

Graphene 2017

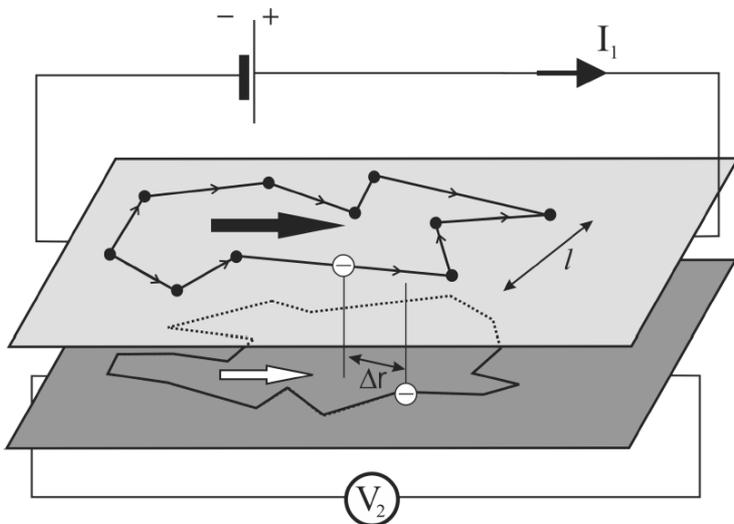
Coulomb Drag in Graphene

-Toward Exciton Condensation

Philip Kim

Department of Physics, Harvard University

Coulomb Drag



Drag Resistance:

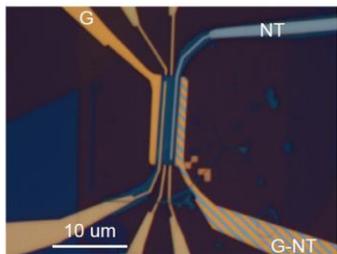
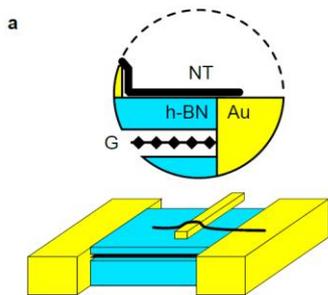
$$R_D = V_2 / I_1$$

Onsager Reciprocity

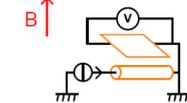
$$V_2(B) / I_1 = V_1(-B) / I_2$$

Price *et al.*, Science (2012)

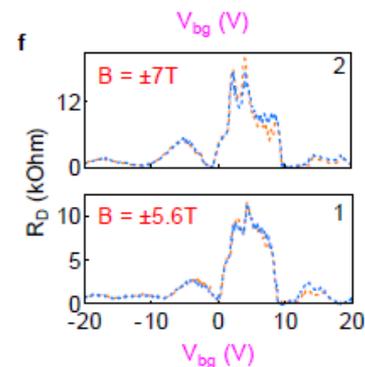
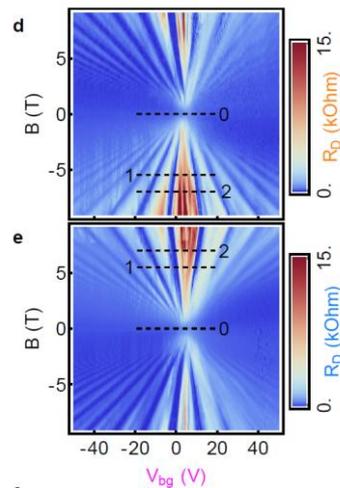
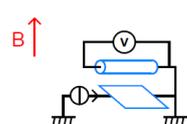
Coulomb drag between graphene/nanotube



Drag in graphene



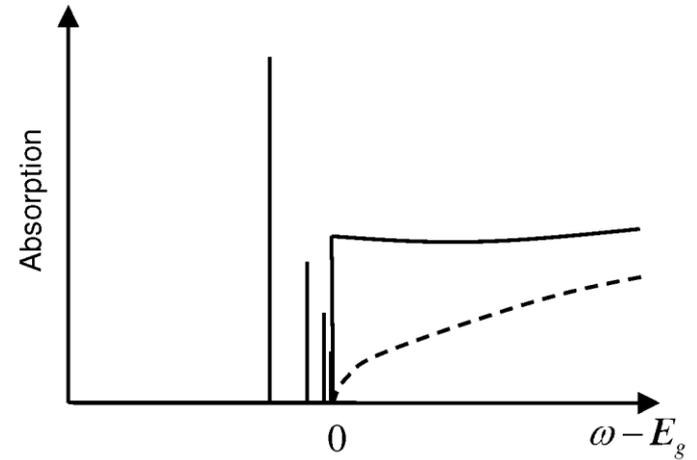
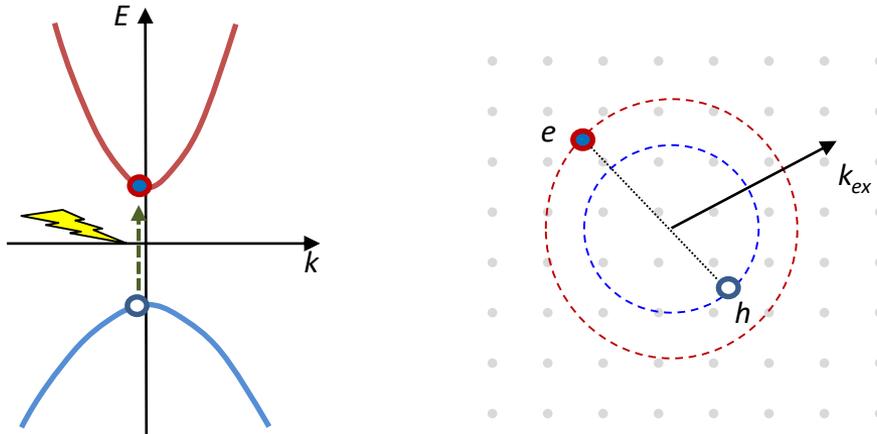
Drag in nanotube



Pillet *et al.*, arXiv:1612.05992

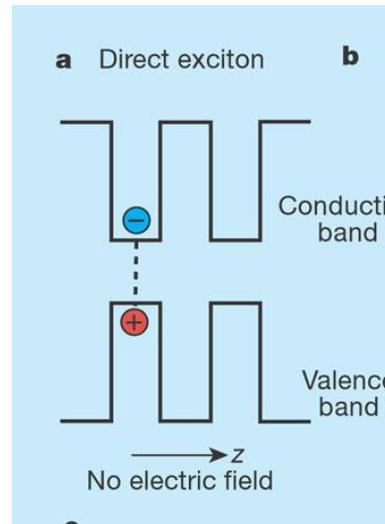
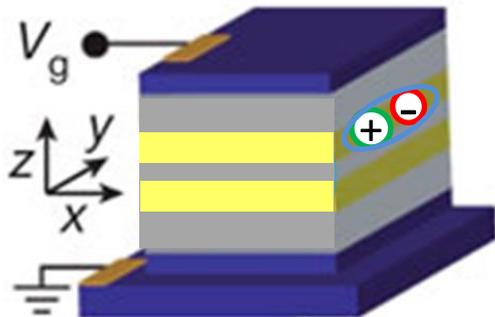
Excitons

Excitons in Semiconductors

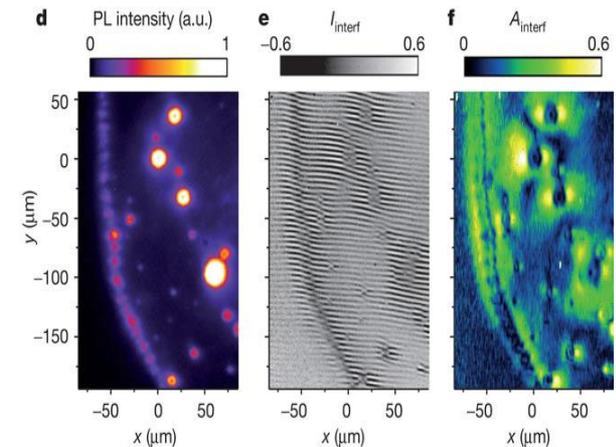


Direct and indirect excitons in semiconducting quantum wells

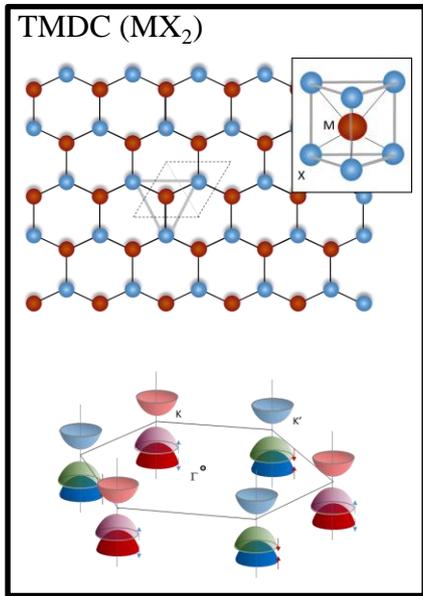
Semiconductor heterostructure



Spontaneous coherence



Excitons in 2D Materials



- Strong confinement
- Lack of strong screening effect
- Spin polarized band
- Spin-valley locking

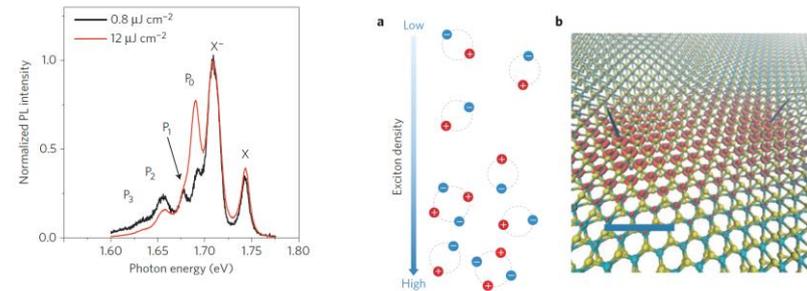
nature
physics

LETTERS

PUBLISHED ONLINE: 11 MAY 2015 | DOI: 10.1038/NPHYS3324

Observation of biexcitons in monolayer WSe_2

Yumeng You^{1,2†}, Xiao-Xiao Zhang^{2†}, Timothy C. Berkelbach³, Mark S. Hybertsen⁴, David R. Reichman³ and Tony F. Heinz^{2*‡}



RESEARCH | REPORTS

688 12 FEBRUARY 2016 • VOL 351 ISSUE 6274

VALLEYTRONICS

Valley-polarized exciton dynamics in a 2D semiconductor heterostructure

Pasqual Rivera,^{1*} Kyle L. Seyler,^{1*} Hongyi Yu,² John R. Schaibley,¹ Jiaqiang Yan,^{3,4} David G. Mandrus,^{3,4,5} Wang Yao,² Xiaodong Xu^{1,6‡}

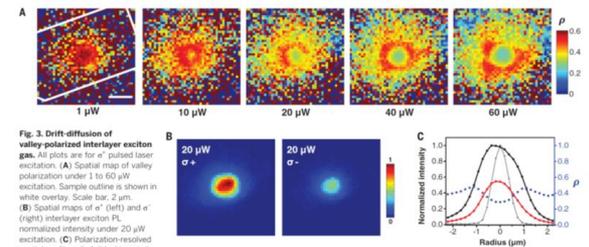
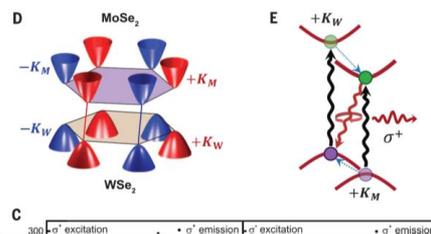
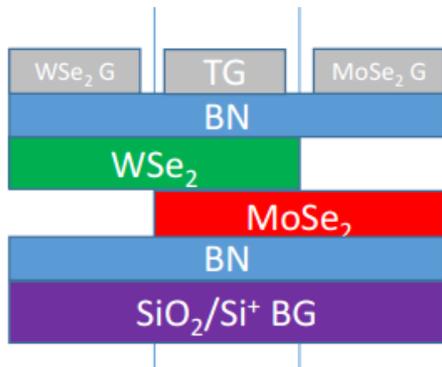
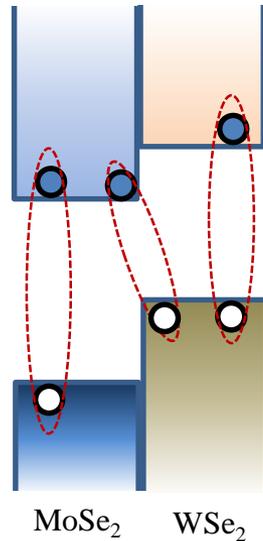
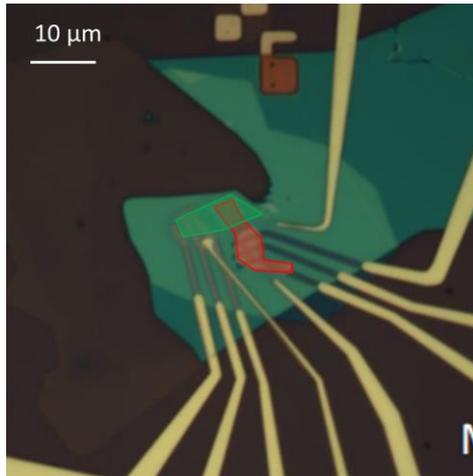


Fig. 3. Drift-diffusion of valley-polarized interlayer exciton gas. All plots are for e^- pulsed laser excitation. (A) Spatial map of valley polarization under 1 to 60 μW excitation. Sample outline is shown in white overlay. Scale bar, 2 μm . (B) Spatial maps of σ^+ (left) and σ^- (right) interlayer exciton PL normalized intensity under 20 μW excitation. (C) Polarization-resolved spatial profiles of σ^+ (black) and σ^- (red) components of interlayer exciton PL under 40 μW excitation. The spatial distribution of valley polarization is shown in blue, and the laser excitation profile is shown in gray. Line cuts are radially averaged through the excitation center, and curves are added as guides to the eye.

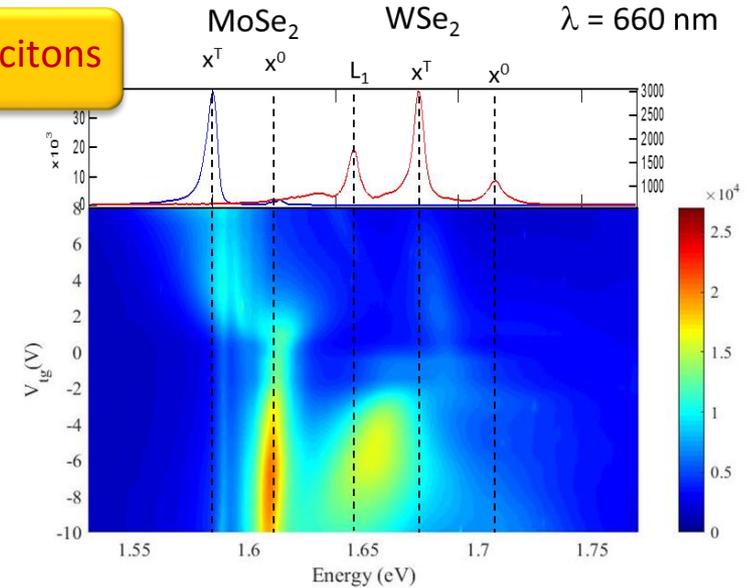
Interlayer Excitons in TMDCs

MoSe₂/WSe₂ heterostructures

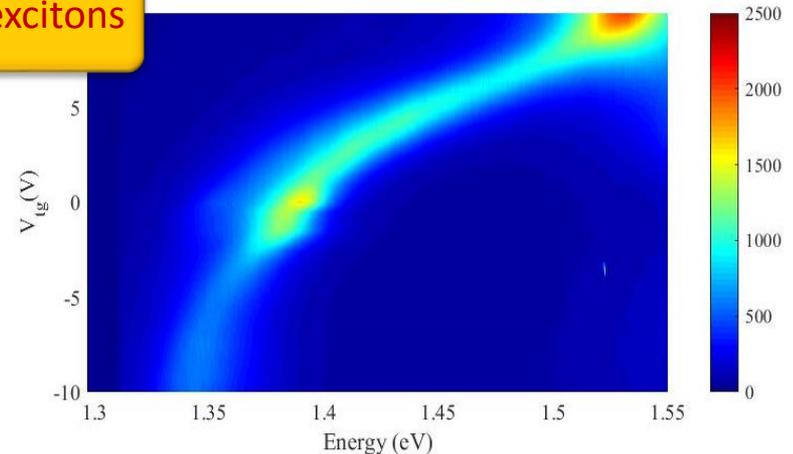


Gate dependent photo luminescent

Intralayer excitons



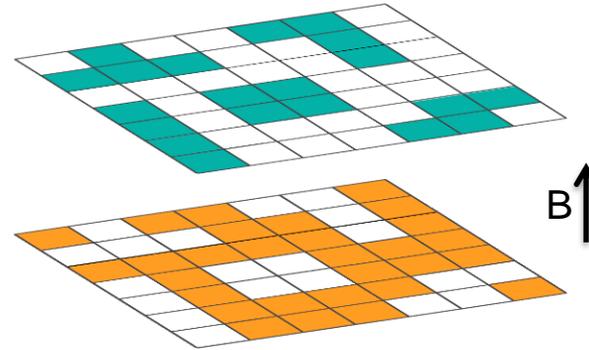
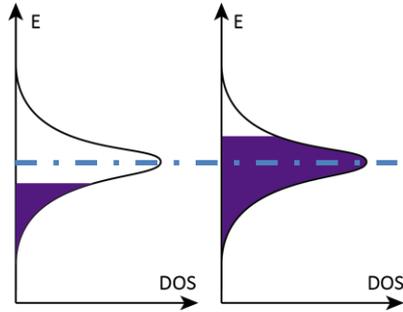
Interlayer excitons



Exciton condensation between Landau levels

J. P. Eisenstein, Annu. Rev. Condens. Matter Phys. **5**, 159 (2014).

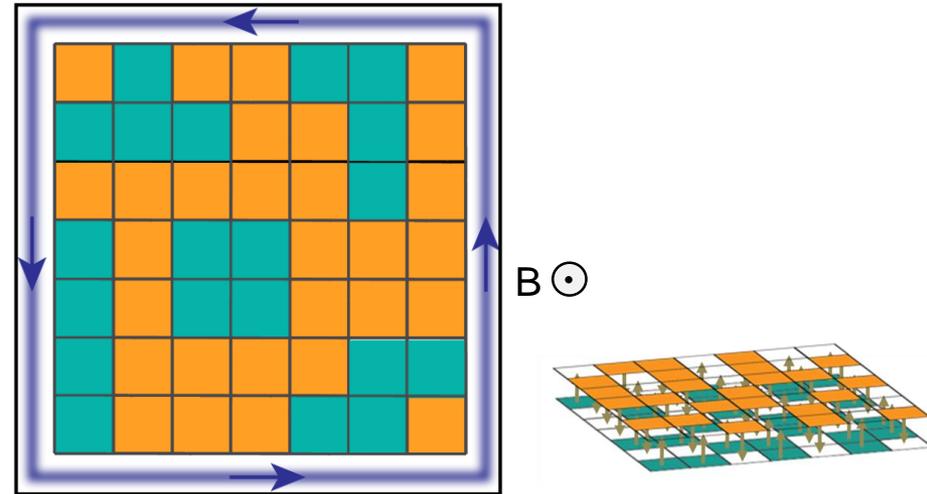
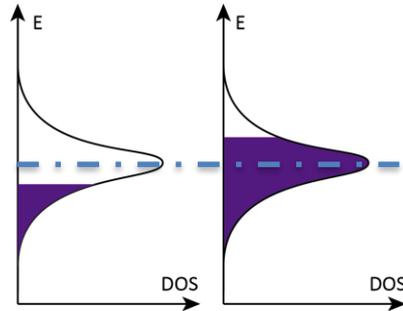
Two partially filled
Landau levels



Exciton condensation between Landau levels

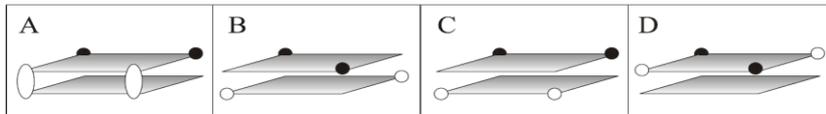
J. P. Eisenstein, Annu. Rev. Condens. Matter Phys. **5**, 159 (2014).

Two partially filled Landau levels

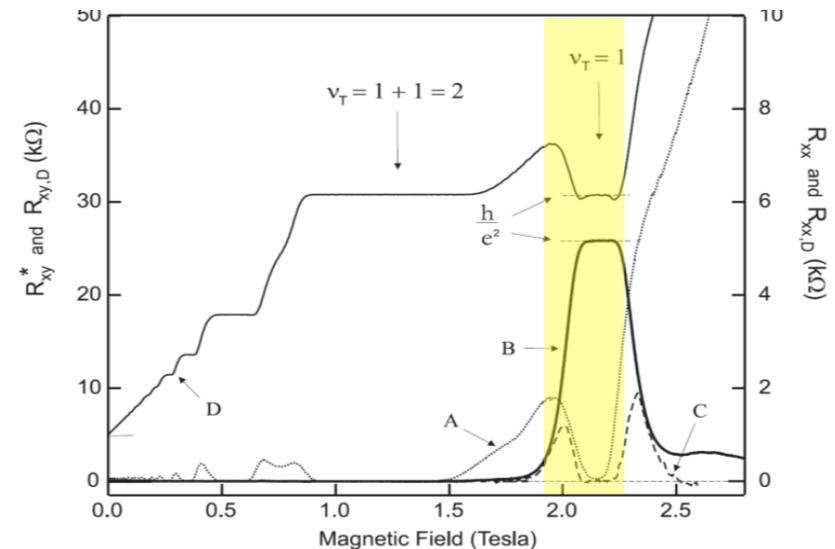


Total Landau level quantum Hall effect

GaAs Double Quantum Well



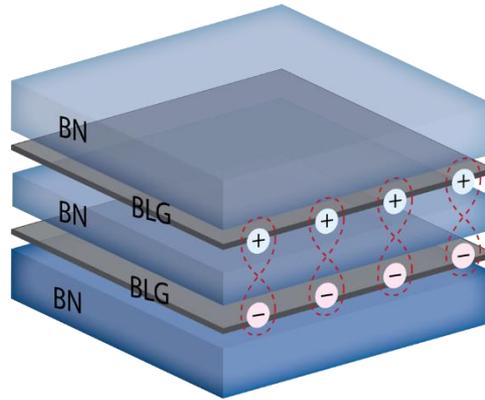
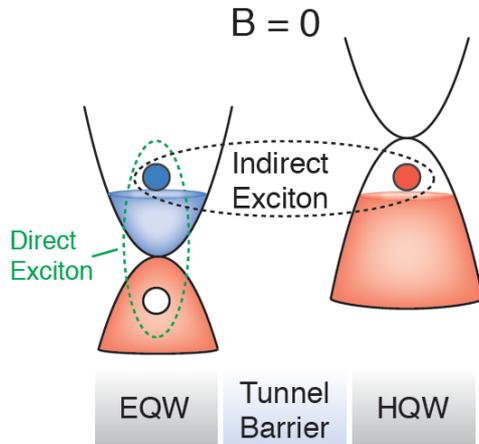
- Quantum Hall effect for two partially filled complementary LLs
- Quantized drag Hall



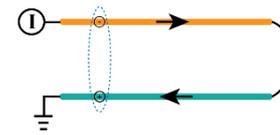
M. Kellogg, et. al, PRL (2002)

Exciton Condensation: Currents

Exciton condensation (exciton insulator)

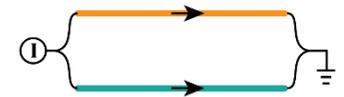


Counter Flow



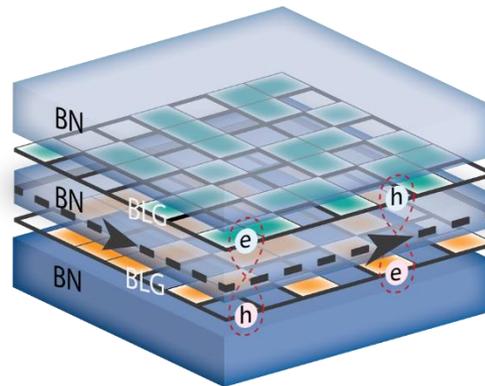
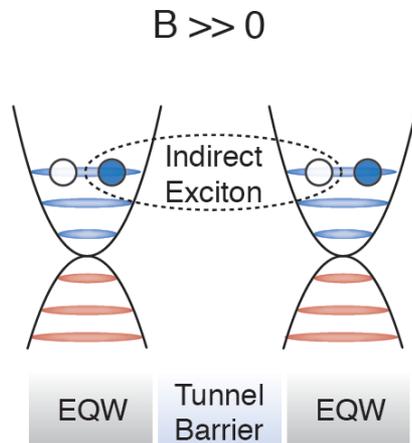
$$R_{CF} = 0$$

Symmetric Flow



$$R_{sym} = \infty$$

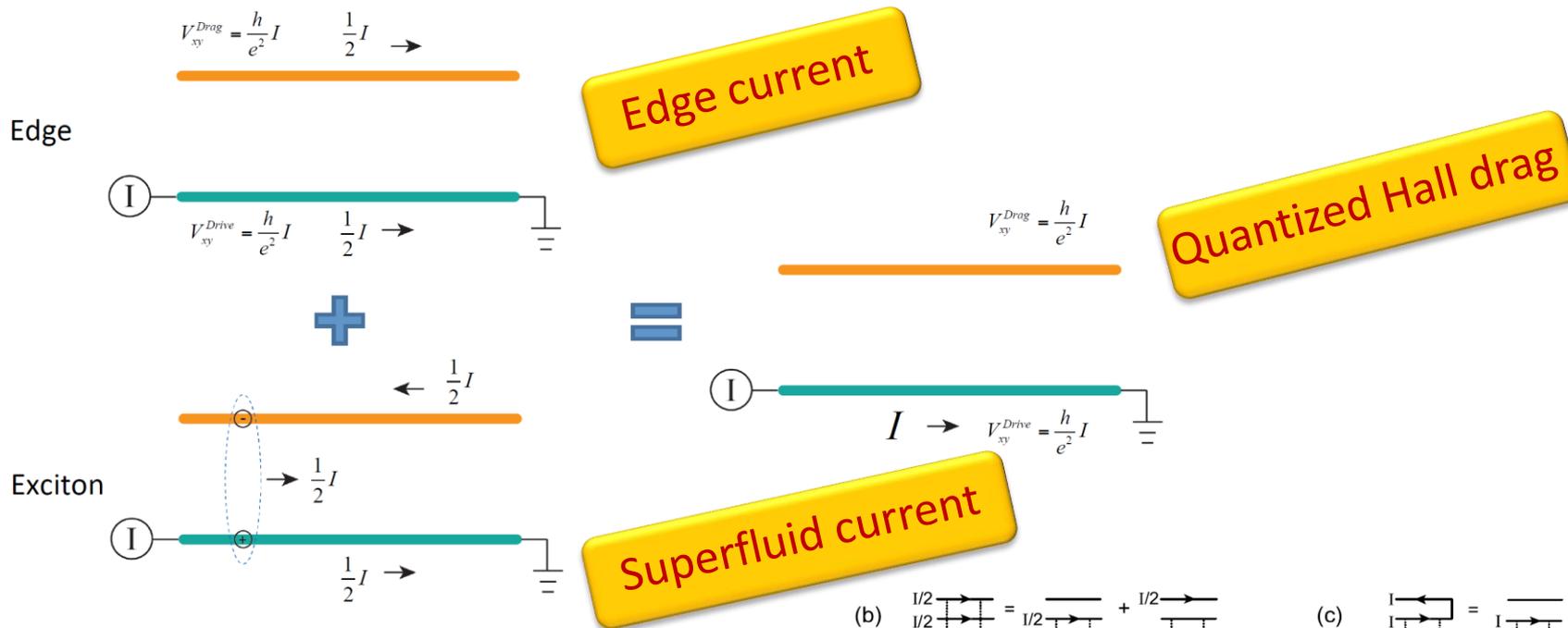
Exciton condensation between LL (topological exciton insulator)



$$R_{xx}^{CF} = 0 \quad R_{xx}^{sym} = 0$$

$$R_{xy}^{CF} = 0 \quad R_{xy}^{sym} = \frac{h}{\nu_{tot} e^2}$$

Exciton Current and Quantized Drag



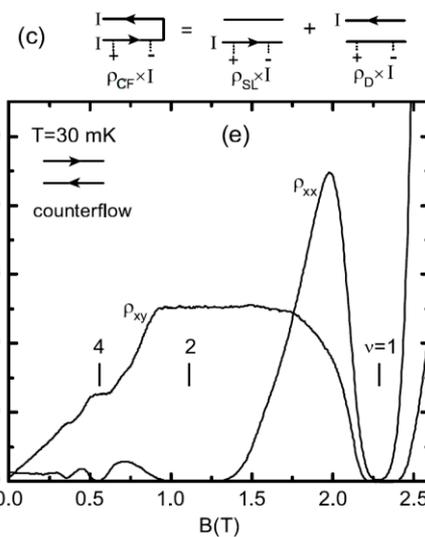
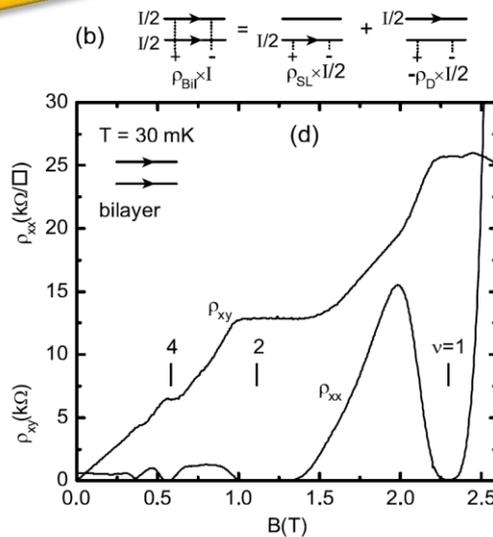
NUMBER 3 PHYSICAL REVIEW LETTERS 16.

Counterflow Measurements in Strongly Correlated GaAs Hole Bilayers: Evidence for Electron-Hole Pairing

E. Tutuc and M. Shayegan

Department of Electrical Engineering, Princeton University, Princeton, New Jersey 08544, USA

D. A. Huse



Outlines

Frictional magneto-Coulomb drag in graphene hexa boron nitride heterostructure

Xiaomeng Liu,¹ Lei Wang,² Kin Chung Fong,³ Yuanda Gao,² Patrick Maher,⁴
Kenji Watanabe,⁵ Takashi Taniguchi,⁵ James Hone,² Cory Dean,⁴ and Philip Kim¹

arXiv:1612.08308

Frictional magneto drag effect in semiclassical regime

Quantum Hall Drag of Exciton Superfluid in Graphene

Xiaomeng Liu¹, Kenji Watanabe², Takashi Taniguchi², Bertrand I. Halperin¹, Philip Kim¹

Quantized drag in excitonic regime

arXiv:1608.03726

Coulomb Drag Experiment in Graphene

PHYSICAL REVIEW B 83, 161401(R) (2011)



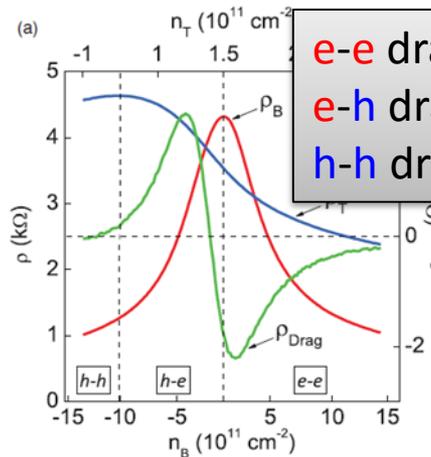
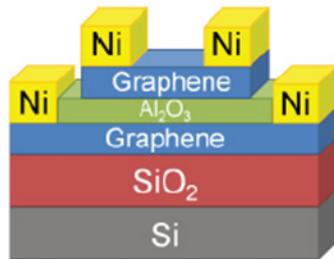
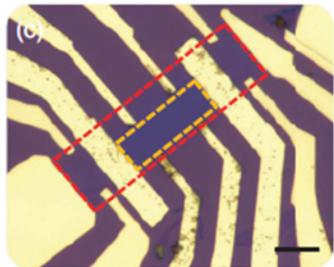
Coulomb drag of massless fermions in graphene

Seyoung Kim,¹ Insun Jo,² Junghyo Nah,¹ Z. Yao,² S. K. Banerjee,¹ and E. Tutuc^{1,*}

¹Microelectronics Research Center, The University of Texas at Austin, Austin, Texas 78758, USA

²Department of Physics, The University of Texas at Austin, Austin, Texas 78712, USA

(Received 9 October 2010; revised manuscript received 17 March 2011; published 8 April 2011)



e-e drag: negative
e-h drag: positive
h-h drag: negative

ARTICLES

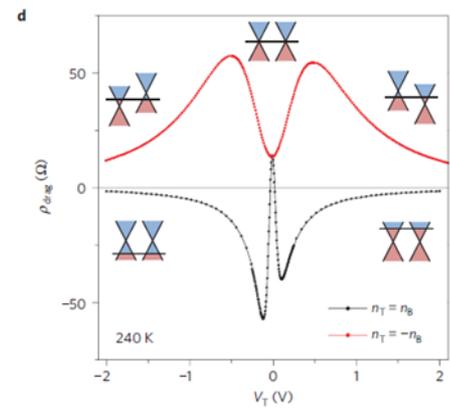
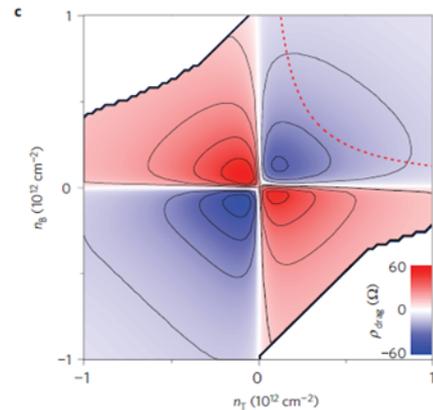
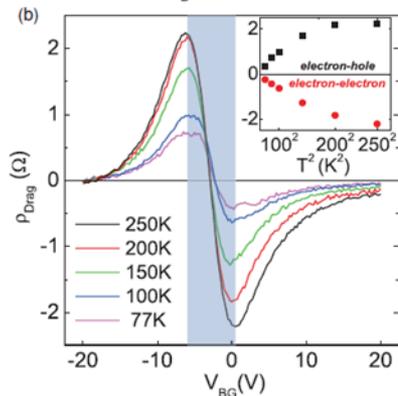
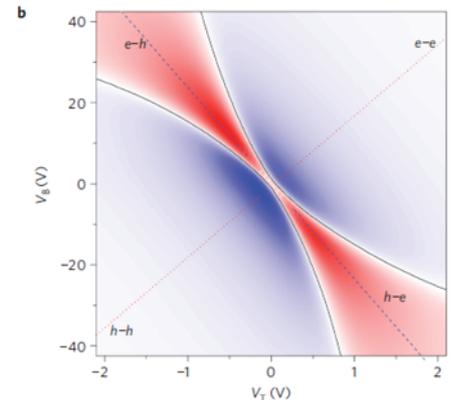
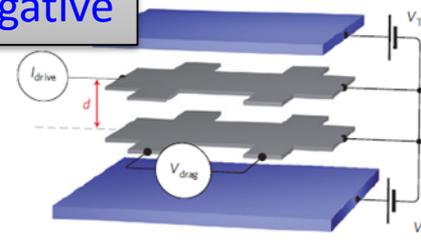
PUBLISHED ONLINE: 14 OCTOBER 2012 | DOI: 10.1038/NPHYS2441

nature
physics

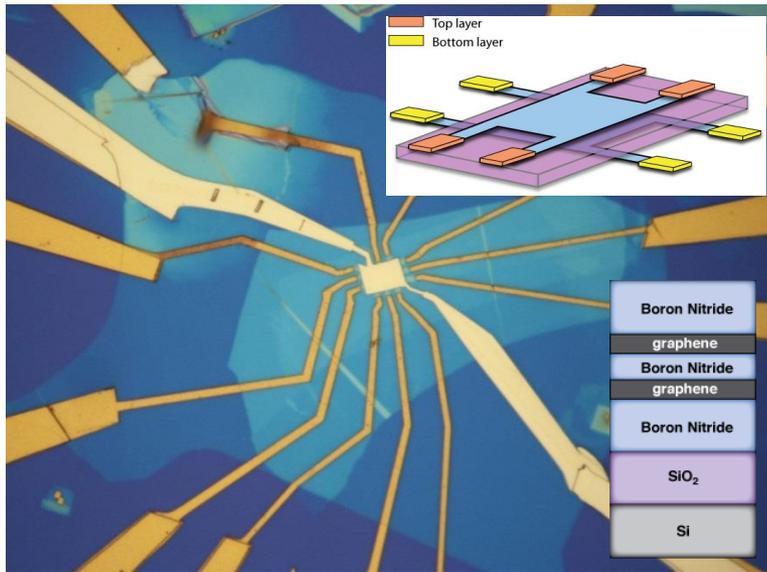
Strong Coulomb drag and broken symmetry in double-layer graphene

R. V. Gorbachev¹, A. K. Geim^{1,2,*}, M. I. Katsnelson³, K. S. Novoselov², T. Tudorovskiy³,
I. V. Grigorieva^{1,2}, A. H. MacDonald⁴, S. V. Morozov^{2,5}, K. Watanabe⁶, T. Taniguchi⁶

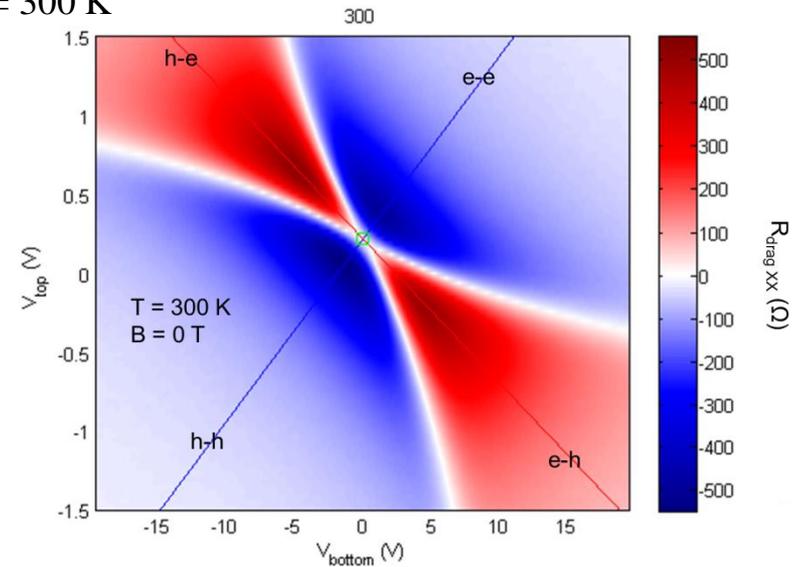
and M. M. Marenko^{1,2,*}



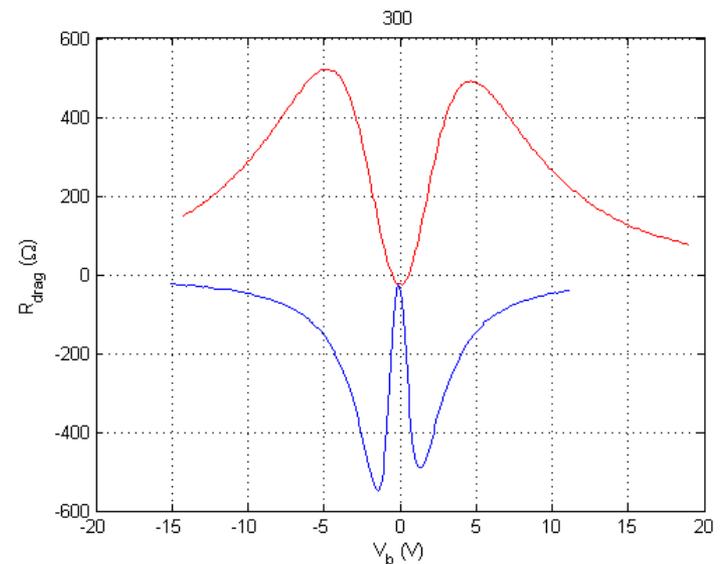
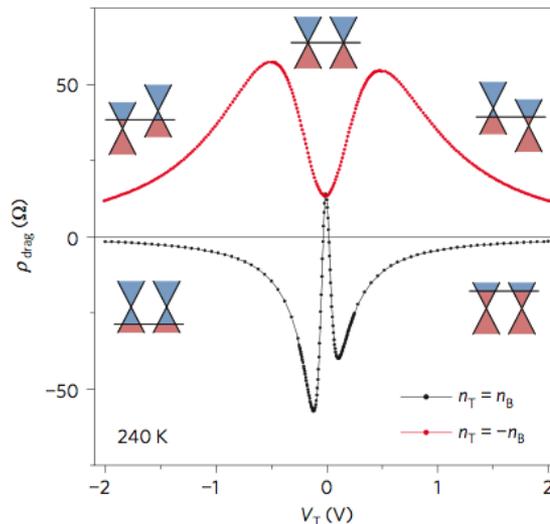
Coulomb Drag in Double Layer Graphene



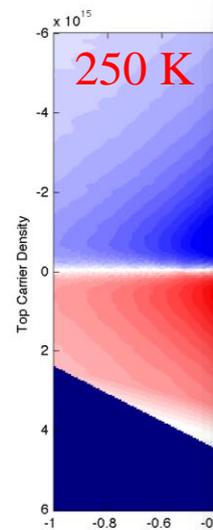
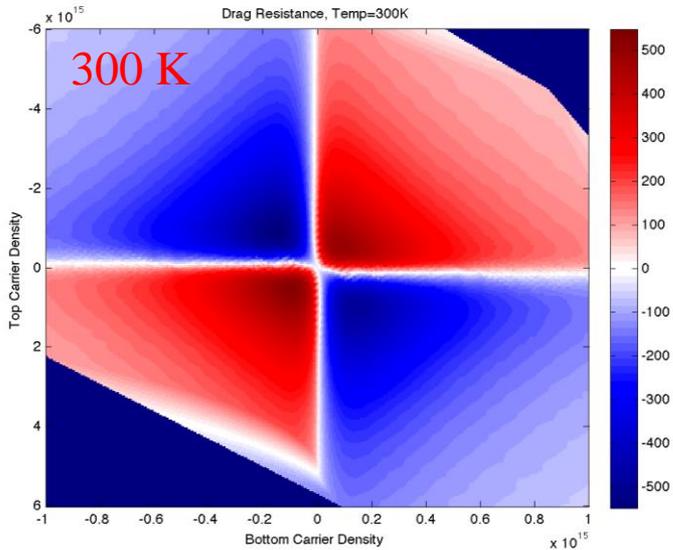
$T = 300$ K



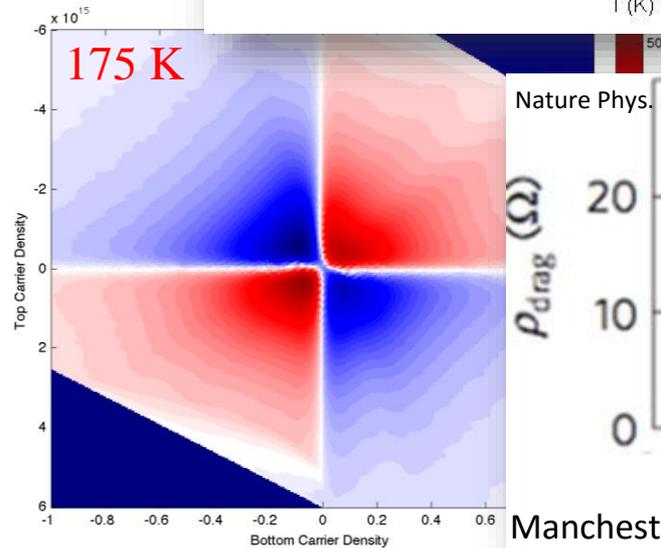
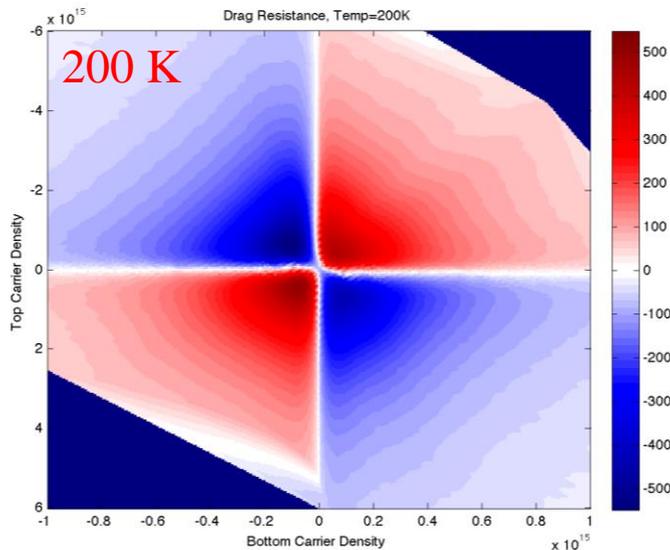
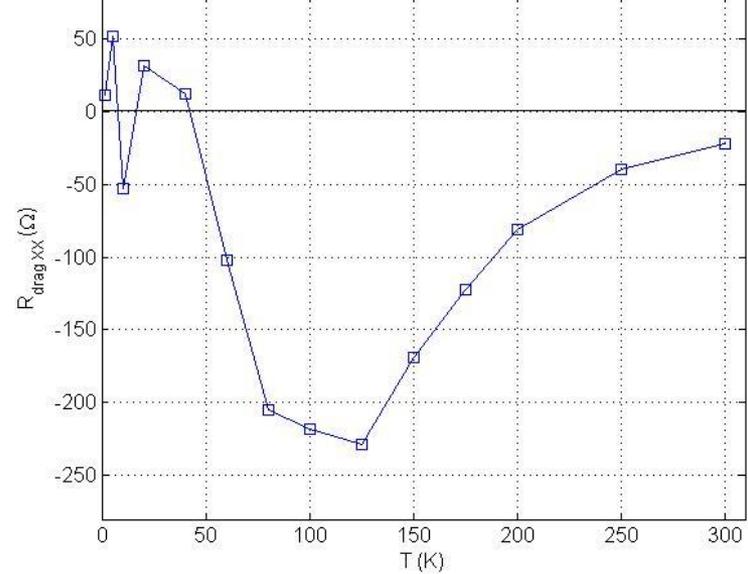
Drag resistance sign reversal at the double neutrality point



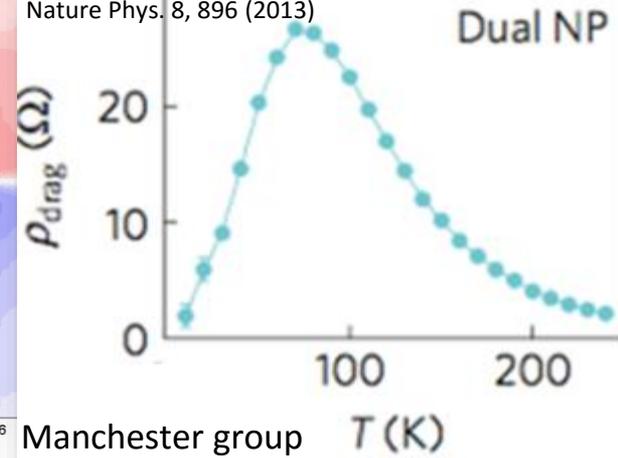
Temperature Dependent Coulomb Drag



Drag Resistance at Double Neutrality



Nature Phys. 8, 896 (2013)



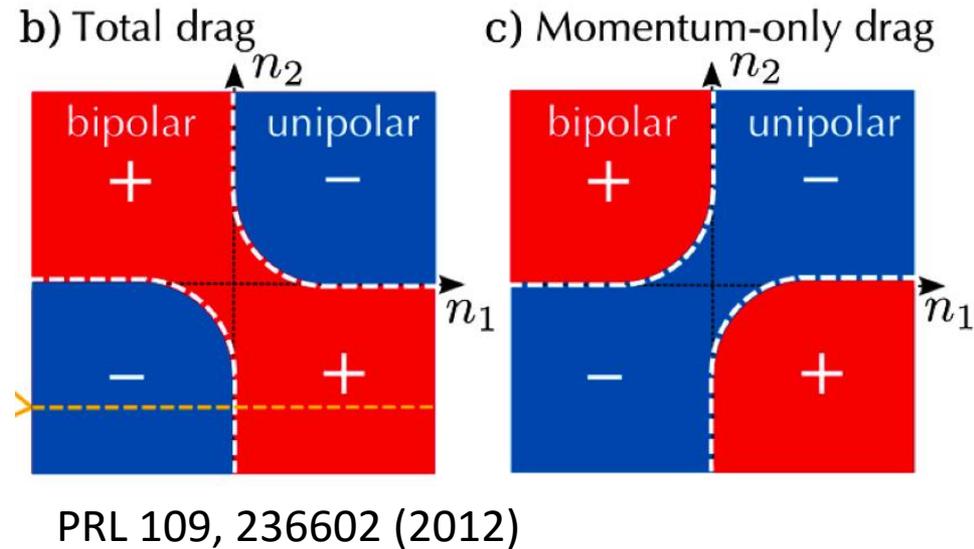
Possible Scenarios

- Interaction dominates at DNP
(Schutt et al, PRL 110, 026601 (2012))

- Energy vs. momentum drag
(Song and Levitov, PRL 109, 236602 (2012))

-correlated disorder

-anti-correlated disorders



- *Strain induced*

(Gibertini et al., Phys. Rev. B 85, 201405 (2012))

Sign of Drag Near the Neutrality

Selected for a Viewpoint in *Physics*
 PHYSICAL REVIEW LETTERS

week ending
 22 JULY 2016

Negative Coulomb Drag in Double Bilayer Graphene

J. I. A. Li,¹ T. Taniguchi,² K. Watanabe,² J. Hone,³ A. Levchenko,⁴ and C. R. Dean¹

PHYSICAL REVIEW LETTERS

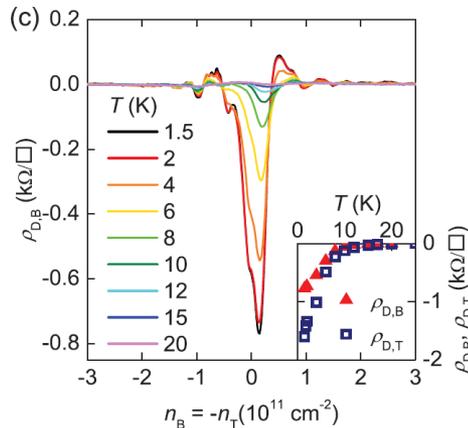
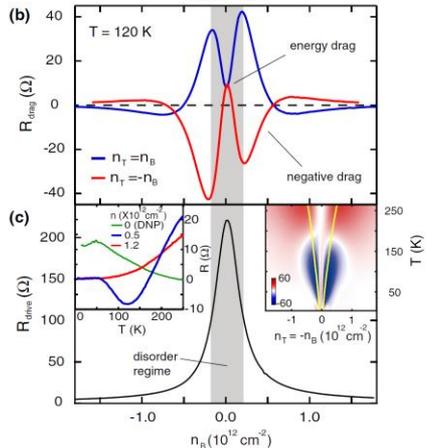
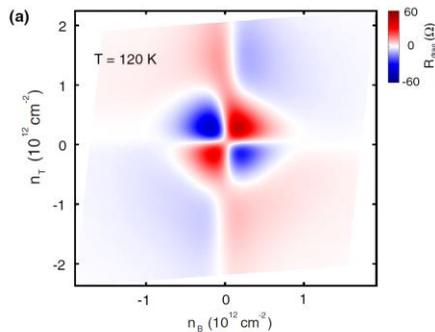
WEEK ENDING
 22 JULY 2016

Giant Frictional Drag in Double Bilayer Graphene Heterostructures

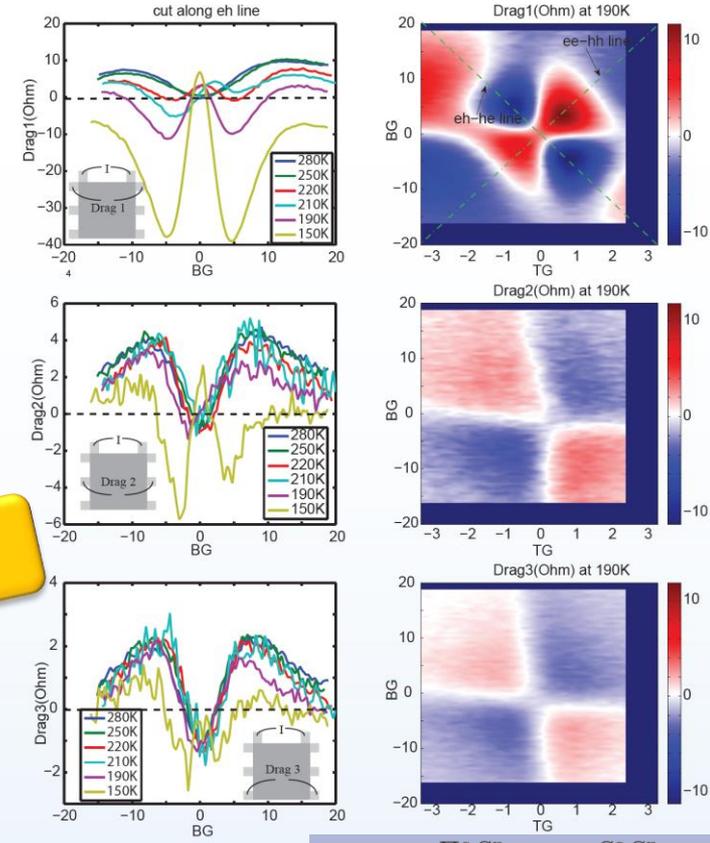
Kayoung Lee,¹ Jiamin Xue,¹ David C. Dillen,¹ Kenji Watanabe,² Takashi Taniguchi,² and Emanuel Tutuc¹

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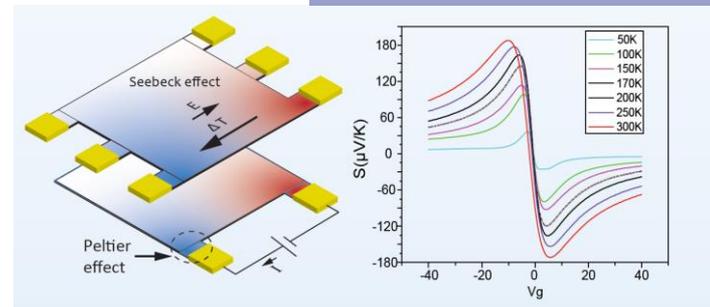
WEEK ENDING
 22 JULY 2016



Thermoelectric origin?

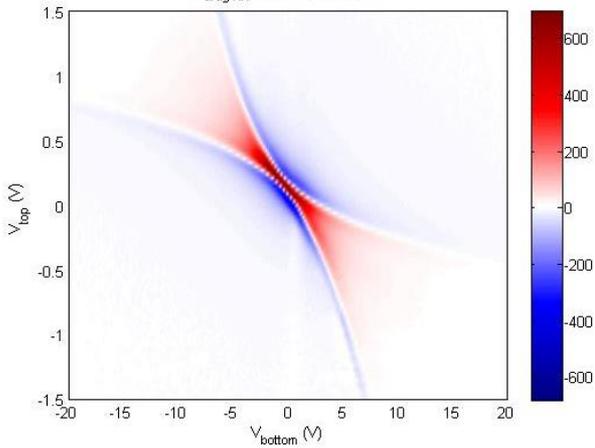


$$R_{drag} = \frac{\Pi^a S^p}{\kappa^a + \kappa^p} = \frac{S^a S^p}{L(\sigma^a + \sigma^b)}$$

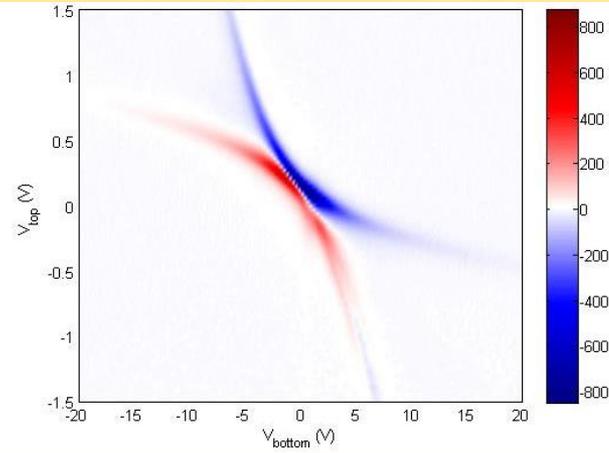


Magneto Drag in Graphene

$R_{\text{drag}}^{xx} (\Omega)$

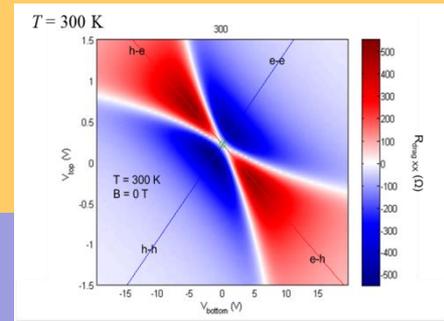


$R_{\text{drag}}^{xy} (\Omega)$

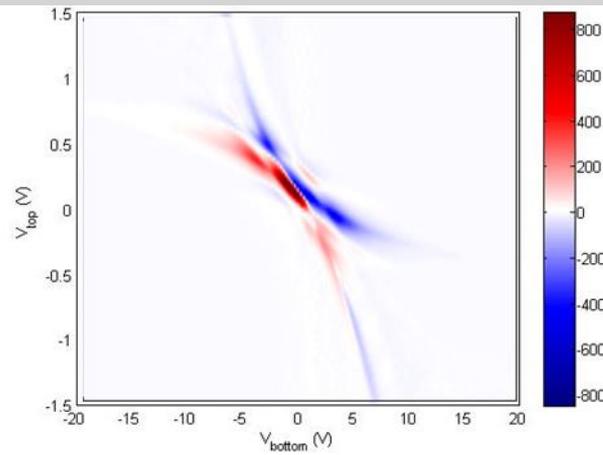
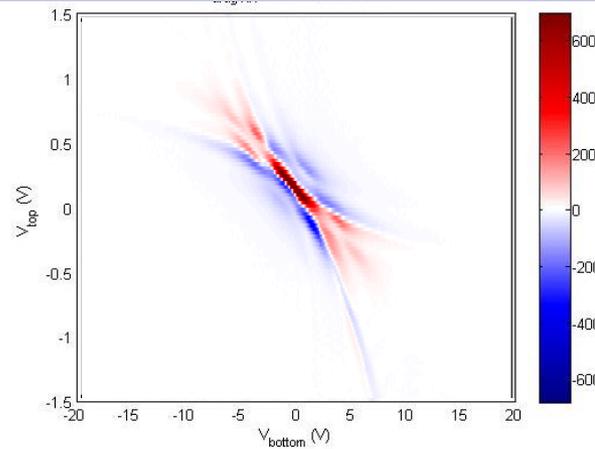


$B = 1 \text{ T}$

$T = 150 \text{ K}$



$T = 70 \text{ K}$



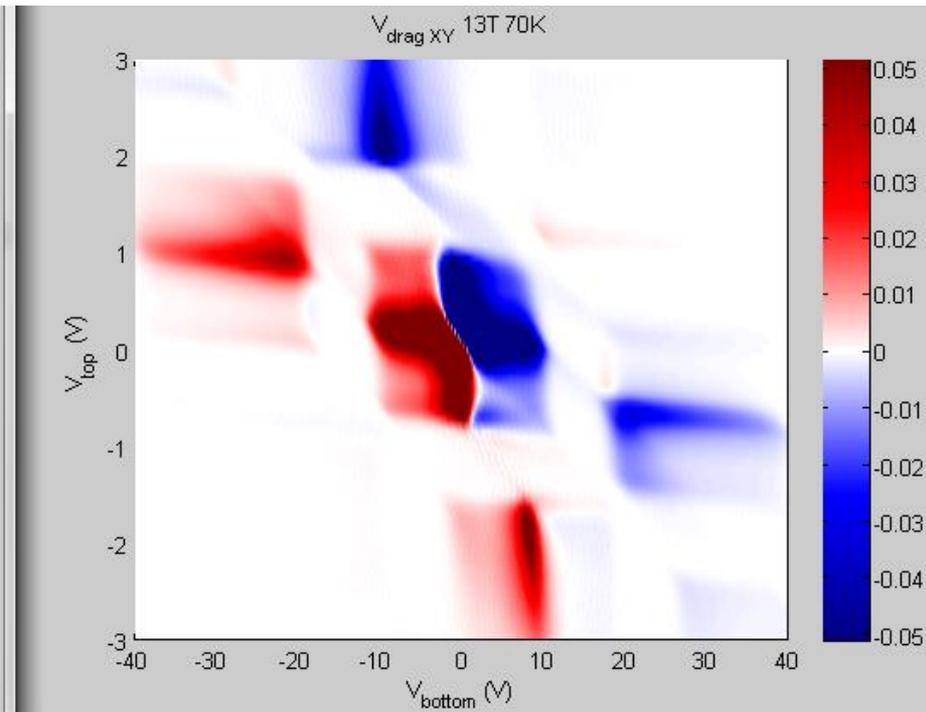
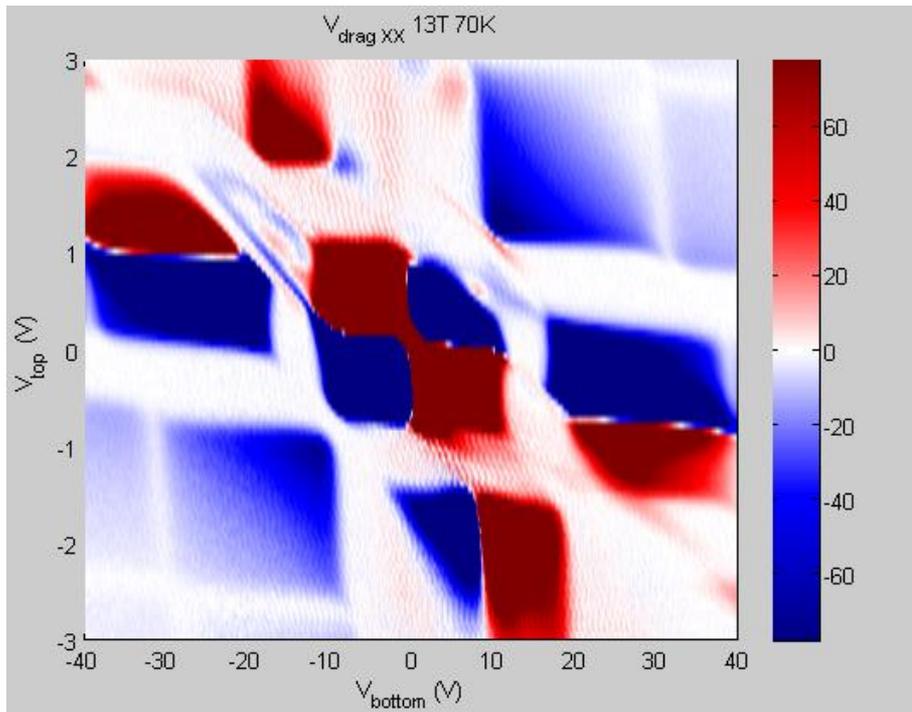
Magneto Drag in Graphene at High Magnetic Field

$B = 13 \text{ T}$

$T = 70 \text{ K}$

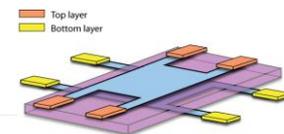
$R_{\text{drag}}^{xx} (\Omega)$

$R_{\text{drag}}^{xy} (h/e^2)$



**Finite drag signal only when the LLs are partially filled.
No drag if either layer becomes incompressible**

Correlation Between Drag and Transport



Oscillating Sign of Drag in High Landau Level

Felix von Oppen,¹ Steven H. Simon,² Ady Stern³

¹Institut für Theoretische Physik, Freie Universität Berlin, Arnimallee 14, 14

²Lucent Technologies, Bell Labs, Murray Hill, New Jersey, 0

³Department of Condensed Matter Physics, The Weizmann Institute of Science

(Received 30 April 2001; published 20 August 2001)

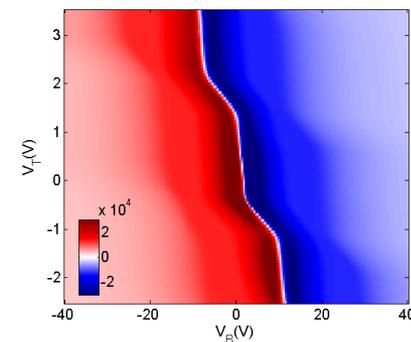
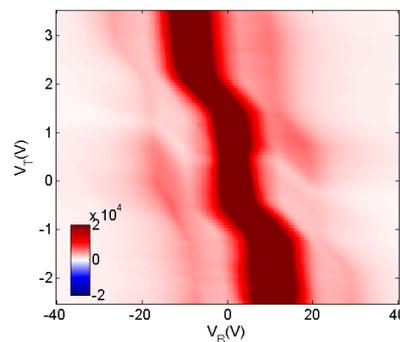
$$\rho_{xx}^D \sim \rho_{xy}^p \left\{ \frac{d\sigma_{yy}^p}{d(en)} \frac{d\sigma_{yy}^a}{d(en)} + \frac{d\sigma_{yx}^p}{d(en)} \frac{d\sigma_{xy}^a}{d(en)} \right\} \rho_{yx}^a$$

$$\rho_{xy}^D \sim \rho_{xy}^p \left\{ \frac{d\sigma_{yy}^p}{d(en)} \frac{d\sigma_{yx}^a}{d(en)} + \frac{d\sigma_{yx}^p}{d(en)} \frac{d\sigma_{xx}^a}{d(en)} \right\} \rho_{xy}^a$$

Measured Magneto and Hall Resistance

$\rho_{xx}^{bottom} (\Omega)$

$\rho_{xy}^{bottom} (\Omega)$

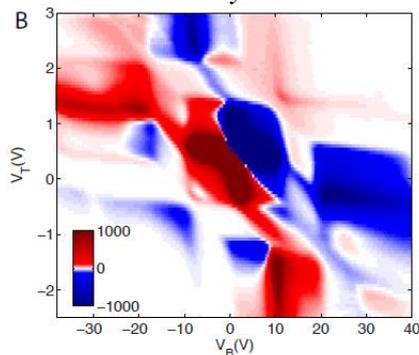
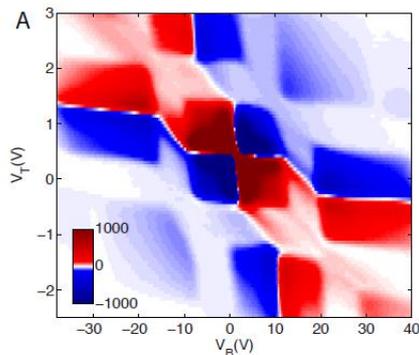


Feng et al. DDT (1000)

Measured magneto and Hall drag

$\rho_{xx}^D (\Omega)$

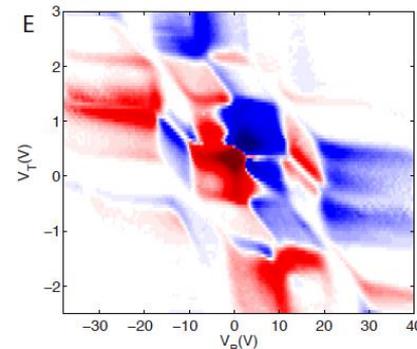
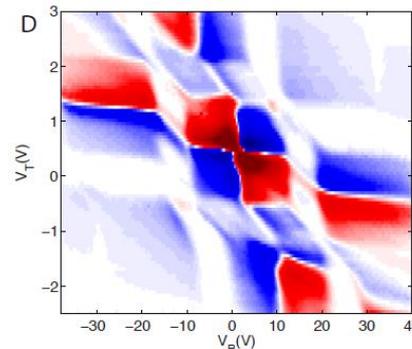
$\rho_{xy}^D (\Omega)$



Calculated magneto and Hall drag

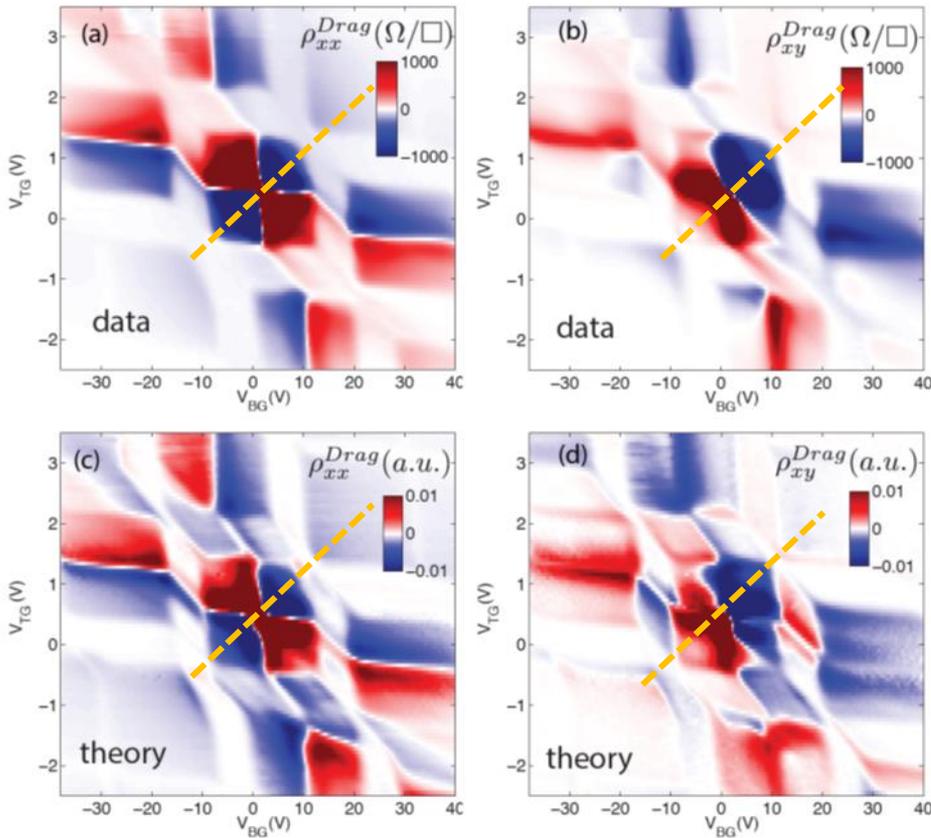
$\rho_{xx}^D (\Omega)$

$\rho_{xy}^D (\Omega)$

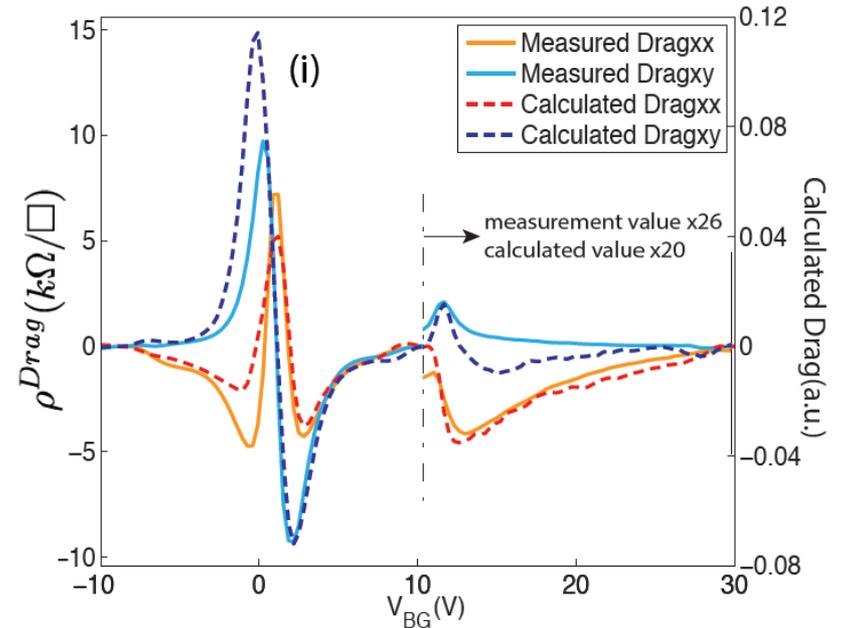


Quantitative Comparison

OSS PRL (2001)



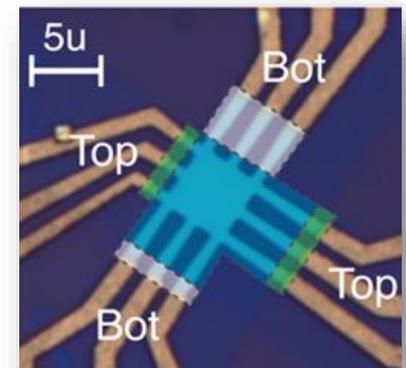
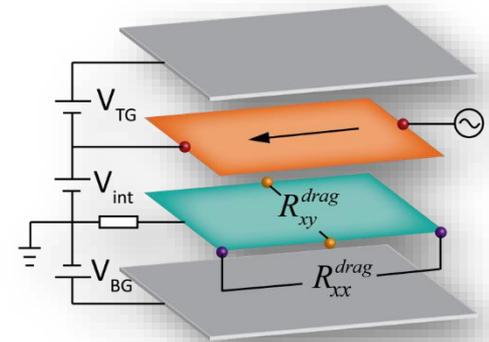
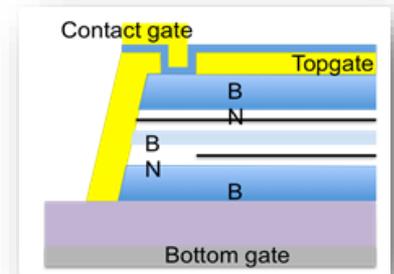
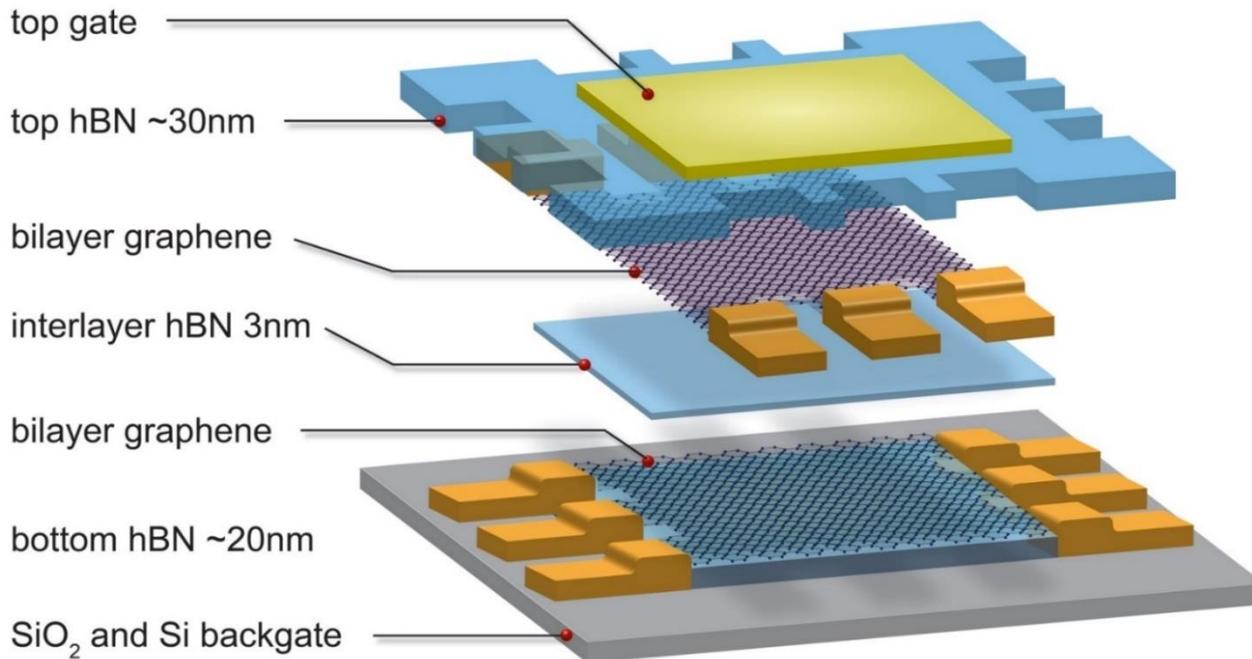
$$\hat{\rho}^{Drag} \sim -\hat{\rho}^p \frac{d\hat{\sigma}^p}{d(en^p)} \frac{d\hat{\sigma}^a}{d(en^a)} \hat{\rho}^a$$



Components of drag tensor follows a similar scaling prefactor.

Double Bilayer Graphene Drag Device

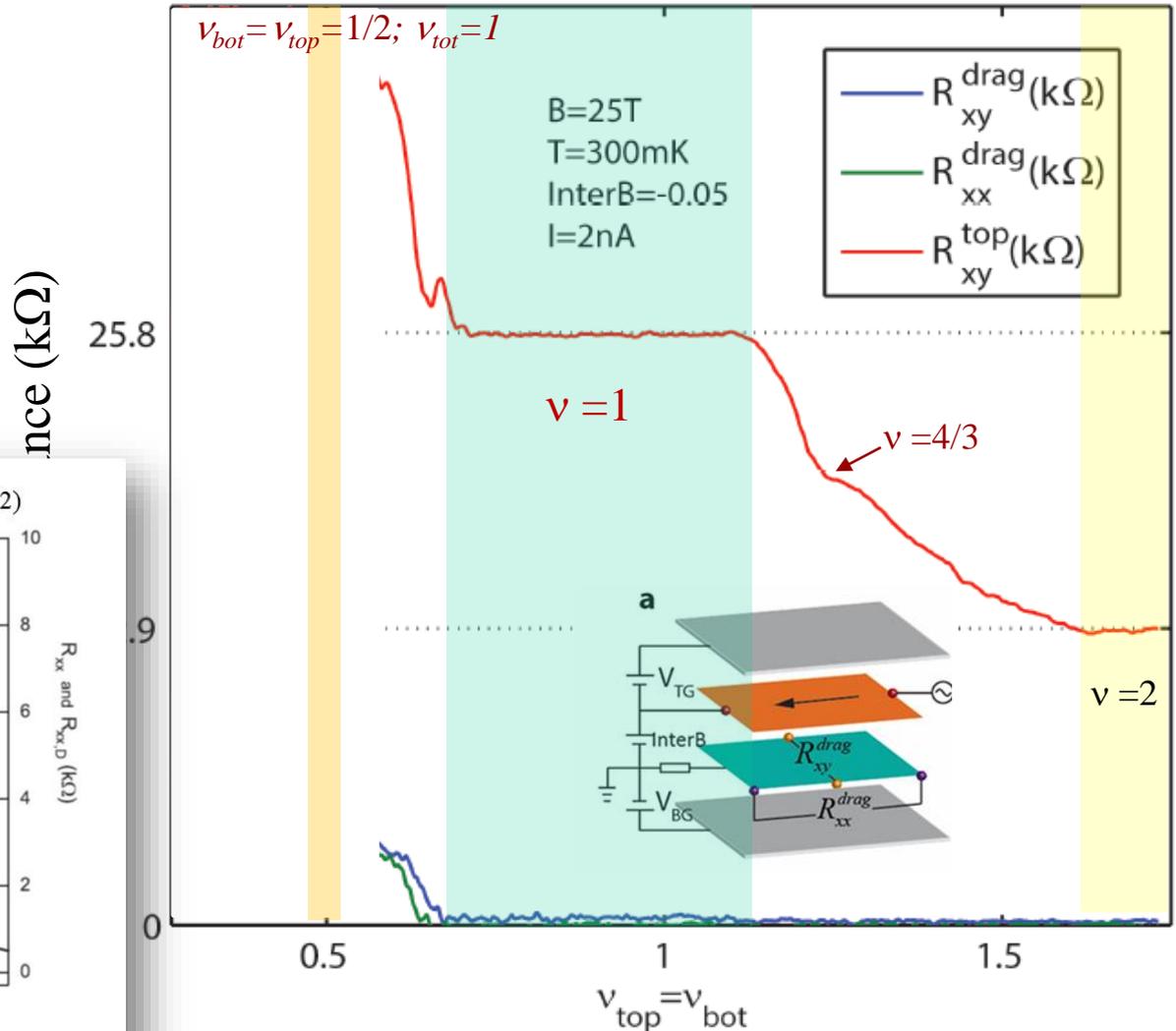
- Mobility $\sim 10^6$ cm²/Vsec
- hBN thickness $d = 3$ nm
- top and bottom gate
- contact gate
- interlayer bias



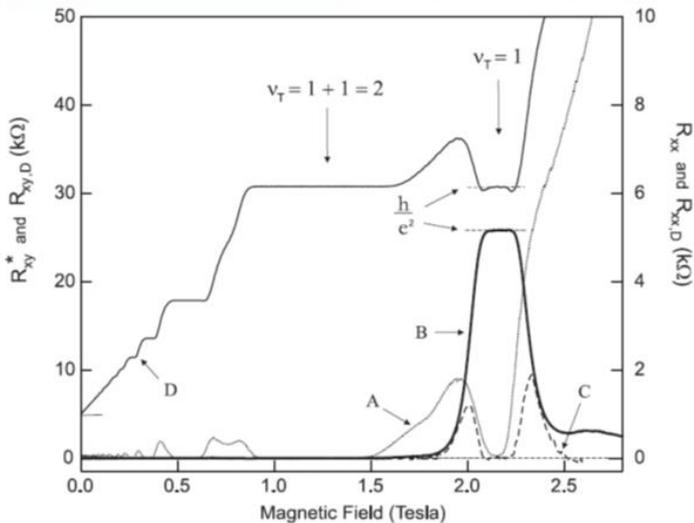
Bilayer Graphene/hBN/Bilayer Graphene: Quantized Hall Drag

- Mobility $\sim 10^6$ cm²/Vsec
- hBN thickness $d = 3$ nm
- top and bottom gate
- contact gate
- interlayer bias

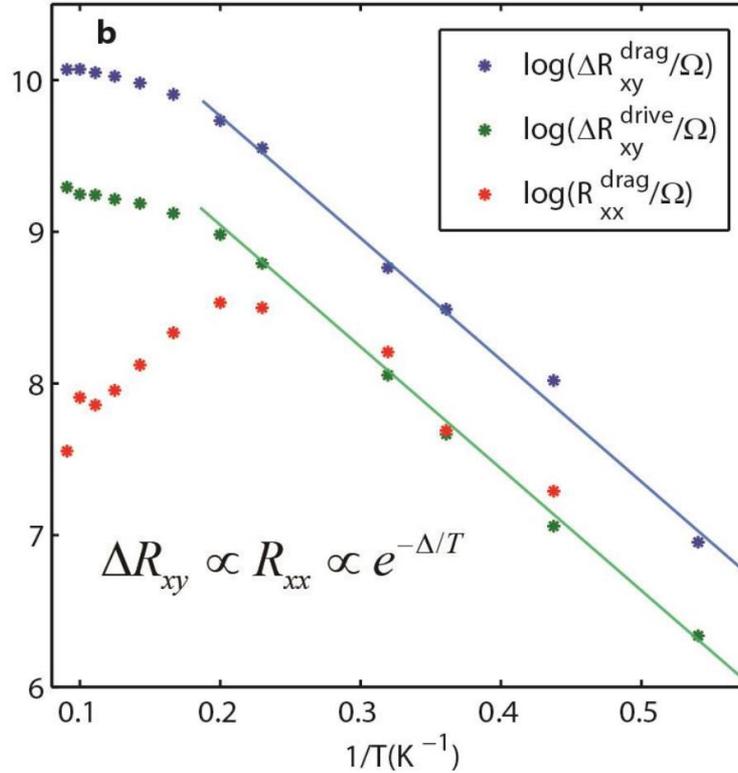
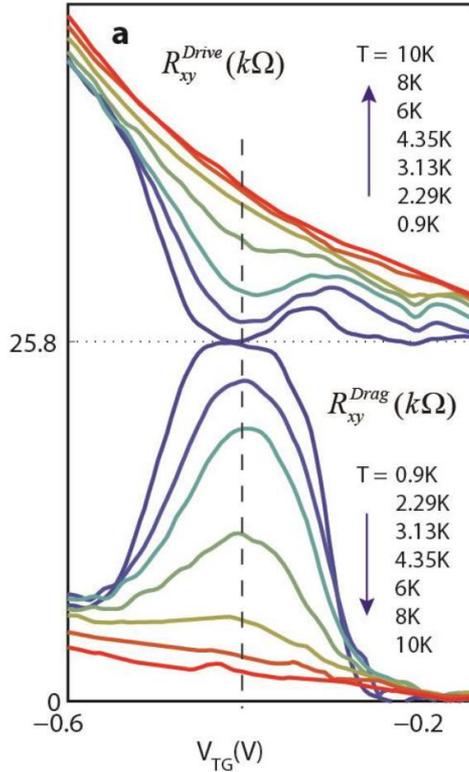
Hall Resistance, Magneto Drag, and Hall Drag



GaAs Double Quantum Well: Kellogg et al. PRL (2002)

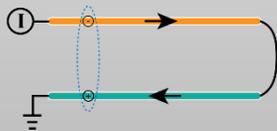


Exciton BEC Energy Scale and Counter Flow



$\Delta \sim 8 \text{ K.}$

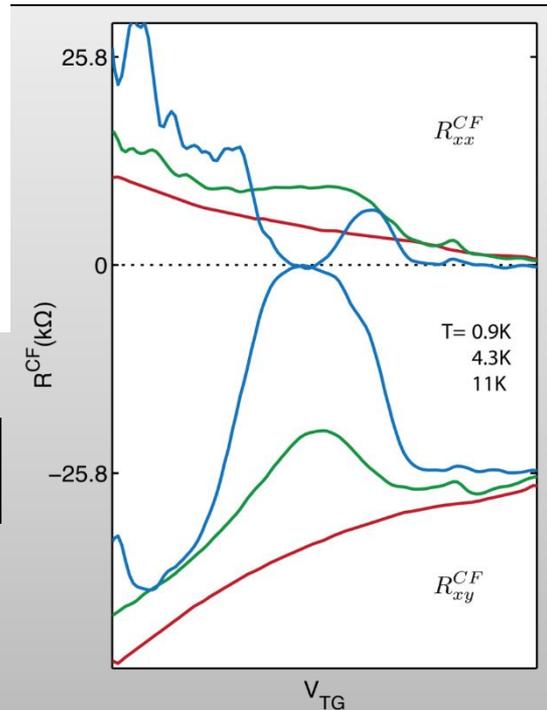
Superfluidic Counter flow



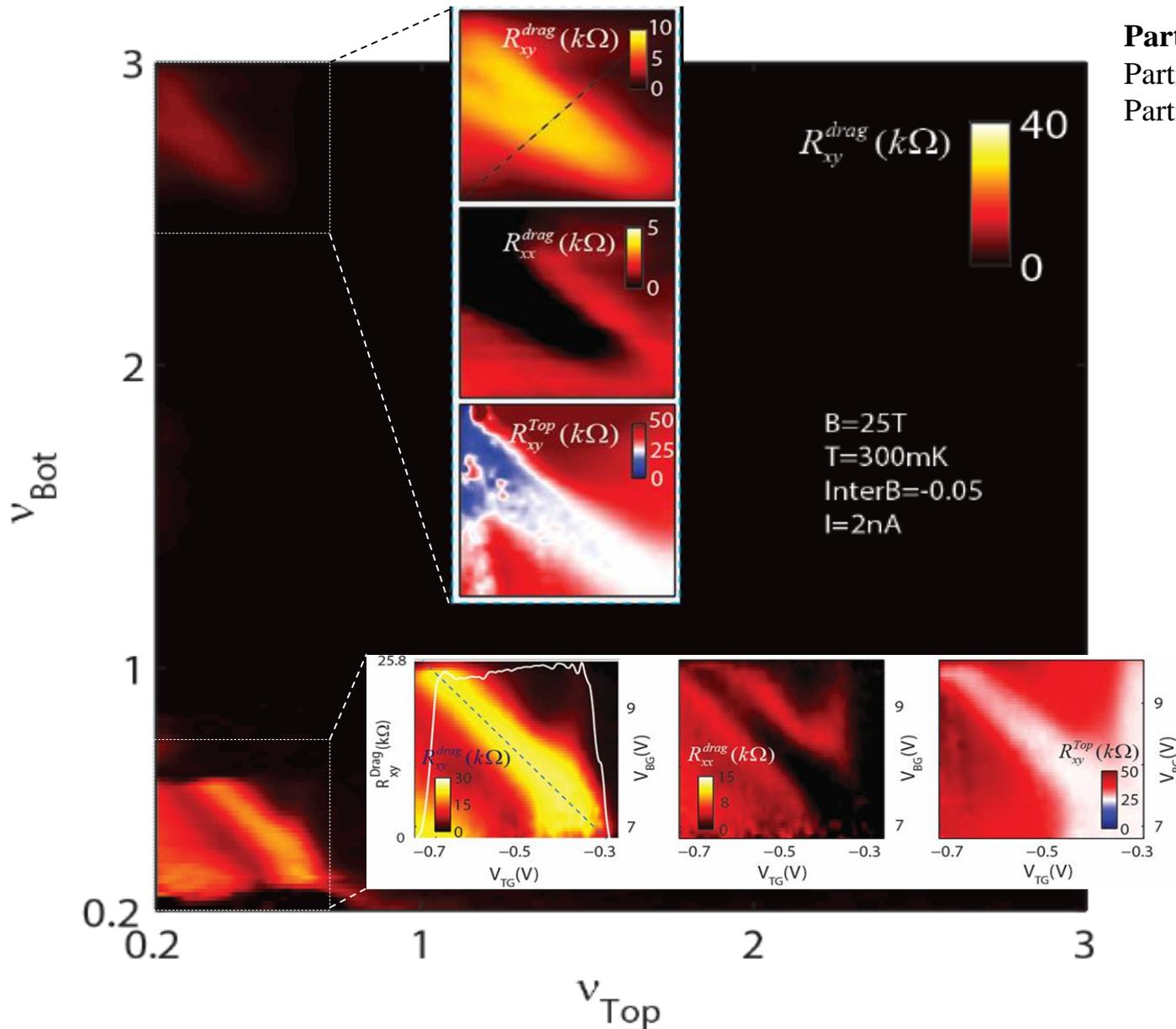
$$\begin{bmatrix} V_1 \\ V_2 \end{bmatrix} = \begin{bmatrix} R_{11} & R_{12} \\ R_{21} & R_{22} \end{bmatrix} \times \begin{bmatrix} I_1 \\ I_2 \end{bmatrix}$$

$$V_{top}^{CF} = R_{top}I - R_{drag}I$$

$$V_{bot}^{CF} = R_{drag}I - R_{bot}I$$

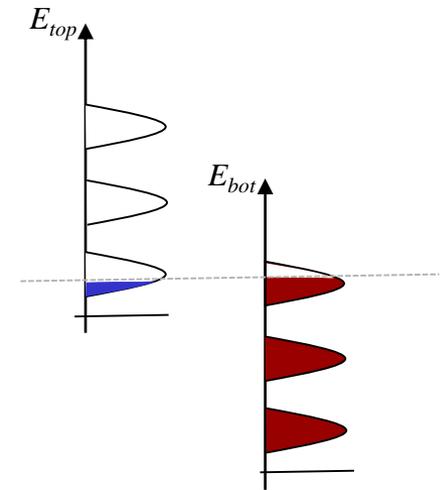


Quantized Hall Drag for $\nu_{tot} = 1$ and 3



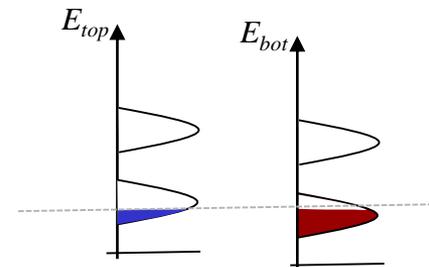
Partial coherent exciton current:

Partially filled $N_{top}=1$
 Partially filled $N_{bot}=3$

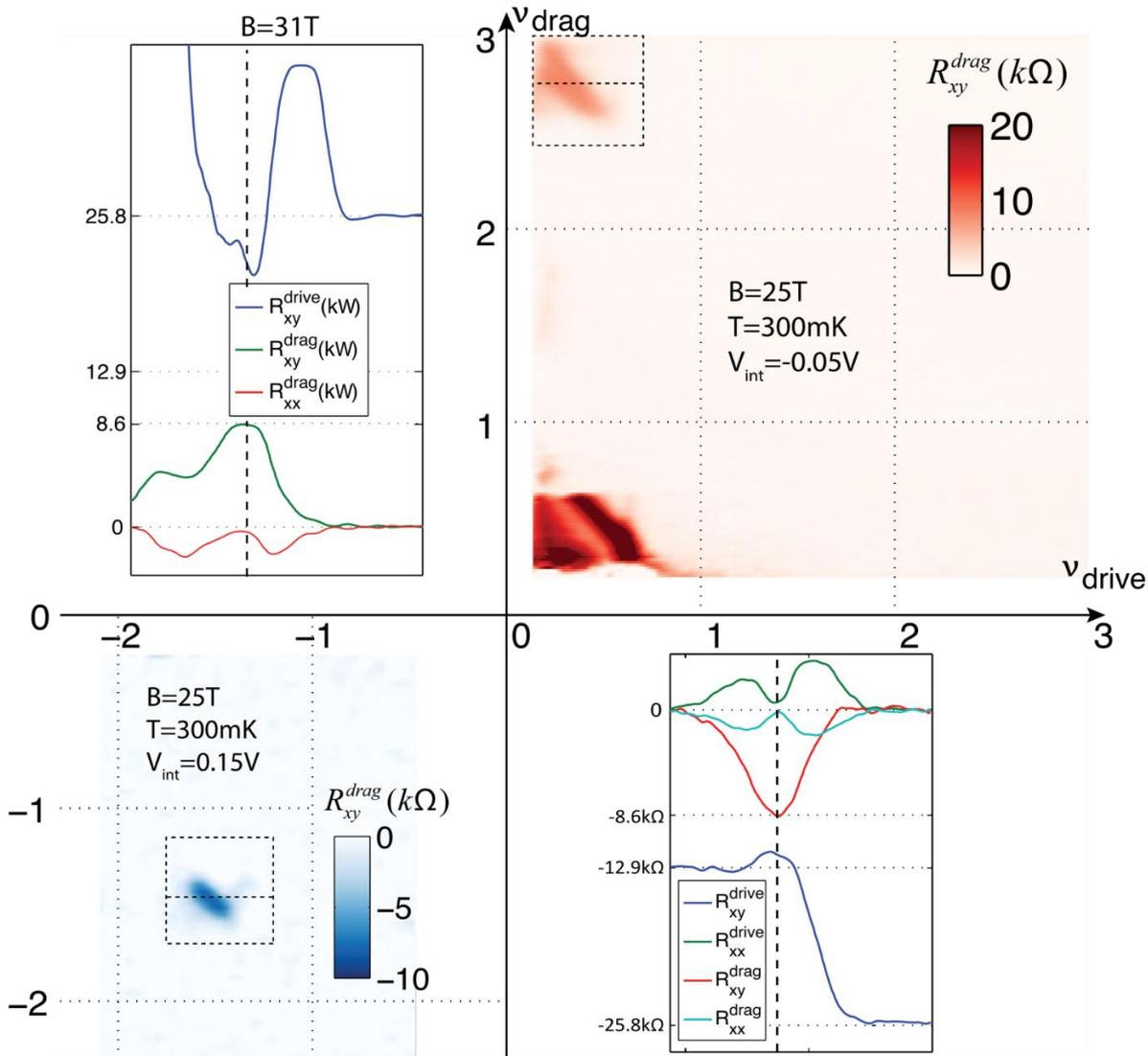


Coherent exciton current:

Partially filled $N_{top}=1$
 Partially filled $N_{bot}=1$



Magneto Exciton Condensation in Different LLs



Observed Exciton condensations

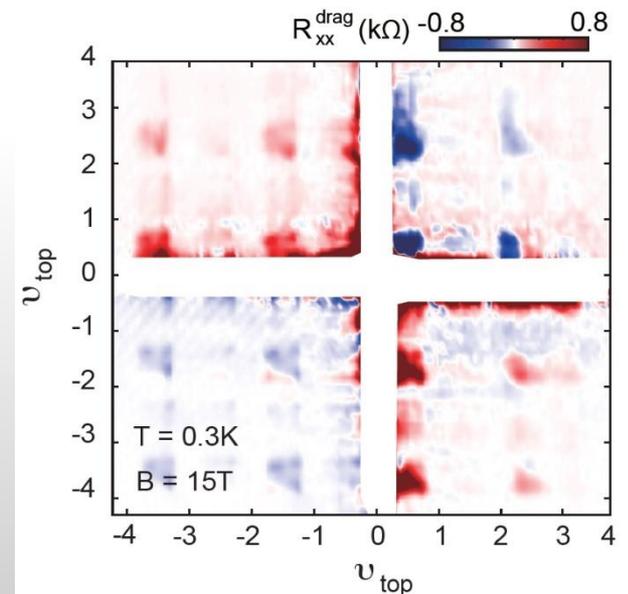
(0.5, 0.5)

(0.5, 2.5); (2.5, 0.5)

(-0.5, -0.5)

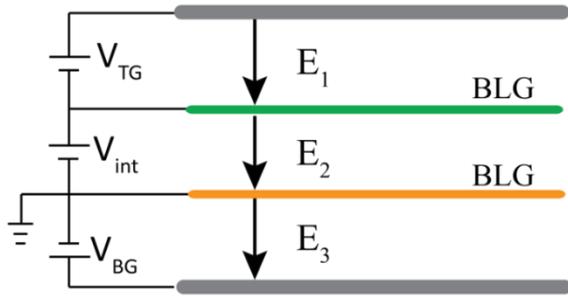
Possible $\nu \rightarrow \nu + 2$ symmetry
in bilayer graphene double layer

Other possible Exciton Condensation



Li et. al., arXiv:1608.05846

Effect of Interlayer Bias Voltage



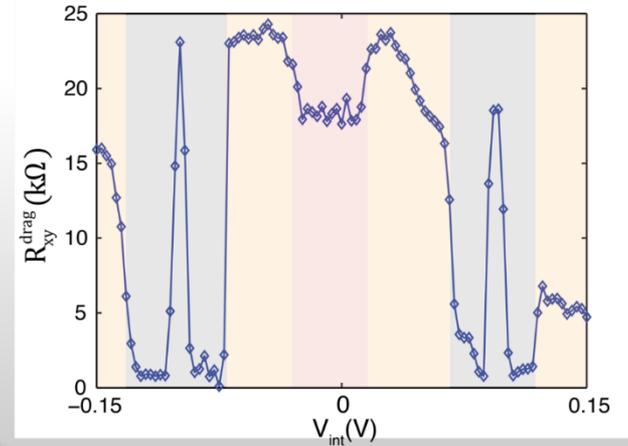
$$n_{top} = C_{TG}V_{TG} - C_{int}V_{int}$$

$$n_{top} = C_{BG}V_{BG} + C_{int}V_{int}$$

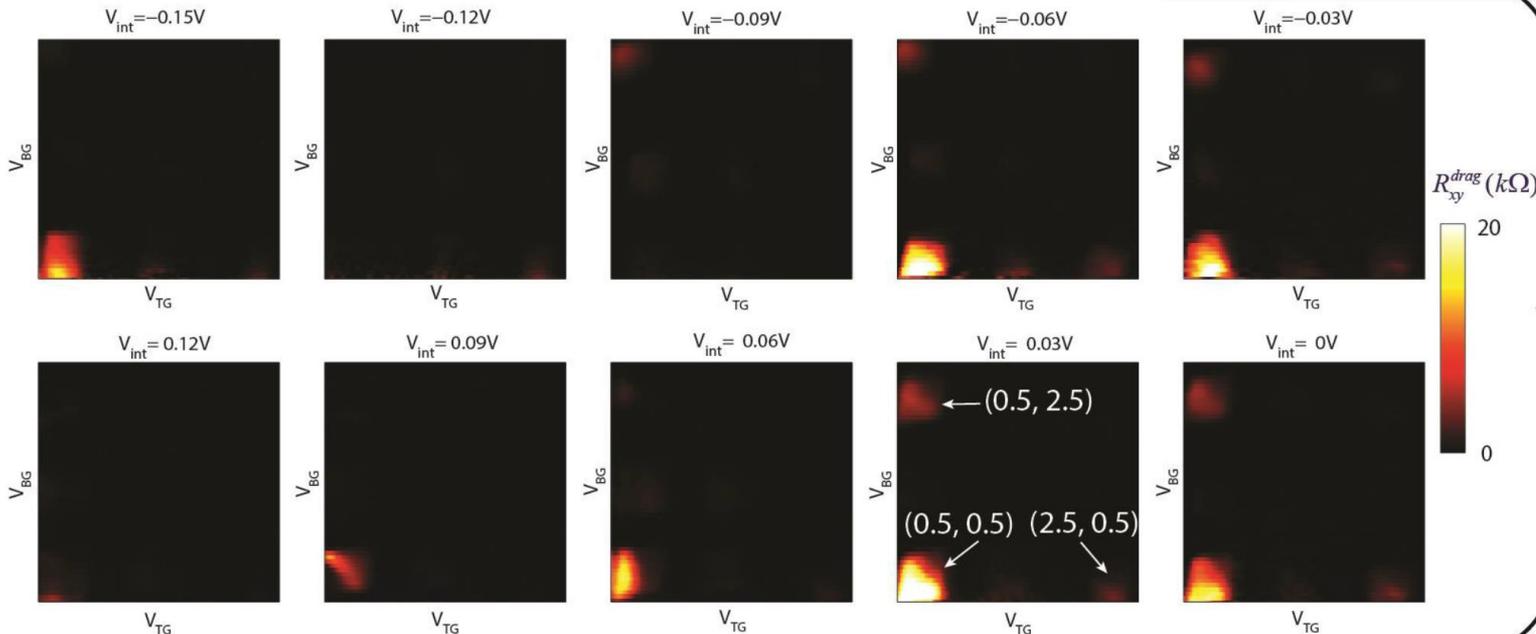
$$D_{top} = (C_{int}V_{int} + n_{top})/2$$

$$D_{bot} = (C_{int}V_{int} - n_{bot})/2$$

Interlayer Bias voltage dependence (0.5, 0.5)

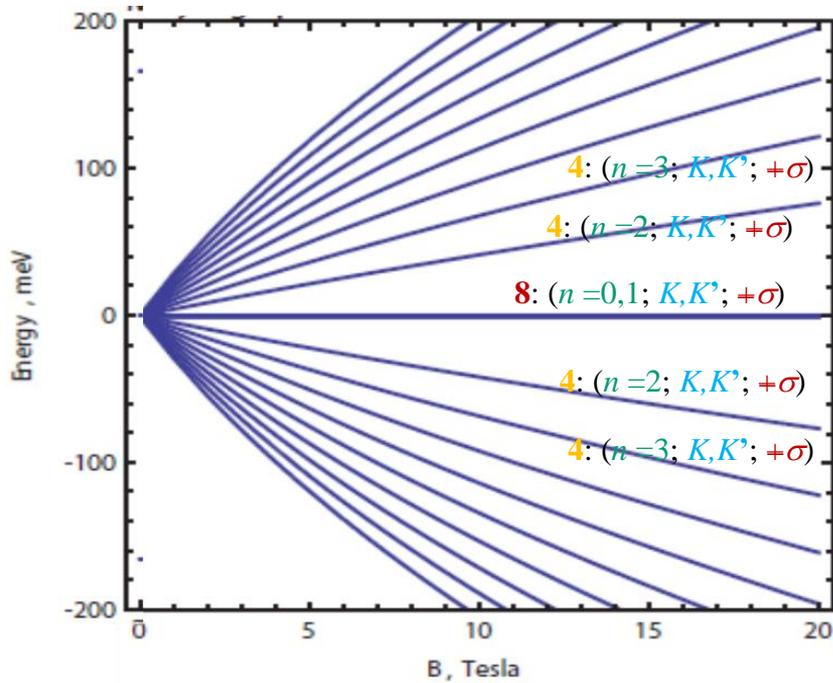


Interlayer Bias voltage dependence R_{xy}^{drag}



Quantum Hall Ferromagnetic Phase Transition in Bilayer Graphene

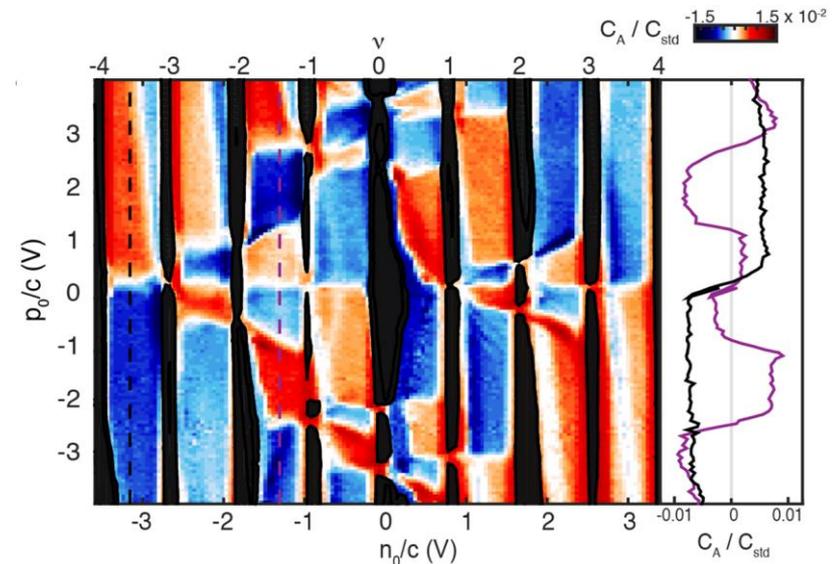
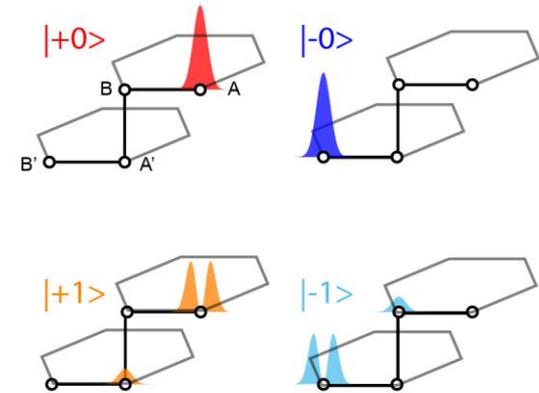
Bilayer Landau level spectrum: SU(4) and SU(8)



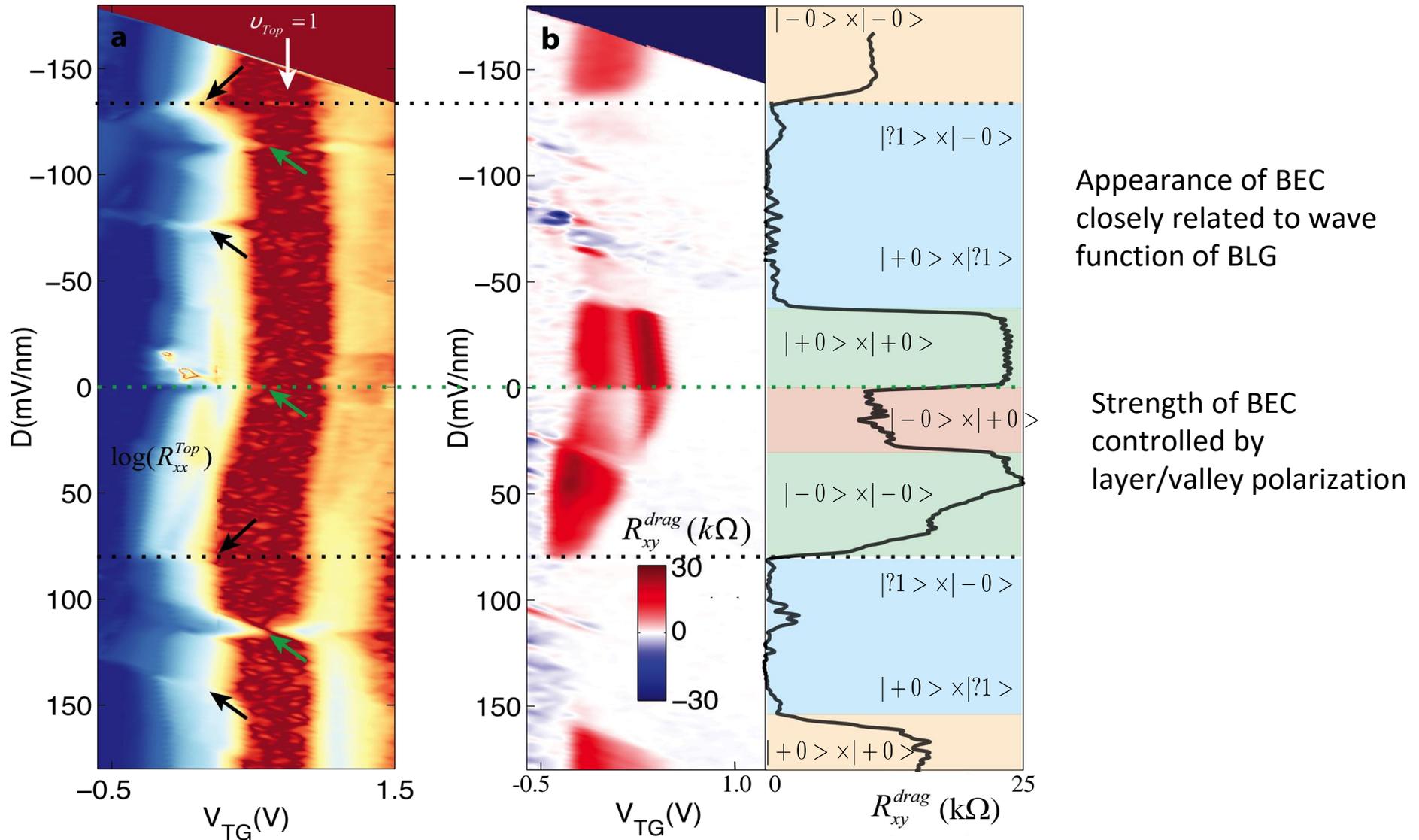
Broken Symmetry Gap in Bilayer due to Interaction:
Tuned by displacement field (pseudo magnetic field)

B. M. Hunt etc. (2016).

Each Landau level is degenerate for spin and valley except zero energy LL where there is an additional 'accidental' degeneracy $n=0,1$.



Exciton BEC Phase Transition: Internal Degree of Freedom of Exciton



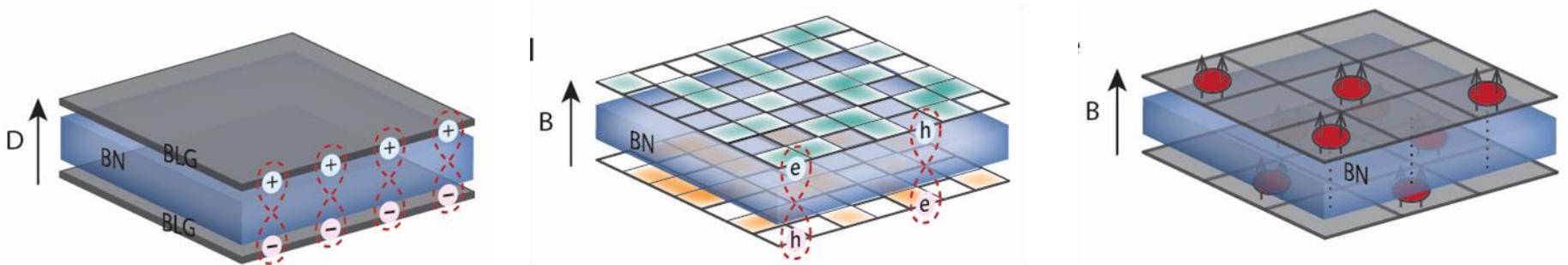
Internal degree of freedom of excitons can be controlled and incur phase transitions in the BEC.

Summary and Outlook

- Semiclassical Coulomb drag in the presence of thermal fluctuations
- Quantum drag Hall effect and robust magneto exciton condensation
- Exciton condensation between different Landau levels
- Phase transition between different exciton condensations

Moving forward:

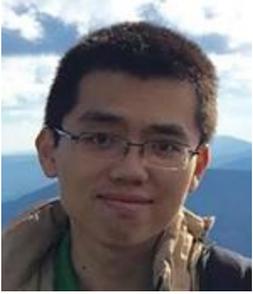
Double monolayer, resonance tunneling, phase transition in BEC, exciton insulator ($B=0$), fractional excitons,



Acknowledgements

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Bert Halperin

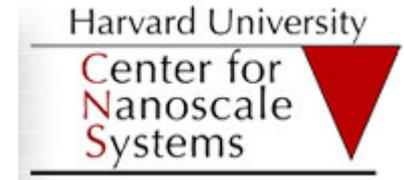


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Kenji Watanabe



CNS
(Fabrication facility)



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