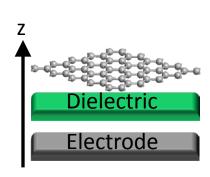
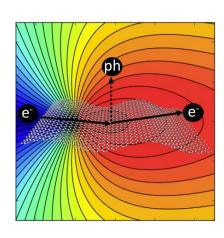






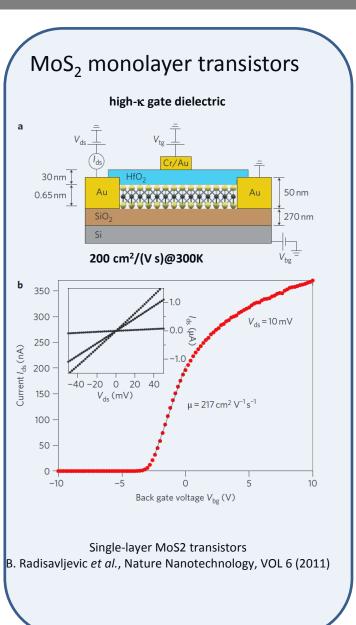
# Flexural-phonon scattering induced by electrostatic gating in graphene



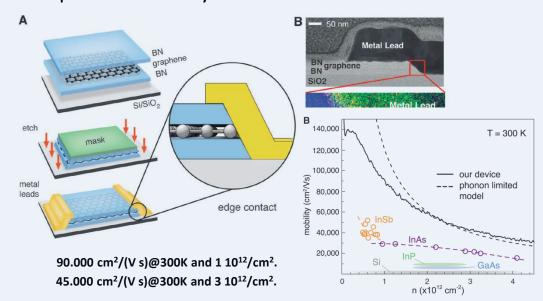


Graphene device – a vibrating membrane in an electrostatic and dielectric environment. Electron-phonon coupling and mobility limitations from the environment.

# Motivation – novel 2D materials and el-ph scattering

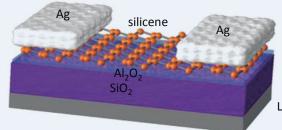


#### Graphene monolayer transistors



One-Dimensional Electrical Contact to a Two-Dimensional Material L. Wang et al., Science, VOL 342 (2013) Encapsulation, edge contacting suspend in high-κ liquid-> screen charged impurities

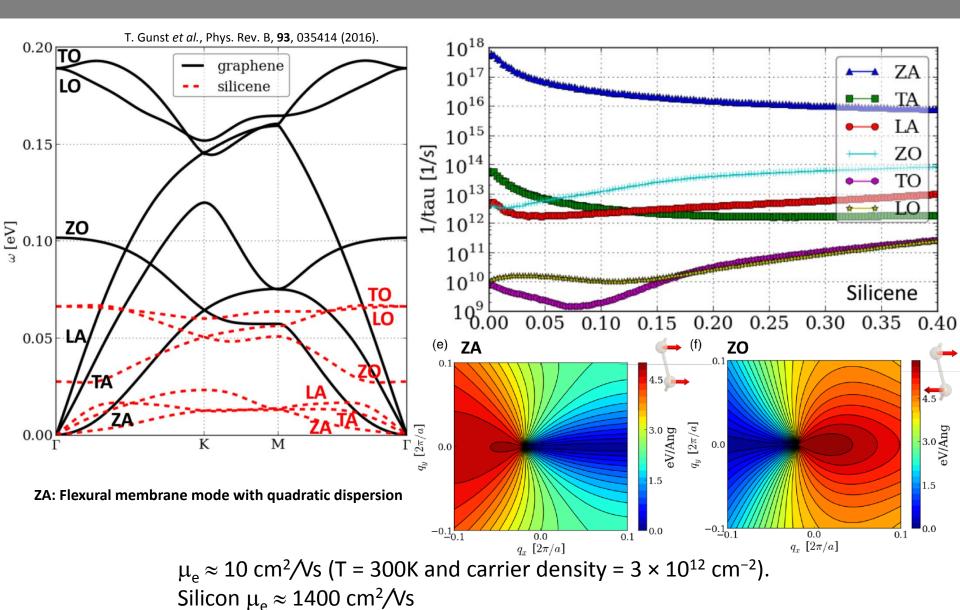
#### Silicene monolayer transistors



Grown at Ag+flip+pattern 100 cm<sup>2</sup>/(V s)@300K

Silicene field-effect transistors operating at room temperature Li Tao et al., Nature Nanotechnology, VOL 10 (2015)

#### Motivation - Silicene (a 2D material with broken planar mirror symmetry)

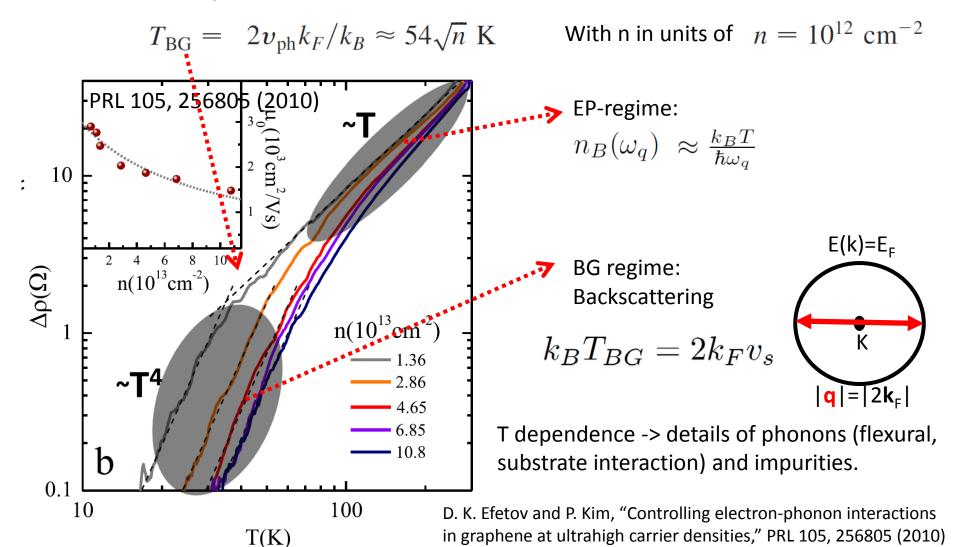


Putting graphene in a gated device stack can also break the planar mirror symmetry!

#### Temperature regimes

T << T<sub>BG</sub>: phonon system is degenerate [Bloch-Grüneisen (BG)]

 $T \gg T_{BG}$ : phonon system is nondegenerate [equipartition (EP)].



# **Summary of semiclassical BTE method**

Phonon scattering rates (from k to k') – absorption/emission of a phonon (FGR):

$$P_{\mathbf{k}\mathbf{k}'}^{\lambda nn'} = \frac{2\pi}{\hbar} g_{\mathbf{k}\mathbf{k}'}^{\lambda nn'} \left[ n_{\mathbf{q}}^{\lambda} \delta \left( \epsilon_{\mathbf{k}'n'} - \epsilon_{\mathbf{k}n} - \hbar \omega_{q\lambda} \right) \delta_{\mathbf{k}',\mathbf{k}+\mathbf{q}} + \left( n_{-\mathbf{q}}^{\lambda} + 1 \right) \delta \left( \epsilon_{\mathbf{k}'n'} - \epsilon_{\mathbf{k}n} + \hbar \omega_{-q\lambda} \right) \delta_{\mathbf{k}',\mathbf{k}-\mathbf{q}} \right]$$

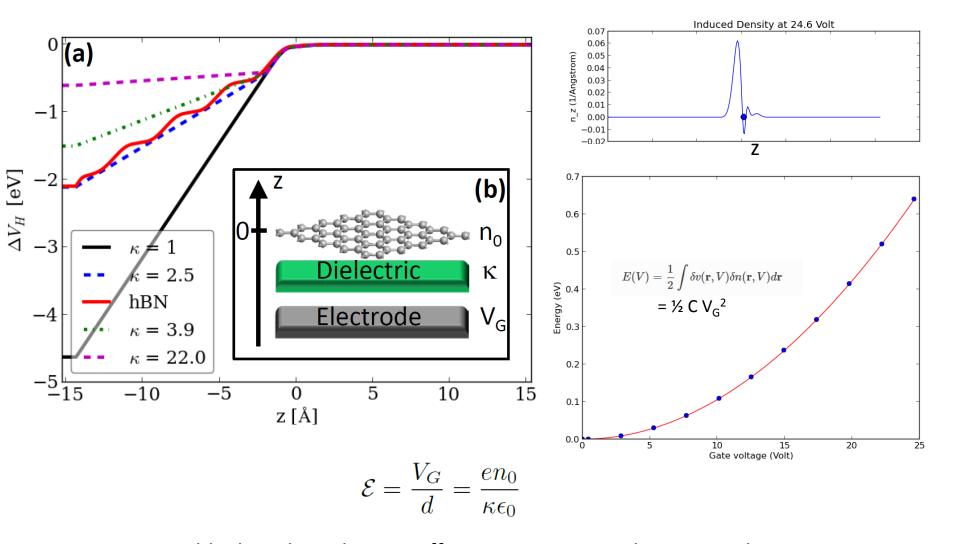
Transport relaxation time:

$$\frac{1}{\tau_{\mathbf{k}n}} = \sum_{\mathbf{k}'n'} \frac{(1 - f_{\mathbf{k}'n'}^0)}{(1 - f_{\mathbf{k}n}^0)} \left(1 - \cos(\theta_{\mathbf{k}\mathbf{k}'})\right) P_{\mathbf{k}\mathbf{k}'}^{nn'} \quad , \quad \cos(\theta_{\mathbf{k}\mathbf{k}'}) = \frac{\mathbf{v}_{\mathbf{k}'n'} \cdot \mathbf{v}_{\mathbf{k}n}}{|\mathbf{v}_{\mathbf{k}'n'}||\mathbf{v}_{\mathbf{k}n}|}$$

Mobility: 
$$\mu_e = -2q \frac{\sum_{\mathbf{k}n \in c} |\mathbf{v}_{\mathbf{k}n}|^2 \frac{\partial f_{\mathbf{k}n}^0}{\partial \epsilon_{\mathbf{k}n}} \tau_{\mathbf{k}n}}{\sum_{\mathbf{k}n \in c} f_{\mathbf{k}n}^0}$$

T. Gunst, T. Markussen, K. Stokbro, M. Brandbyge, "First-principles method for electron-phonon coupling and electron mobility: Applications to 2D materials", Phys. Rev. B, **93**, 035414 (2016).

# A graphene capacitor

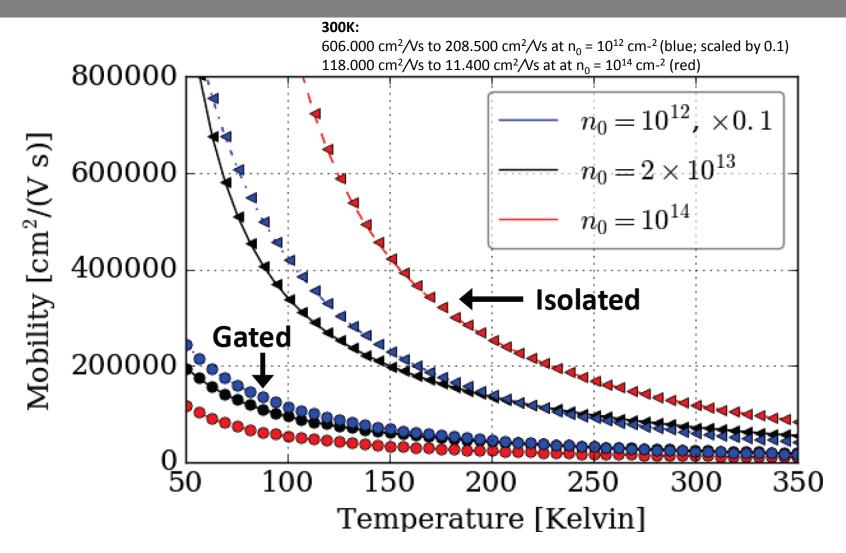


Highly doped graphene is efficient in screening the potential.

Behaves much like a metal-plate capacitor.

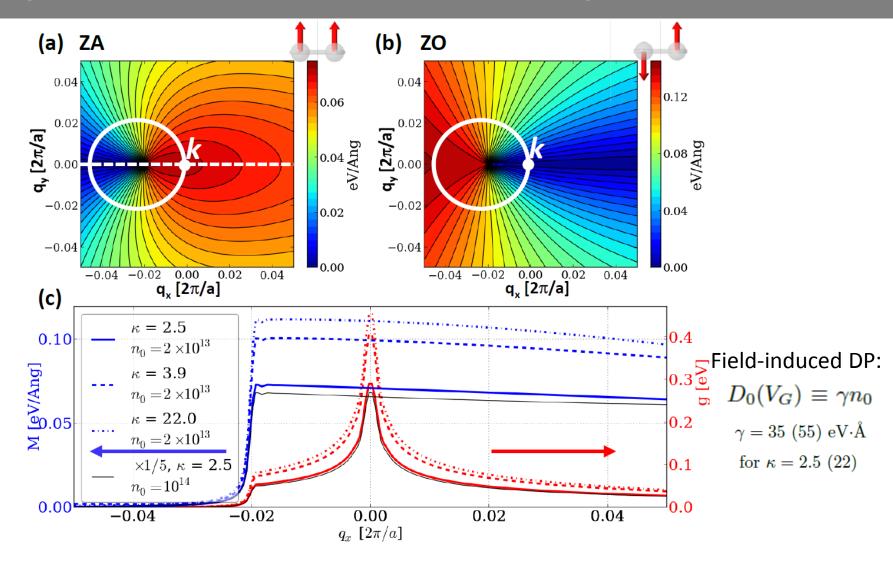
Gate field breaks planar mirror symmetry!

# **Mobility degradation**



Significant (field-induced) degradation!
Modified scaling with temperature and carrier density!

## Ingredients - Field-induced flexural scattering



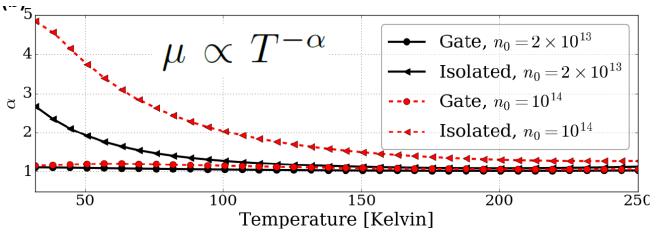
High occupation of (low frequency, constant DOS) ZA mode.

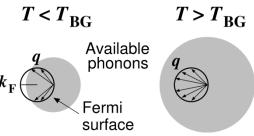
Nonzero coupling induced by the gate.

Coupling is linear in the induced carrier density and increases with dielectric constant.

## Temperature scaling

Dirac model: 
$$\frac{1}{\tau_k^{ZA}} = \frac{v_F D_0^2(V_G)}{2\pi \rho b^2} \varepsilon_k^{-2} \frac{k_B T}{q_c}$$





#### BG regime:

$$q_{\rm max} = 2k_F$$
  
 $k_B T_{\rm BG} = \hbar \omega_{q_{\rm max}}$   
 $T_{\rm BG}^{\rm LA} = 57 \sqrt{\tilde{n}} \, {\rm K}$   
 $T_{\rm BG}^{\rm ZA} = 0.46 \, \tilde{n} \, {\rm K}$   
 $\tilde{n} = n/10^{12} \, {\rm cm}^{-2}$ 

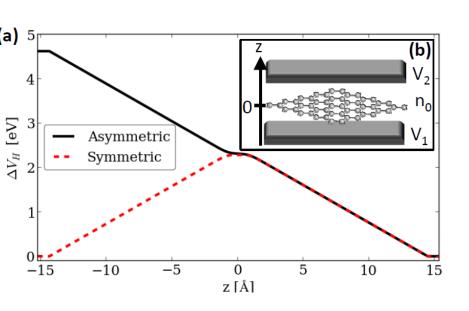
Dominant mechanism at room temperature and even more at low T!

Modified scaling with temperature!

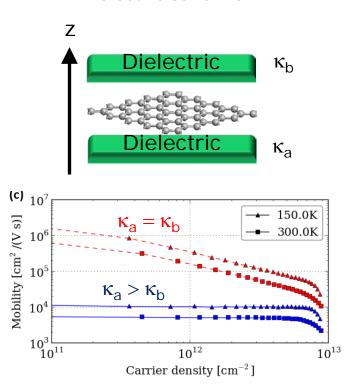
Room T: Indistinguishable from in-plane scattering (Higher DP).

## Electrode configuration – controlling the effect

#### Double gate stack



#### Dielectric sandwich



#### Breaking the symmetry or not?

Sandwiching graphene between two layers with different dielectric constants or electrodes at different potentials degrades mobility significantly.

#### **Experimental signatures**

Modified carrier density scaling:

$$\mu \sim n_0^{-3/2}$$
 vs.  $\mu \sim 1/n_0 \ (\sqrt{n_0})$ 

When gate-induced flexural scattering dominates.

Modified temperature scaling:

$$\mu \propto T^{-\alpha}$$

$$\alpha$$
=1 Vs.  $1 < \alpha \lesssim 5$ 

When gate-induced flexural scattering dominates transport in the in-plane BG regime.

T< 
$$T_{\rm BG} \approx 57\sqrt{\tilde{n}} \text{ K}$$
  $\tilde{n} = n/10^{12} \text{ cm}^{-2}$ 

Gating strategy:

Double gate setup can be utilized to turn the effect on and off by generating a symmetric or asymmetric potential across the device.

#### **Conclusions:**

- New "Field induced flexural-phonon" scattering mechanism defines a relation between device symmetry and resulting mobility.
- Protecting the planar mirror symmetry is of utmost importance to fully exploit the unique transport properties of graphene.
- T. Gunst, K. Kaasbjerg, M. Brandbyge, "Flexural-phonon scattering induced by electrostatic gating in graphene", Physical Review Letters, **118**, 046601 (2017).
- T. Gunst, T. Markussen, K. Stokbro, M. Brandbyge, "First-principles method for electron-phonon coupling and electron mobility: Applications to 2D materials", Phys. Rev. B, **93**, 035414 (2016).