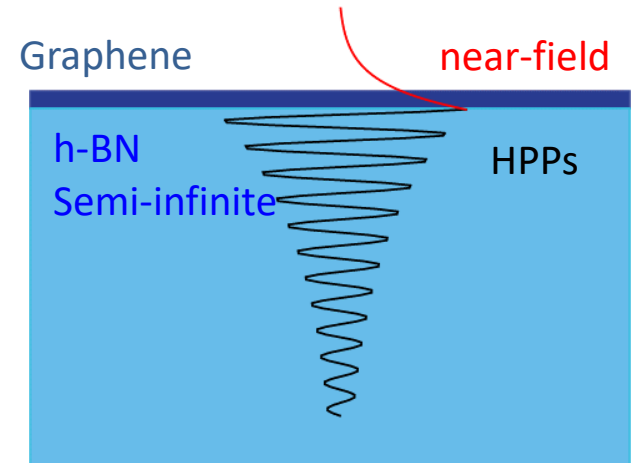
Impedance matching

$$P = \frac{n}{4\pi^2} \frac{\hbar\omega\Delta\omega}{\exp[\hbar\omega/k_B T] - 1} \times M$$

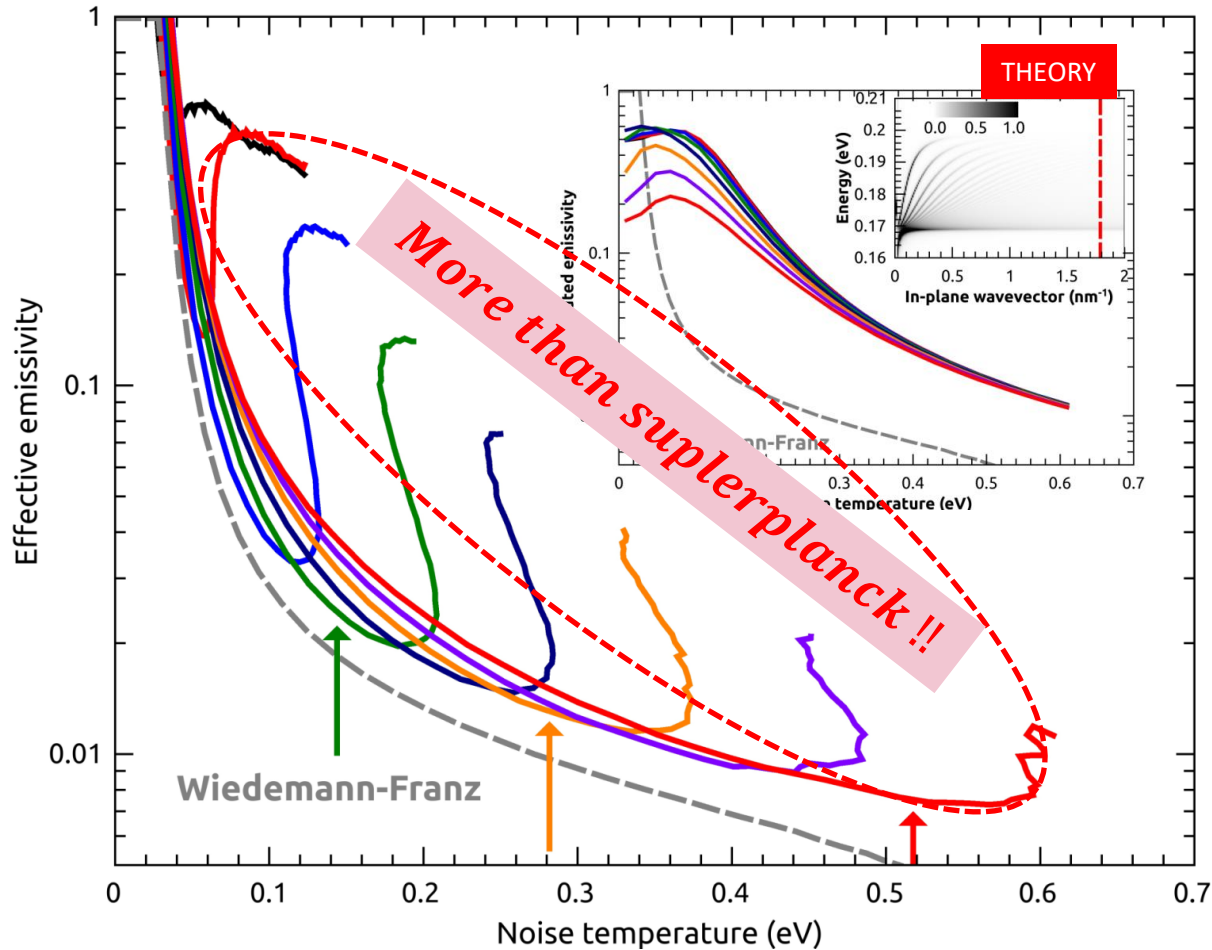
$$M = \left[\frac{4\text{Re}(Y_0)4\text{Re}(\sigma)}{|Y_0 + \sigma|^2} \right] \quad (\text{non-local emissivity})$$

$$\sigma(q, \omega) \quad (\text{non-local graphene conductivity})$$

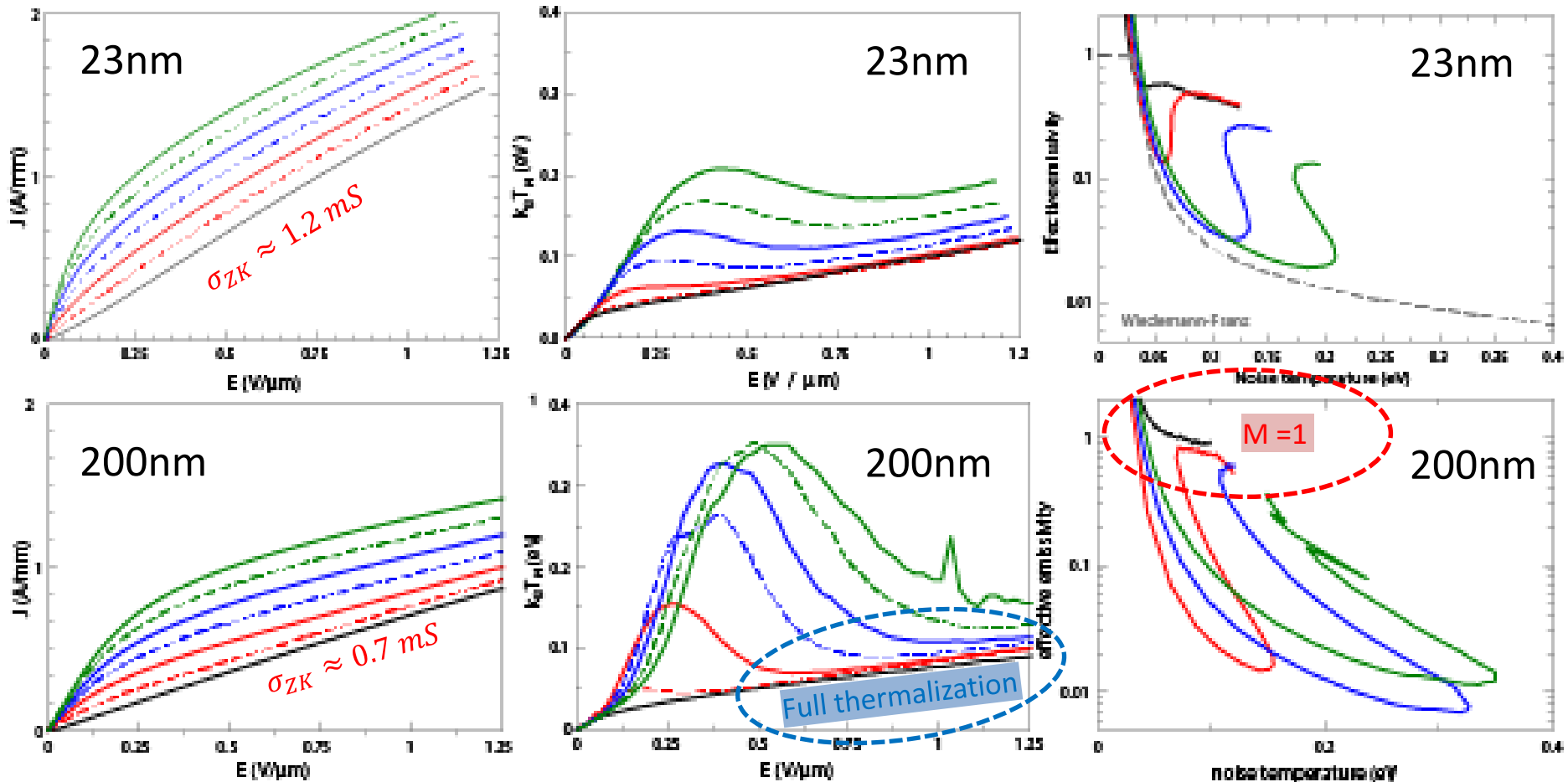
HPPs are propagative modes



EXPERIMENT

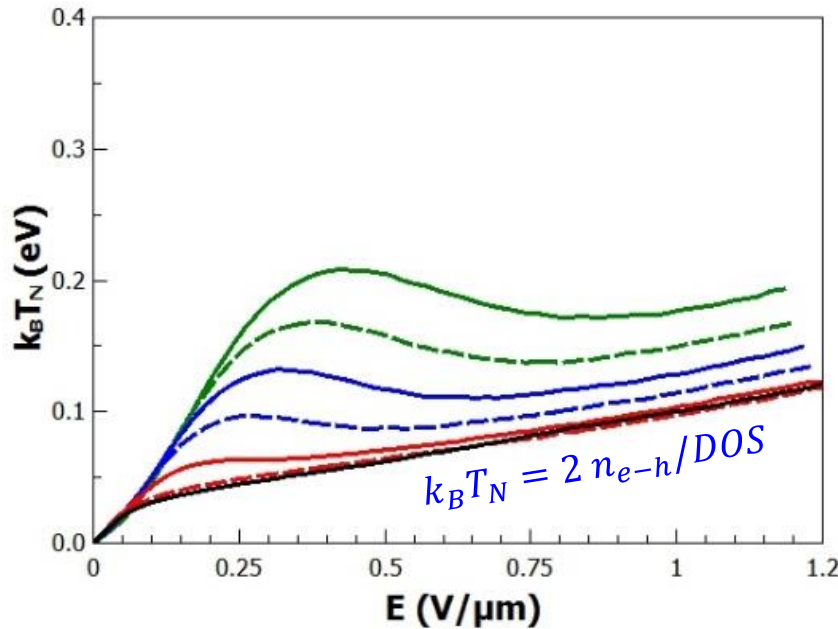


From thin to thick h-BN

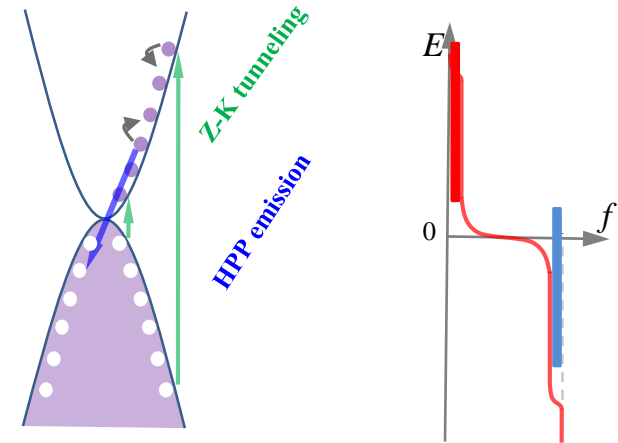


- Same J_{sat}
- **Smaller σ_{ZK}**

- HPP cooling \gg Joule Power in ZKT regime
- **Emissivity ≈ 1 in ZKT regime**



Stationary e-h pair density



e-h pumping rate (th.)

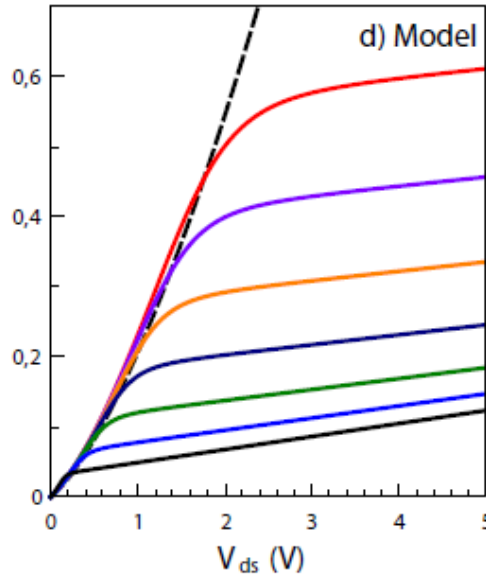
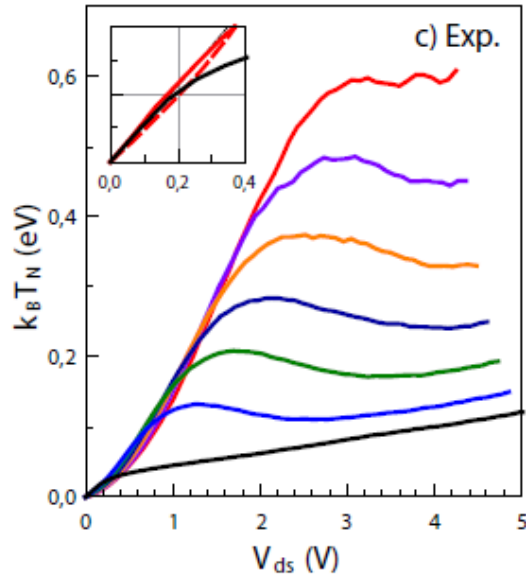
$$\dot{n}_{e-h}^{ZK} = \frac{e k_F}{\pi^2 \hbar} (E - E_{ZK}) = \frac{n_{e-h}}{\tau_{HPP}}$$

e-h density (exp.):

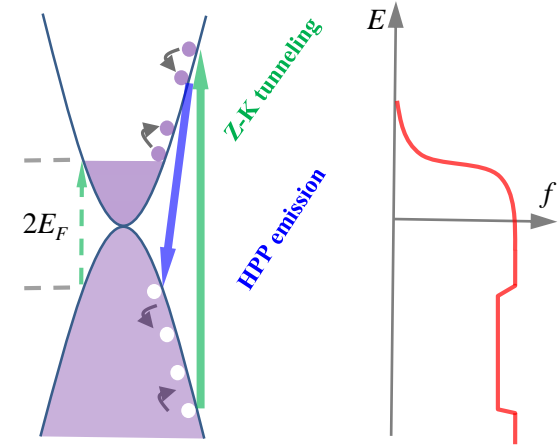
$$n_{e-h} = \text{DOS} \times k_B \Delta T_N / 2$$

HPP cooling rate :

$$\tau_{HPP} \leq 0.46 \text{ ps} \quad \left(\frac{4\pi^2}{\Delta\Omega_{HPP2}} \approx 0.8 \text{ ps} \right)$$



HPP cooling doped regime



$$\text{ZK current : } J_{zk} = \alpha \left(\frac{4e^2}{h} \frac{k_F l_{zk}}{4\pi} \right) (E - E_{zk})$$

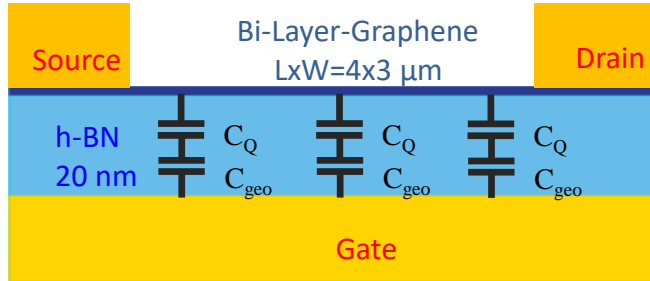
$$\text{ZK pumping : } \dot{n}_{e-h}^{ZK} = \frac{e k_F}{\pi^2 \hbar} (E - E_{ZK})$$

$$\text{HPP cooling : } P_{HPP} = \hbar \Omega \dot{n}_{e-h}^{HPP} = \hbar \Omega \dot{n}_{e-h}^{ZK} = \hbar \Omega \frac{e k_F}{\pi^2 \hbar} (E - E_{zk})$$

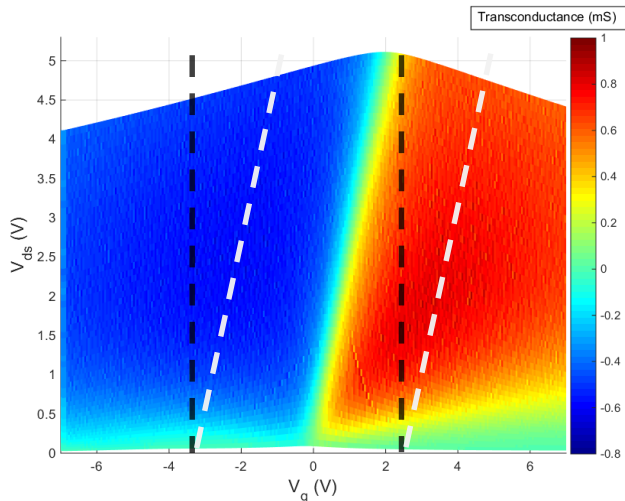
$$\text{Joule Heating : } \Delta P_{Joule} = J_{sat} (E - E_{sat}) = 2\epsilon_{sat} \frac{e k_F}{\pi^2 \hbar} (E - E_{sat})$$

$$\text{in GoBN, where } \hbar \Omega_{II} \approx 2\hbar \Omega_I \approx 200 \text{ meV} \Rightarrow P_{HPP} \approx P_{Joule}$$

ZKT-FETs as power amplifiers with efficient HPP cooling

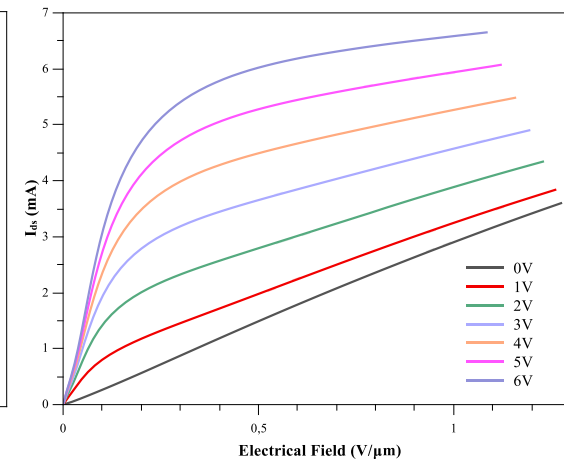
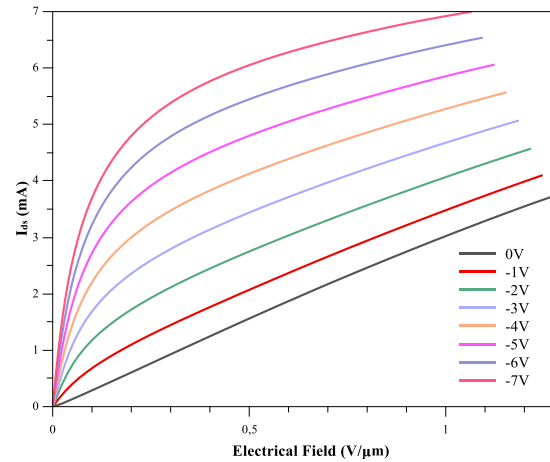


Transconductance

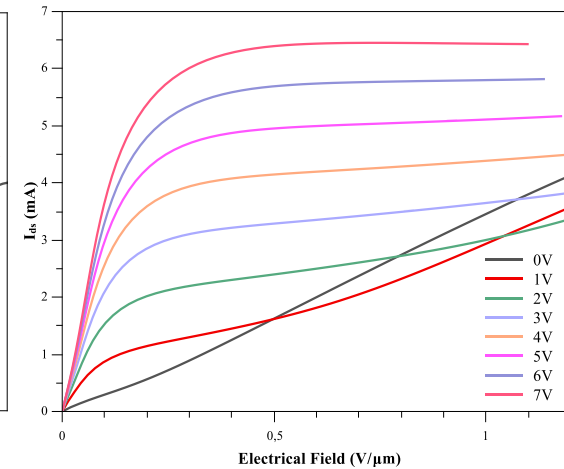
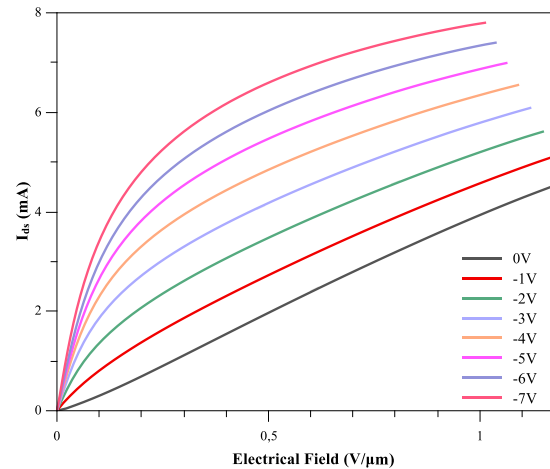


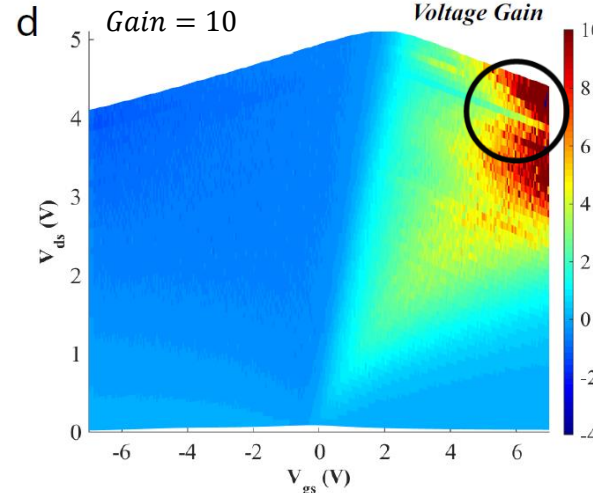
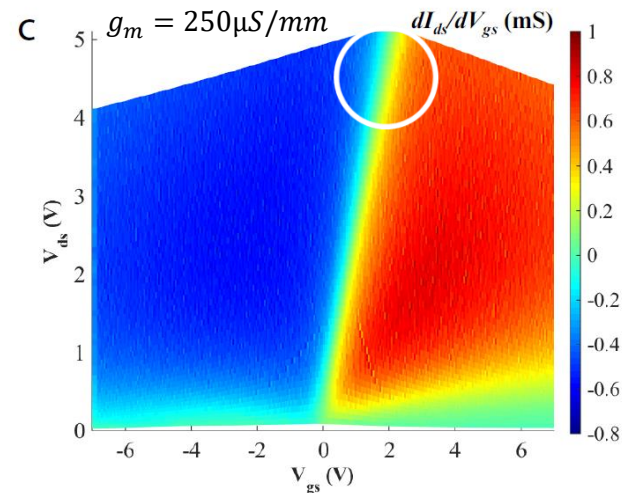
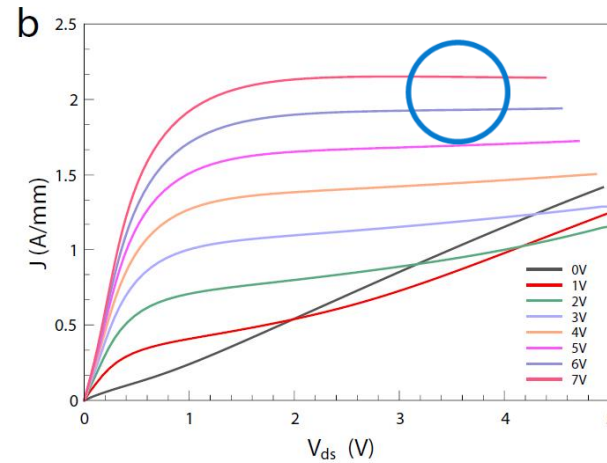
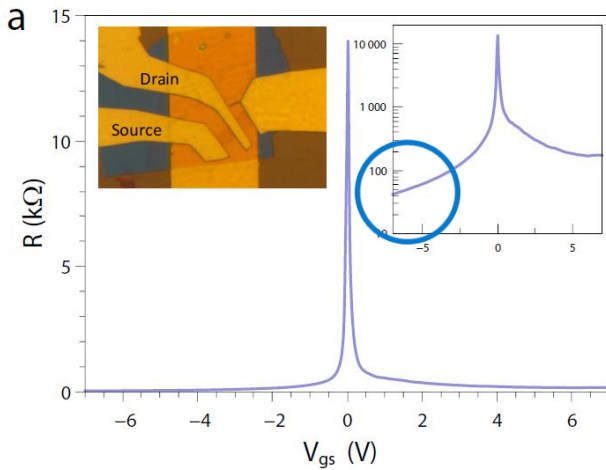
Graphene 2017, opto-electronics

Constant carrier density



Constant gate voltage





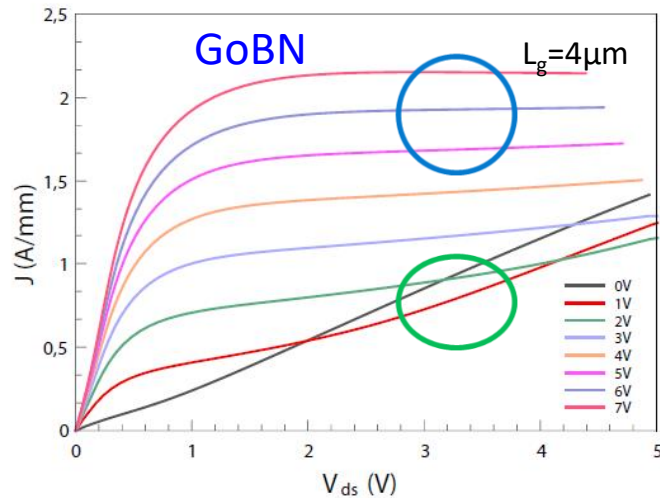
GoBN

- High mobility ($30\,000\text{ cm}^2/\text{V/s}$)
- Low contact resistance
- Current saturation ++
- High-power ++
- Zener-Klein regime operation ++
- Negligible self heating effects

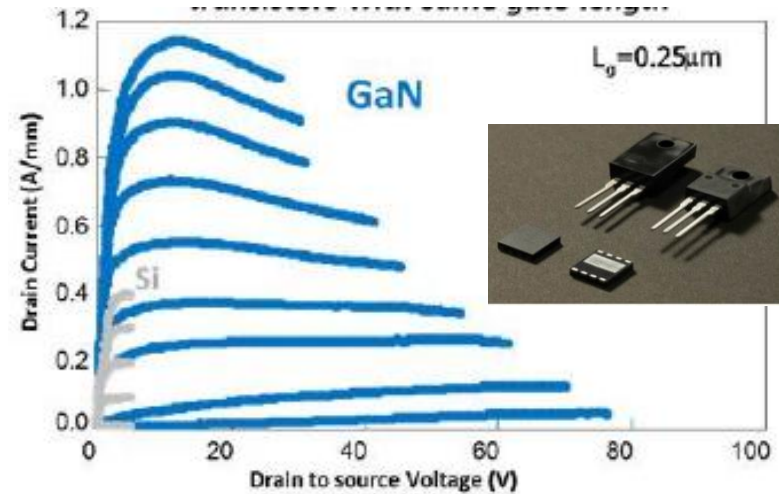
Bottom gating

- Drain gating (bottom gate effect)
- Transconductance ($250\ \mu\text{S}/\text{mm}$)
- Large voltage gain ($G \sim 10$)

GoBN Zener-Klein transistor



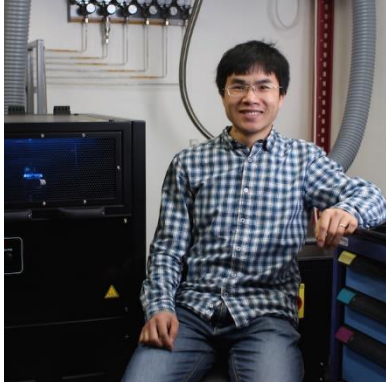
Panasonic : X-GaN Power transistor



5 merits of h-BN

1. High mobility
2. Large saturation currents (power amplification ?)
3. Pinchoff replaced by Zener-Klein tunneling
4. Compensation of ZK tunneling by a bias induced doping depletion
5. No thermal degradation => cooling by hyperbolic hBN phonons !!!

1. *HPP cooling promotes h-BN is the ideal heat sink*
2. *Zener-Klein Tunneling optimizes HPP emission*
3. *ZKT-FETs are promising high power transistors*



W. Yang (post-doc), S. Berthou (PhD student), Q. Wilmart (PhD student)

A. Denis, M. Rosticher (LPA, RF electronics and clean room engineers)

X. Lu, G. Zhang (Beijing, sample fabrication)

T. Taniguchi, K. Watanabe (NIMS, hBN crystals)

G. Fève, J.M. Berroir, BP, C. Voisin, E. Baudin (LPA meso / optics groups)